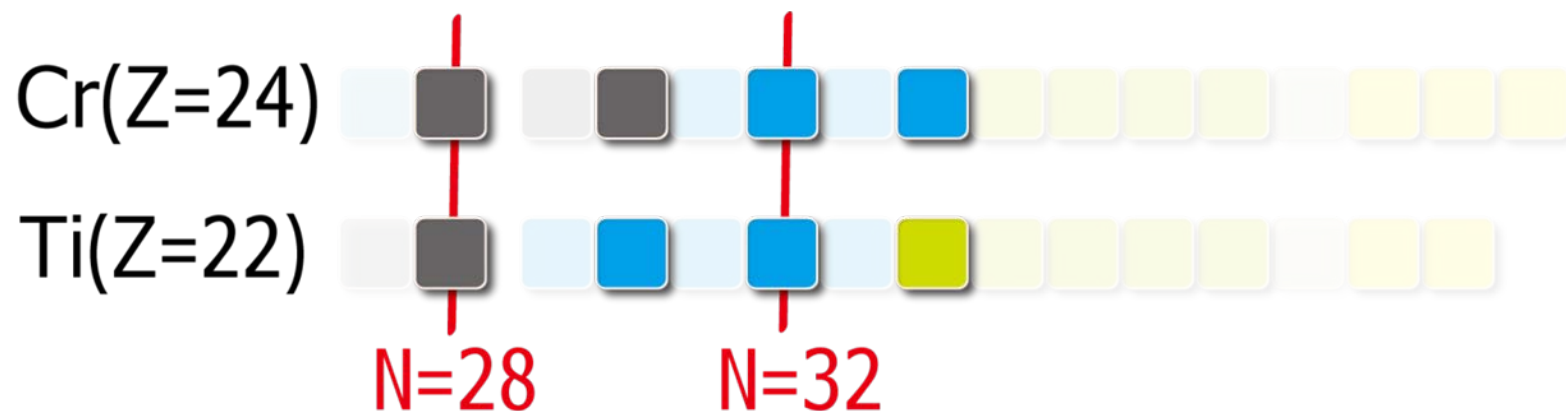


# Shape coexistence and Triaxial deformation in Cr isotopes studied by AMD-HFB model

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## Introduction: $2^+$ and $B(E2)$ of Ti and Cr



Characteristic behavior of  $2^+$  and  $B(E2)$  is reported by exp

- New magic number in  $N \neq Z$  nucleus
- Change of deformation depends on neutron number
- Shape coexistence and triaxial deformation are expected

Shell model : T. Otsuka et al., Eur. Phys. J. A 15, 151 (2002).

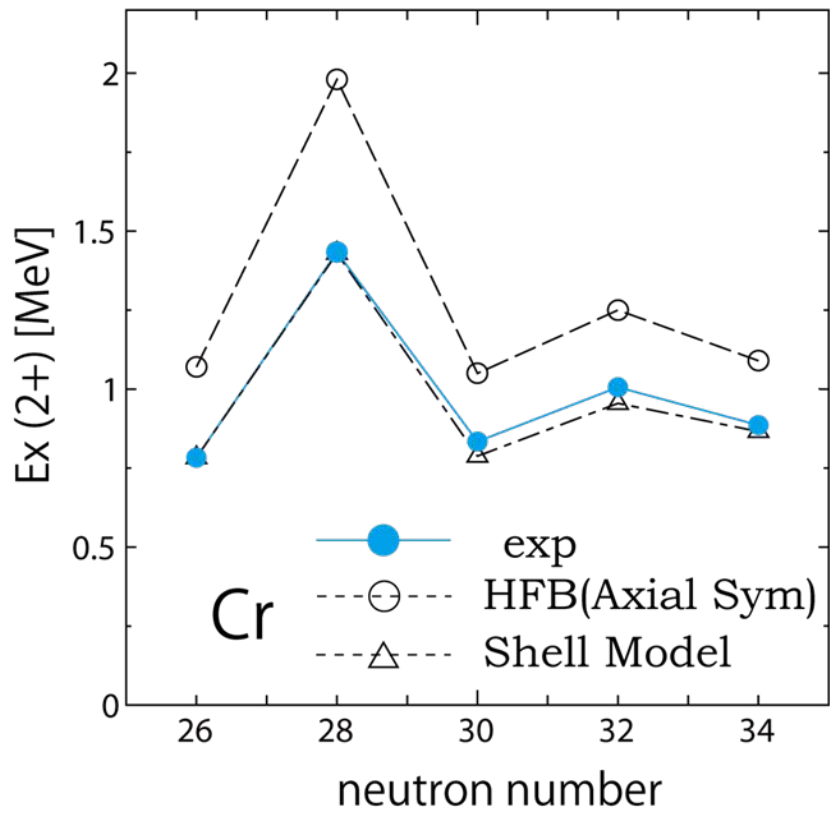
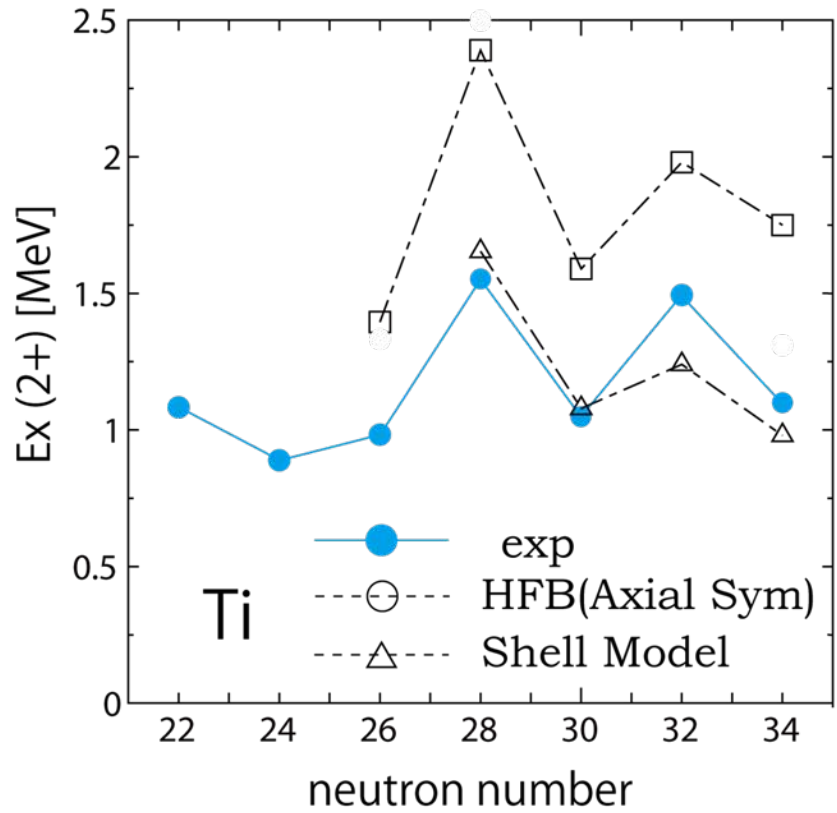
HFB+GCM: T.R. Rodriguez, et. al. PRL99, 062501(2007).

