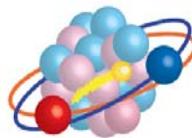


不安定核

—核構造の新天地、核構造とハドロン物理の架け橋—



Takaharu Otsuka
University of Tokyo / RIKEN / MSU

T. Suzuki
M. Honma
Y. Utsuno
Y. Akaishi
H. Grawe

Nihon U.
U. Aizu
JAEA
RIKEN
GSI

A. Schwenk
B. Holt

TRIUMF
TRIUMF

アウトライン

1. 不安定核の物理を見る視点
 - パラダイムシフト
 - 中性子過剰核で浮かび上がる核力の側面
2. テンソル力と殻進化の概観
3. 有効核力と Extended Weinberg Ansatz
4. 3体力が決めるドリップライン、不安定核が明かす3体力の謎
5. まとめと展望

30年前頃の核物理

メゾスコピック系への応用: メタルクラスター、など

核物理としてはやや閉塞感

自分としては「応用」に活路を見出すことには抵抗感

15年前頃の核物理

中性子ハローの発見: ドリップラインこそがフロンティア

原子核の半径は $A^{1/3}$ で決まるというパラダイム
(密度の飽和性)の変更

科学の革新的進歩にはパラダイムシフトが伴う

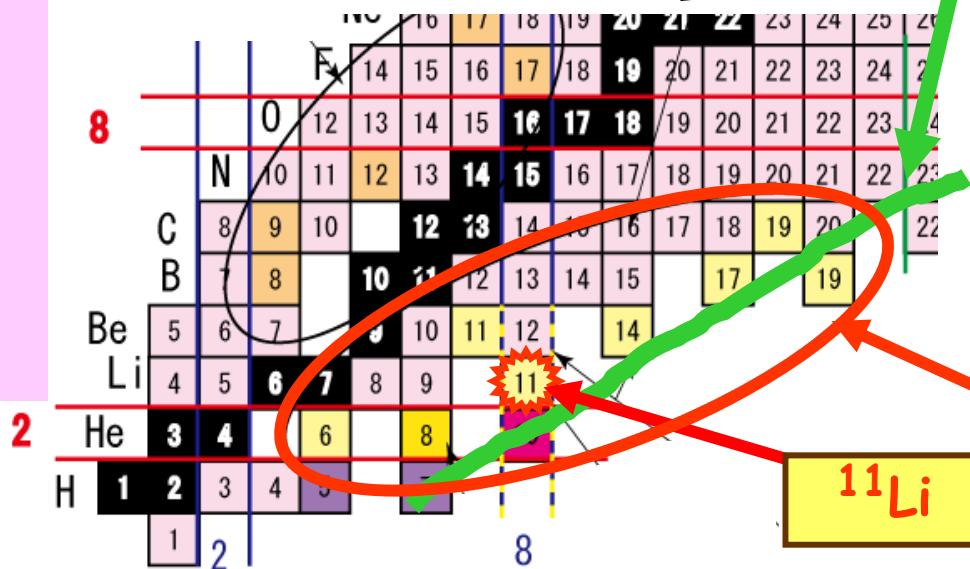
(Thomas Kuhn)

