Magnetism of the antiferromagnetic spin-3/2 dimer compound CrVMoO₇ having an antiferromagnetically ordered state[†]

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In a magnetically ordered state, two types of magnetic excitations exist: gapless transverse-mode (Nambu-Goldstone mode) and gapped longitudinalmode (amplitude Higgs mode) excitations. The former is well known as spin wave excitation. The longitudinal-mode (L-mode) was observed by inelastic neutron scattering experiments in the pressure-induced magnetically ordered state of TlCuCl₃, which is a threedimensional (3D) interacting antiferromagnetic (AF) spin-1/2 dimer compound.^{1–3} While the L-mode has weak intensity and spontaneously decays into a pair of transverse-modes, it is well defined in the ordered state in the vicinity of the quantum critical point for 3D systems.⁴

As for low-dimensional systems, it is difficult to observe the L-mode in longitudinal susceptibility by inelastic neutron scattering, since the longitudinal susceptibility exhibits an infrared singularity that can obscure the amplitude peak at a finite energy. In terms of scalar susceptibility, however, the L-mode can be well defined both in 2D and 3D systems. In magnetic systems, the L-mode was actually observed by Raman scattering experiments (which can measure the scalar susceptibility) in the pressure-induced ordered state of KCuCl₃ and in the magnetic-field-induced ordered state of TlCuCl₃.⁵⁻⁷)

According to the results of theoretical investigations on interacting AF spin-cluster compounds, the L-mode excitations can be observed in the antiferromagnetically ordered state that appears on cooling under atmospheric pressure and zero magnetic field.⁸⁾ A shrinkage of the ordered magnetic moments by quantum fluctuations leads to a large intensity of the L-mode excitations. If the ground state of the isolated spin cluster is a spin-singlet state, the shrinkage of ordered moments can be expected in an ordered state that is generated by introducing intercluster interactions. We expect that an interacting AF spin-3/2 dimer model can be applied to the compound CrVMoO₇, judging from its crystal structure (Fig. 1).^{9,10}

We studied the magnetic properties of CrVMoO₇ in powder form, using magnetization, specific heat, electron spin resonance, neutron diffraction, and inelastic

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Fig. 1. The unit cell of CrVMoO₇. An AF spin-3/2 dimer is formed by two neighboring Cr^{3+} ions with a distance of 3.01 Å.

neutron scattering measurements. An antiferromagnetically ordered state appeared below $T_{\rm N} = 26.5 \pm 0.8$ K. The magnetic susceptibility at high temperatures was close to that calculated for the isolated AF spin-3/2 dimer with an intradimer interaction of $J = 25 \pm 1$ K and $g = 1.92\pm0.02$. We were able to explain the magnetization curves on the basis of the interacting AF spin-3/2 dimer model with an effective interdimer interaction of $J_{\rm eff} = 8.8 \pm 1$ K. The magnitude of the ordered moment was $0.73(2)\mu_{\rm B}$, which is much smaller than the classical value of $\sim 3\mu_{\rm B}$. Using inelastic neutron scattering measurements, the magnetic excitations were observed, and the results were qualitatively explained on the basis of the interacting AF spin-3/2 dimer model.

In conclusion, $CrVMoO_7$ is a rare spin-dimer compound that shows an antiferromagnetically ordered state under atmospheric pressure and zero magnetic field. Though we could not confirm the L-mode in this study, the magnetic excitations of L-mode would be observable in *single crystalline* CrVMoO₇.

References

- 1) Ch. Rüegg et al., Phys. Rev. Lett. 100, 205701 (2008).
- 2) P. Merchant et al., Nat. Phys. 10, 373 (2014).
- 3) M. Matsumoto et al., Phys. Rev. B 69, 054423 (2004).
- 4) I. Affleck, G. Wellman, Phys. Rev. B 46, 8934 (1992).
- 5) H. Kuroe et al., J. Phys.: Conf. Ser. 400, 032042 (2012).
- M. Matsumoto *et al.*, J. Phys. Soc. Jpn. **77**, 033702 (2008).
- 7) H. Kuroe et al., Phys. Rev. B 77, 134420 (2008).
- M. Matsumoto *et al.*, J. Phys. Soc. Jpn. **79**, 084703 (2010).
- 9) X. Wang et al., Inorg. Chem. 37, 3252 (1998).
- 10) K. Knorr *et al.*, Eur. J. Solid State Inorg. Chem. **35**, 161 (1998).