# Introduction: Excitation energy of Ti and Cr 2+ states

- Staggering behavior of observed 2+ energies
- Shell model and HFB calculations reproduces this behavior

Shell model: GXPF1A interaction

HFB: Axial symmetry, Beyond mean-field (N, J projection, GCM)



# Introduction: $B(E2)$ of Ti and Cr

- Staggering behavior of observed  $B(E2)$  (Magicity of  $N=32$ ?)

- Ti

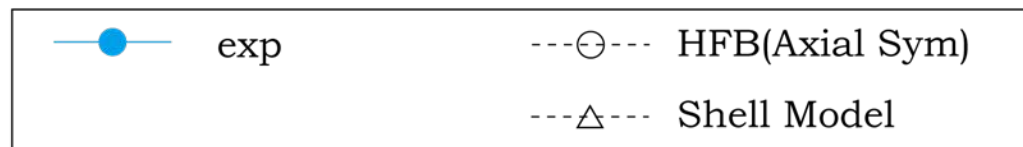
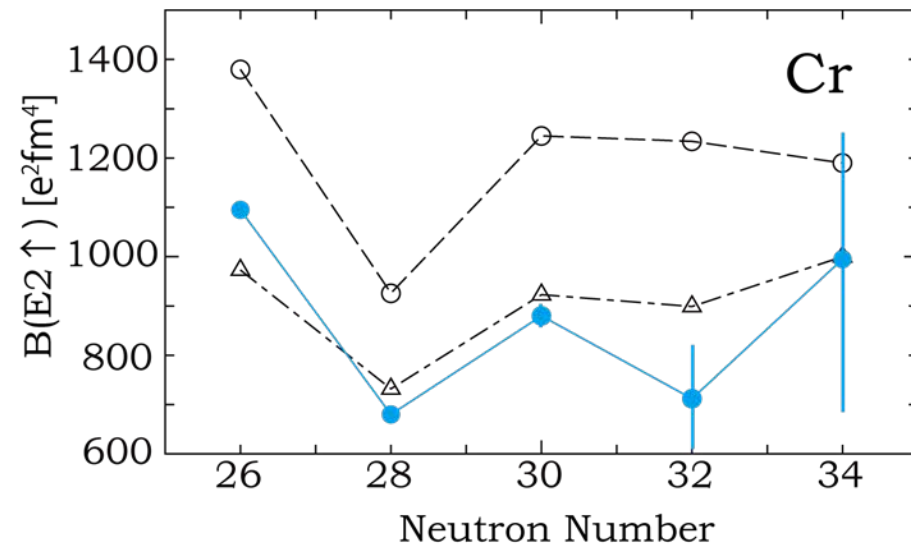
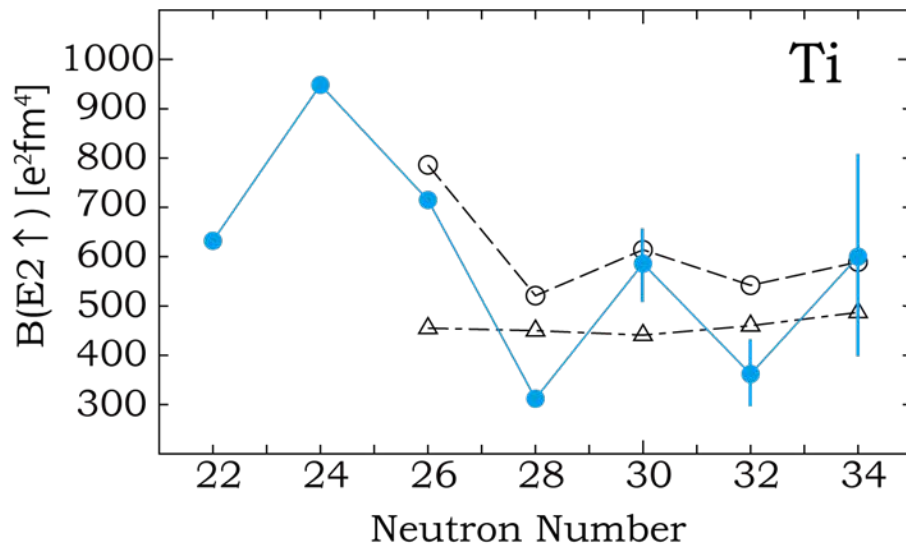
Shell model: Almost constant  $B(E2)$  (constant effective charge)

HFB+GCM: Staggering behavior is reproduced

- Cr

Shell model: Staggering behavior is reproduced

HFB+GCM: Overestimation for all isotopes



# Introduction: Aim of the study

## Structure of Ti and Cr isotopes

Shell model : T. Otsuka et al., Eur. Phys. J. A 15, 151 (2002).

HFB: T.R. Rodriguez, et. al. PRL99, 062501(2007).

- Change of deformation depending on neutron number
- Coexistence of prolate and oblate deformation
- Importance of particle number projection to describe zigzag behavior

- **Possibility of triaxial deformation**



Investigation of deformation of Ti and Cr isotopes

## Construction of theoretical model that can treat

- deformation of nucleus without assumption on its symmetry
- beyond mean-field calculations (N, J projections, GCM,...)
- Clustering phenomena

(extention of Antisymmetrized Molecular Dynamics)

# Framework: AMD+HFB

**Hamiltonian:**  $\hat{H} = \hat{T} + \hat{V}_N + \hat{V}_C - \hat{T}_g$

Effective interaction: Gongy D1S

**Wave function:** HFB type wave function (uses localized Gaussian as basis)

$$|\Psi\rangle = \prod_{i=1}^M \beta_i |0\rangle$$

$$\beta_i = U_{pi}^* a(Z_p) + V_{pi}^* a^\dagger(Z_p)$$

$$\langle \vec{r} | a^\dagger(Z_p) | 0 \rangle = \left( \frac{2\nu}{\pi} \right)^{3/4} \exp \left\{ -\nu \left( \vec{r} - \frac{\vec{Z}_p}{\sqrt{\nu}} \right)^2 + \frac{1}{2} Z_p^2 \right\}$$

$$\{a_i(\vec{Z}_p), a_j^\dagger(\vec{Z}_q)\} = b(\vec{Z}_p, \vec{Z}_q)$$

**Variational calculation:**

- Variation of U and V: ordinary HFB equations (gradient method)
- Variation of basis wave function Z: AMD like equation (equation of motion for wave packets)

**Number and J projection and GCM**

- Numerical Integration is required. Small number of basis wave function reduces CPU time

**AMD**

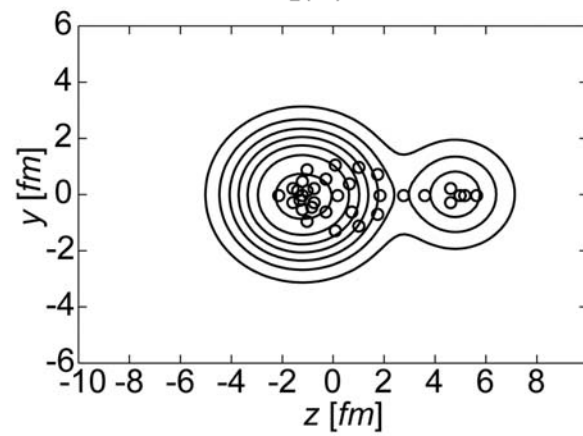
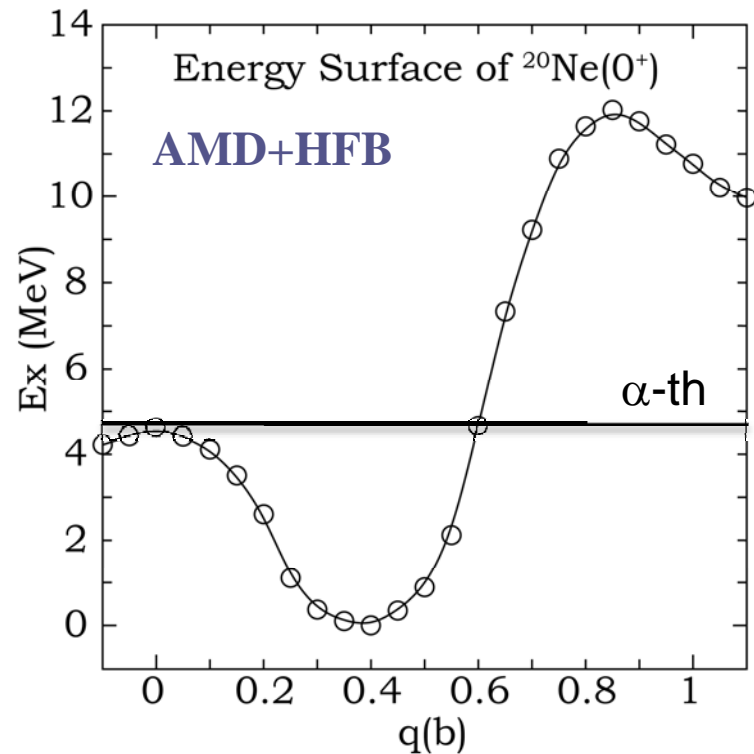
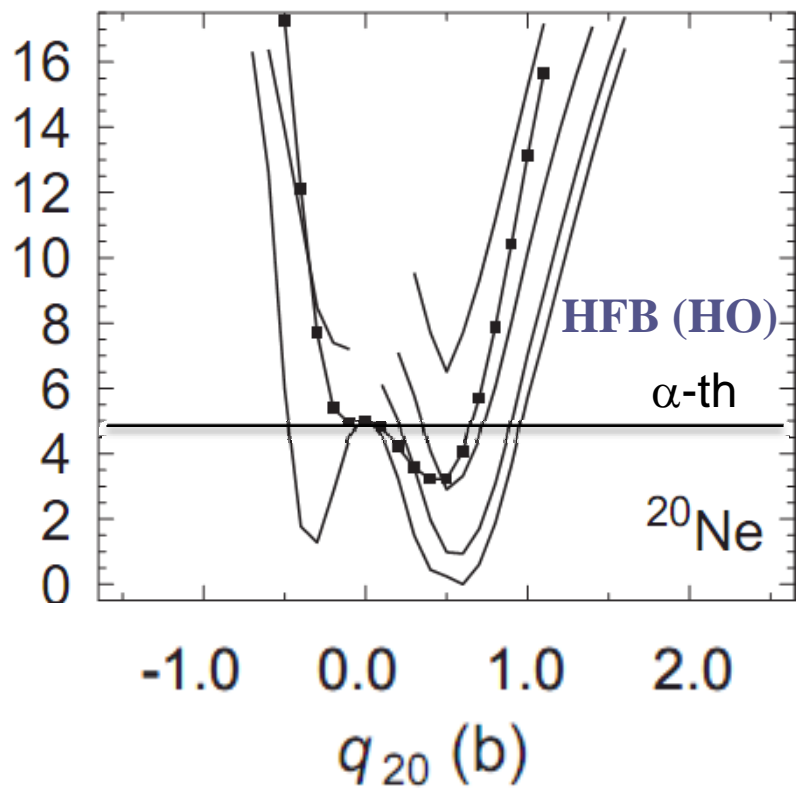
$$|\Psi\rangle = a^+(Z_1) \dots a^+(Z_A) |0\rangle$$

**HFB with H.O. basis**

$$|\Psi\rangle = \prod_{i=1}^M \beta_i |0\rangle,$$

$$\langle r | a^+_i | 0 \rangle = \phi_{njm}(r)$$

# Framework: AMD+HFB



# Results: Energy surface for $\beta$ deformation

## Results of axial symmetric calculation (Projected to $J^\pi=0^+$ and N projection)

- Change of energy surface depending on neutron number

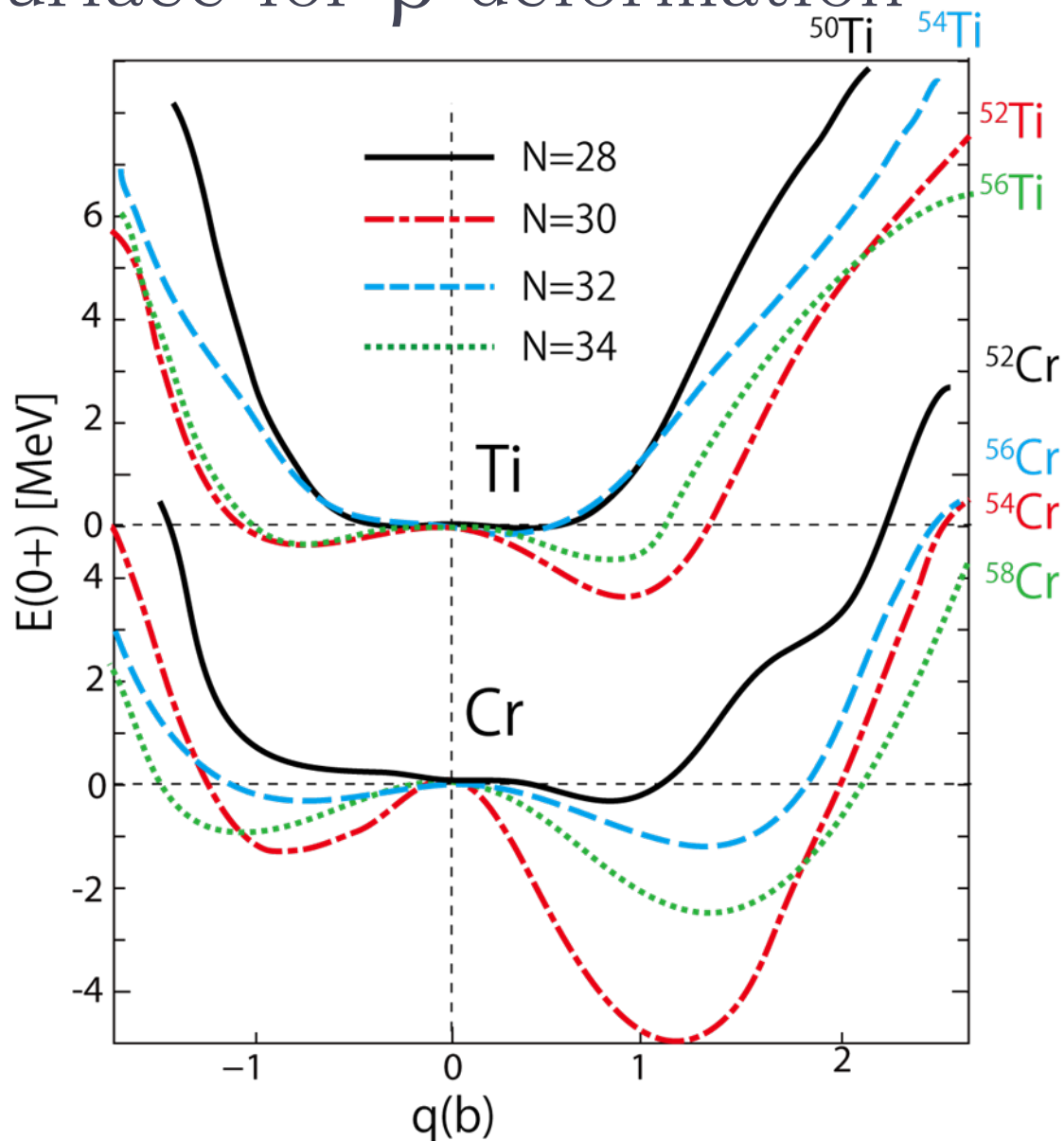
N=30, 34 :

Well developed prolate minimum.  
Coexistence of prolate and oblate minimum.

N=28, 32 :

Shallower prolate minimum  
or spherical

- Larger deformation of Cr than Ti



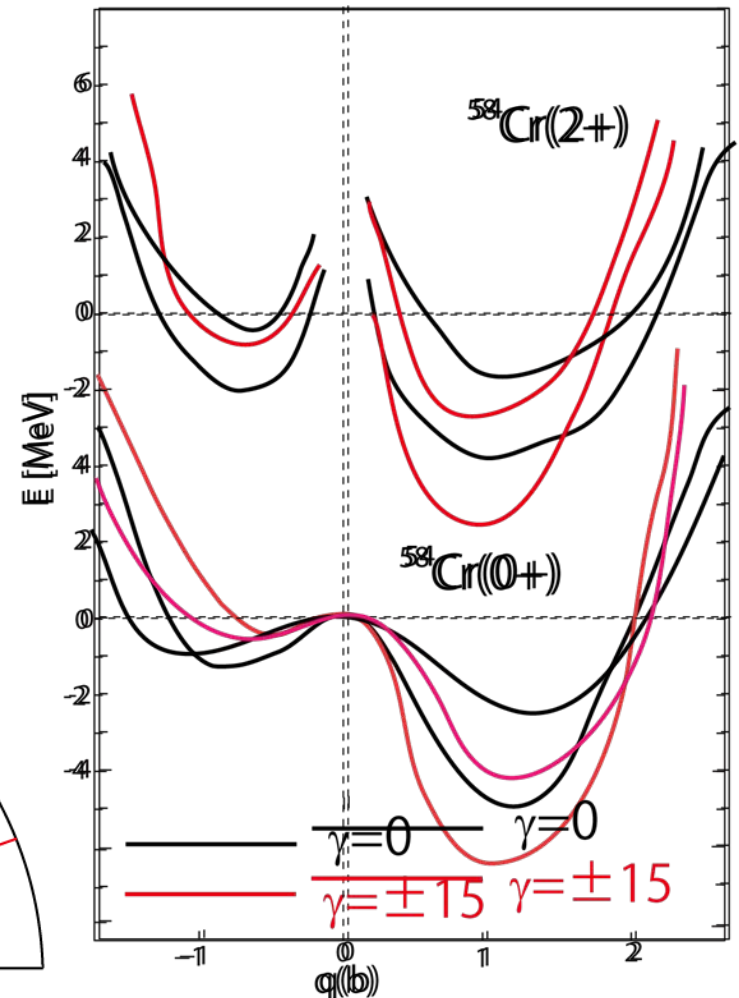
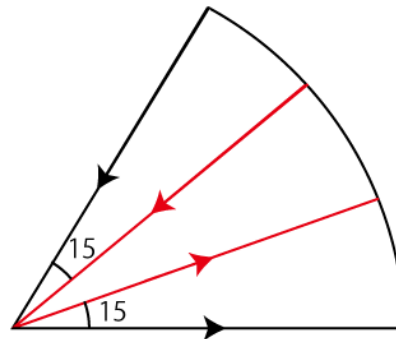


# Results: Energy surface for $\gamma$ deformation

## Results of triaxial calculation $g=15$ and $45$ degree (Projected to $J^\pi=0^+$ and N projection)

- Intrinsic energy (Energy before projection) is higher than axial case
- After the J projection, triaxial state is more deeply bound than axial states in  $^{54}\text{Cr}$  and  $^{58}\text{Cr}$

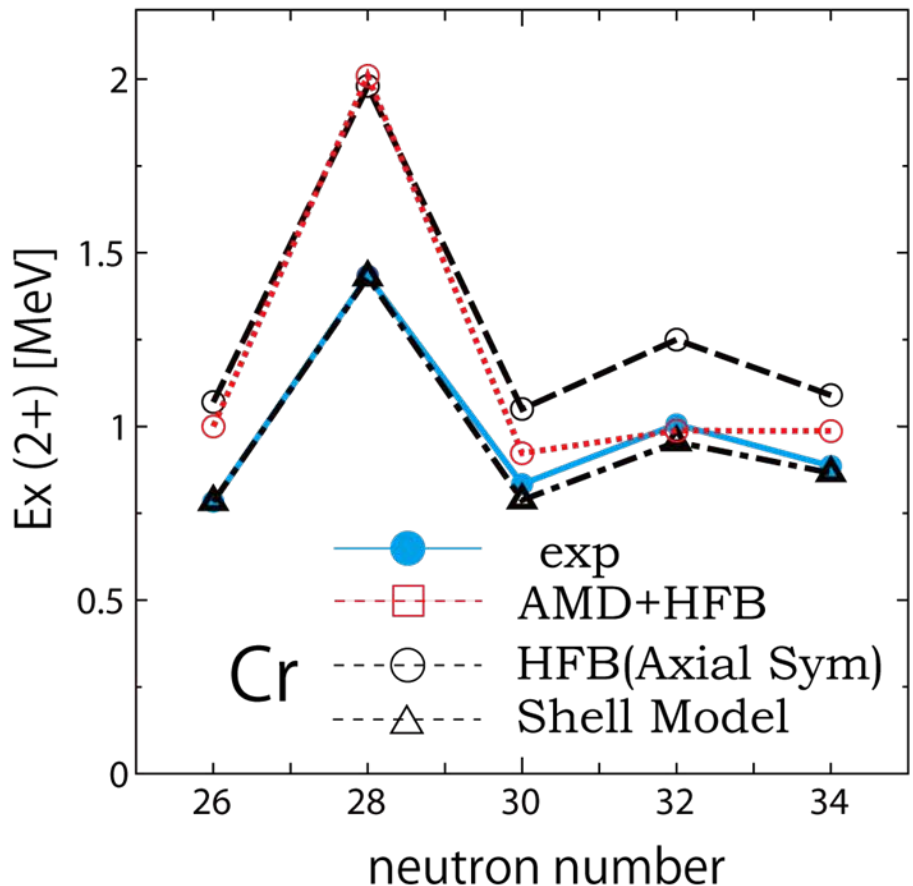
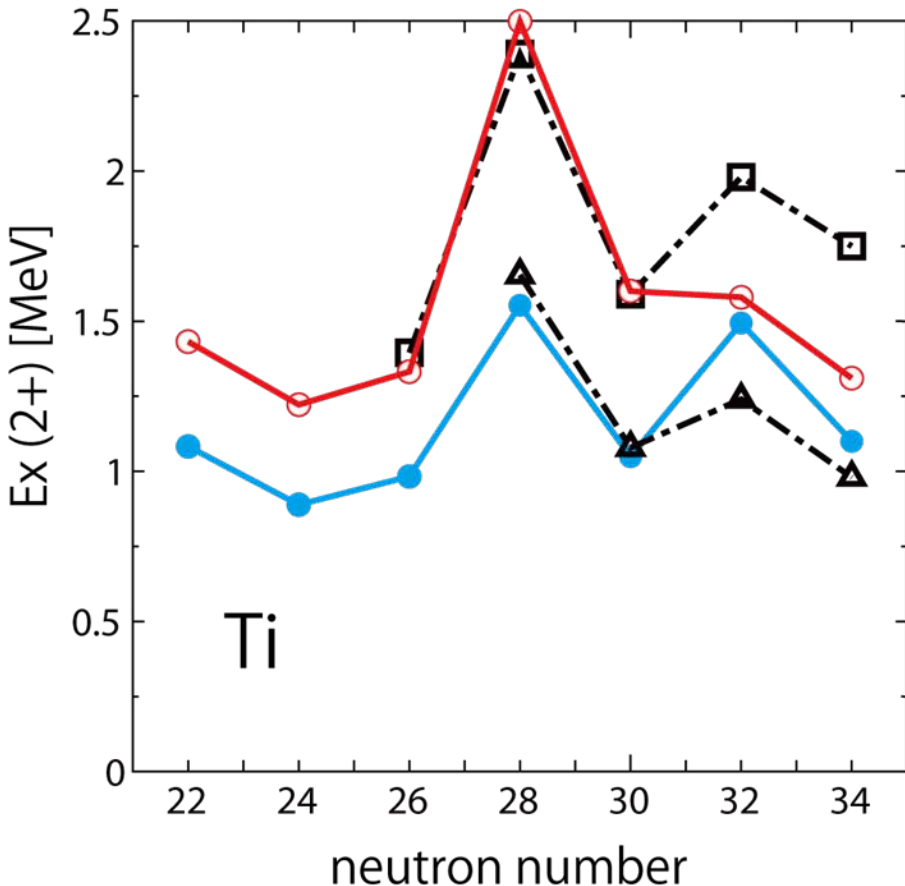
This leads to the mixing of prolate, oblate and triaxial states after GCM calculation



# Results: Excitation energy of $2^+$ states of Ti and Cr

## Results of triaxial calculation (projection and GCM)

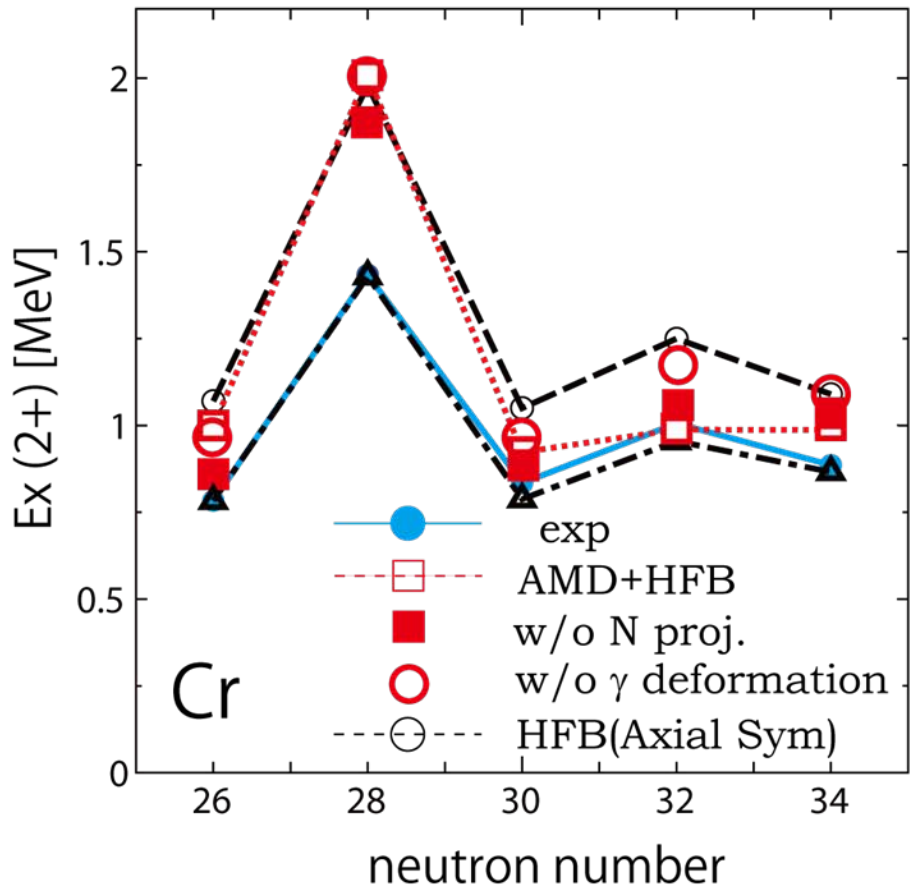
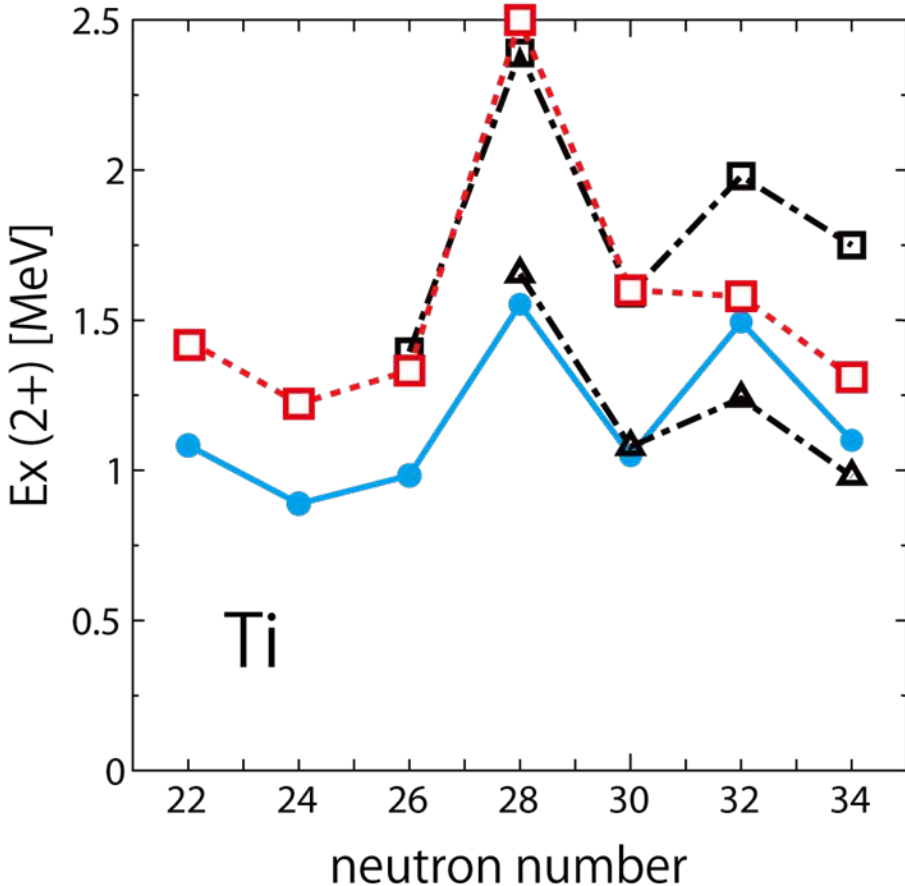
- $2^+$  energies are improved for N=30, 32, and 34 systems
- Triaxial deformation reduces  $2^+$  energies
- Energy gain by the triaxial deformation is rather small.



# Results: Effect of N proj. and triaxality to 2+ energy

## Comparison of axial v.s. triaxial and N projected and unprojected calculations

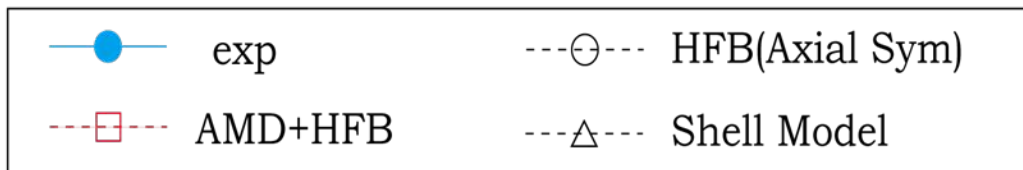
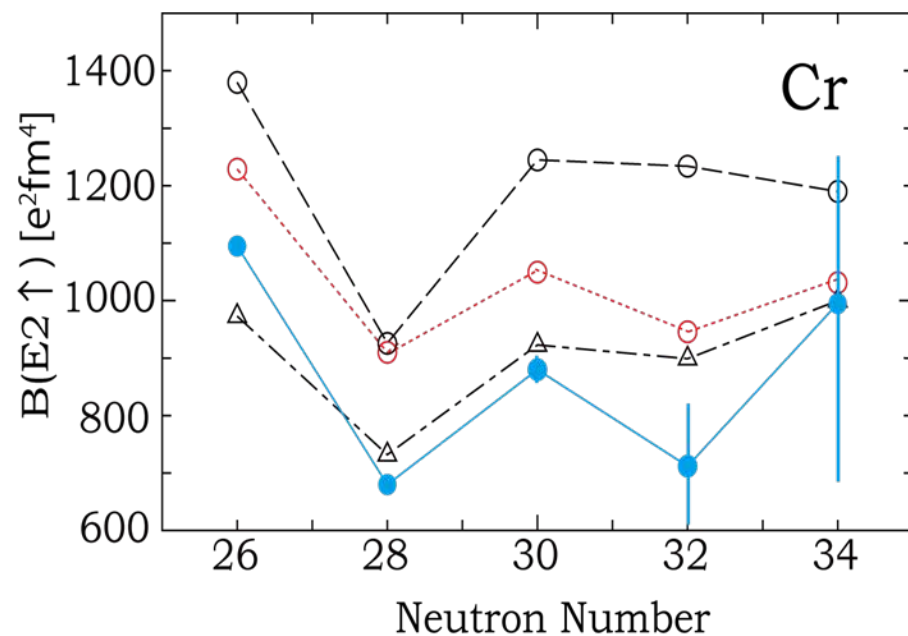
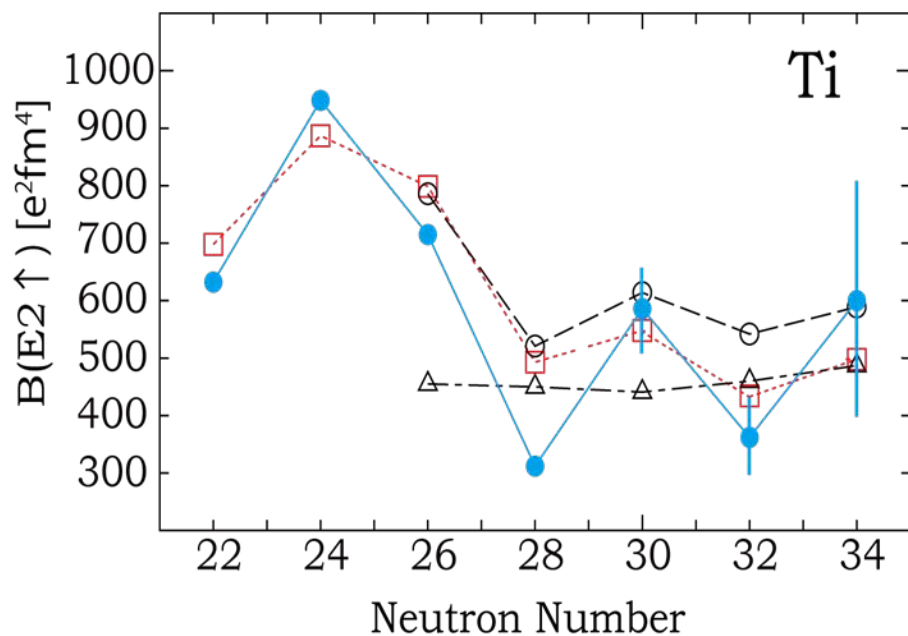
- N projection has weaker effect on 2+ energies
- Triaxial deformation lowers 2+ energies by a 200~300 keV



# Results: $B(E2)$ of Ti and Cr

## Results of triaxial calculation (projection and GCM)

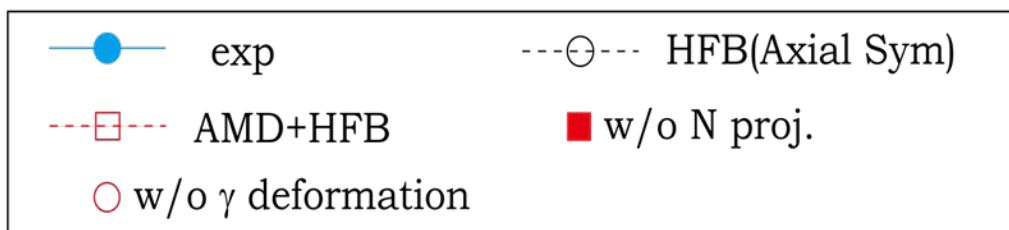
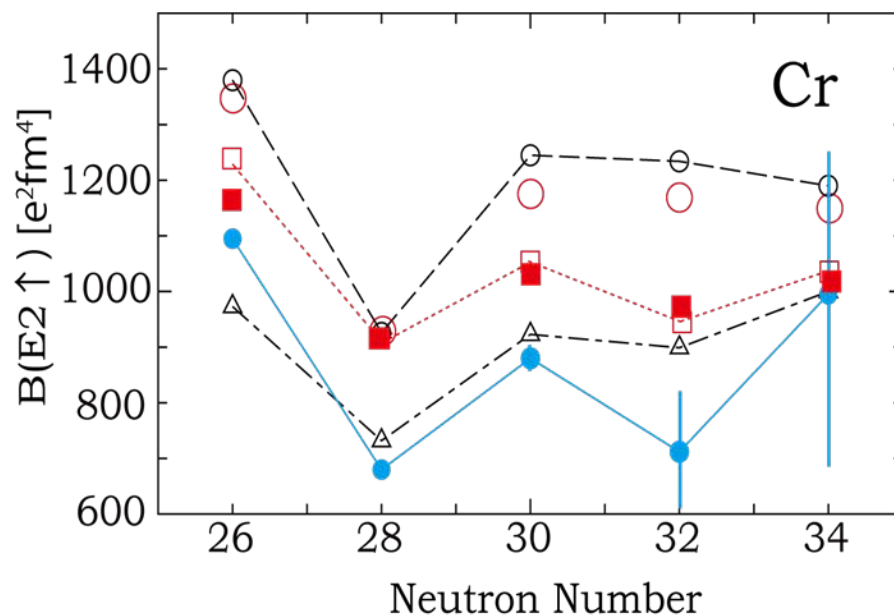
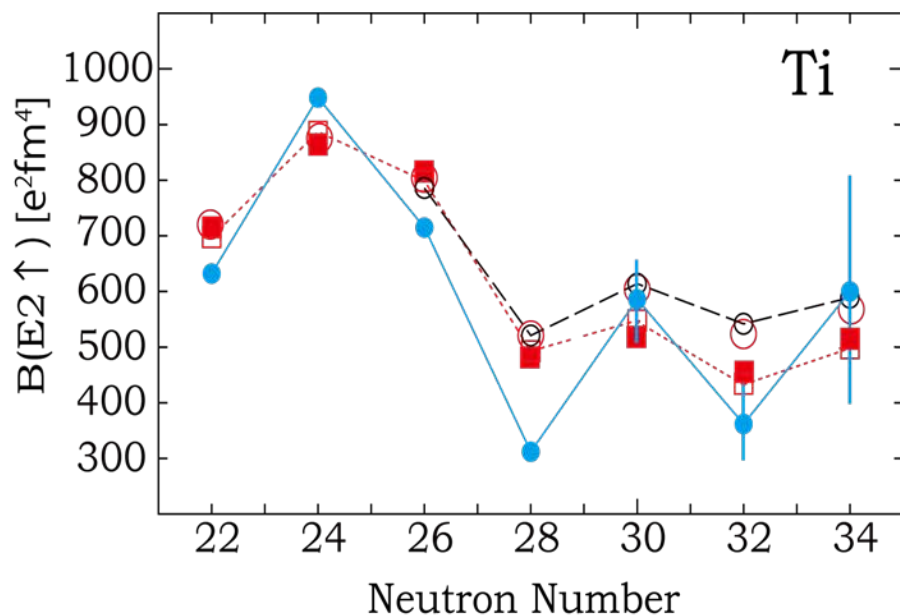
- Reduction of  $B(E2)$  in  $N=30, 32$  and  $34$  System
- Cr isotopes have larger reduction due to their larger triaxial effect that is most apparent in  $^{56}\text{Cr}(N=32)$



# Results: $B(E2)$ of Ti and Cr

## Comparison of axial v.s. triaxial and N projected and unprojected calculations

- N projection does not have large effect
- Triaxial deformation greatly reduces  $B(E2)$  of  $^{54,56,58}\text{Cr}$  (N=30, 32, 34)



# Summary & Outlook

## **Theoretical Framework of AMD+HFB is suggested**

- Deformation of nucleus without assumption on its symmetry
- Beyond mean-field calculations (N, J projections, GCM,...)
- Clustering phenomena  
(extention of Antisymmetrized Molecular Dynamics)

## **Behavior of 2+ and B(E2) of Ti and Cr isotopes**

- Coexistence of prolate and oblate deformation
- Development of triaxial deformation  
(most prominent in  $^{54}\text{Cr}$  and  $^{58}\text{Cr}$ )
- Mixing of prolate, oblate and triaxial wave functions
- This mixing greatly reduces the B(E2) of Cr isotopes ( $N \geq 30$ )

## **Outlook**

Triaxial deformation and shape coexistence in  $^{43}\text{S}$