パラダイムシフトはそうそう起こるものではないが、他にも
あり得るだろうか？

Studies on exotic nuclei in the 80~90's

Left-lower part of the Nuclear Chart

Proton number →



Neutron number →

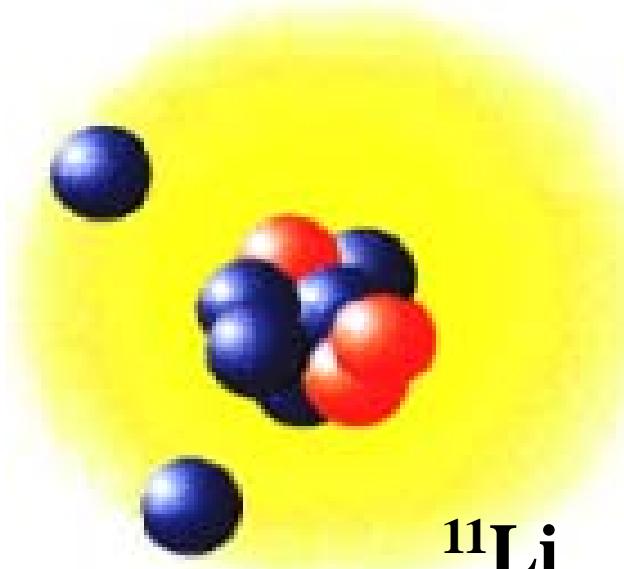
Stability line and drip lines are not so far from each other
→ Physics of loosely bound neutrons, e.g., halo while other issues like ^{32}Mg

neutron halo

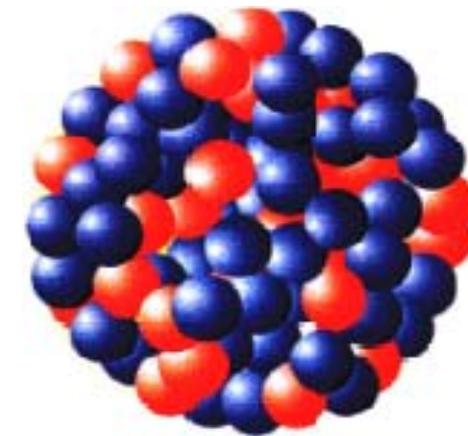
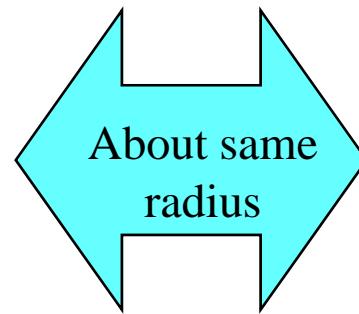
- A nuclei (mass number)
- stable
- exotic
- -- with halo

Neutron halo

Strong tunneling of loosely bound excess neutrons



^{11}Li



^{208}Pb

VOLUME 55, NUMBER 24

PHYSICAL REVIEW LETTERS

Measurements of Interaction Cross Sections and Nuclear Radii in the

I. Tanihata,^(a) H. Hamagaki, O. Hashimoto, Y. Shida, and N. Yosh

Institute for Nuclear Study, University of Tokyo, Tanashi, Tokyo 188, Japan

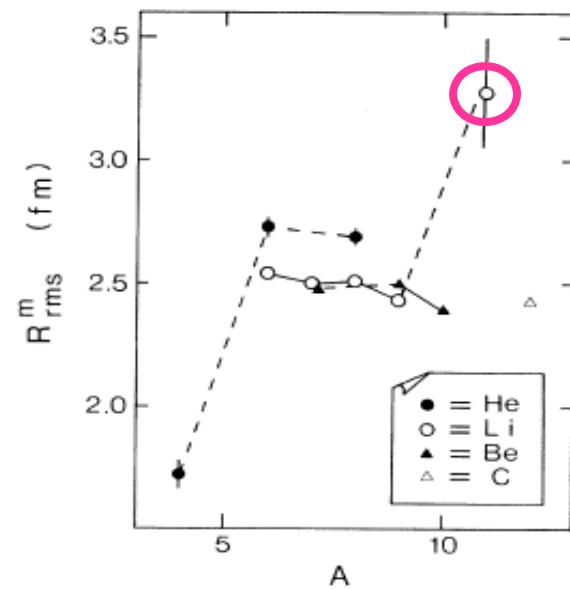
K. Sugimoto,^(b) O. Yamakawa, and T. Kobayashi

Nuclear Science Division, Lawrence Berkeley Laboratory, University of California, Berkeley

and

N. Takahashi

14:46

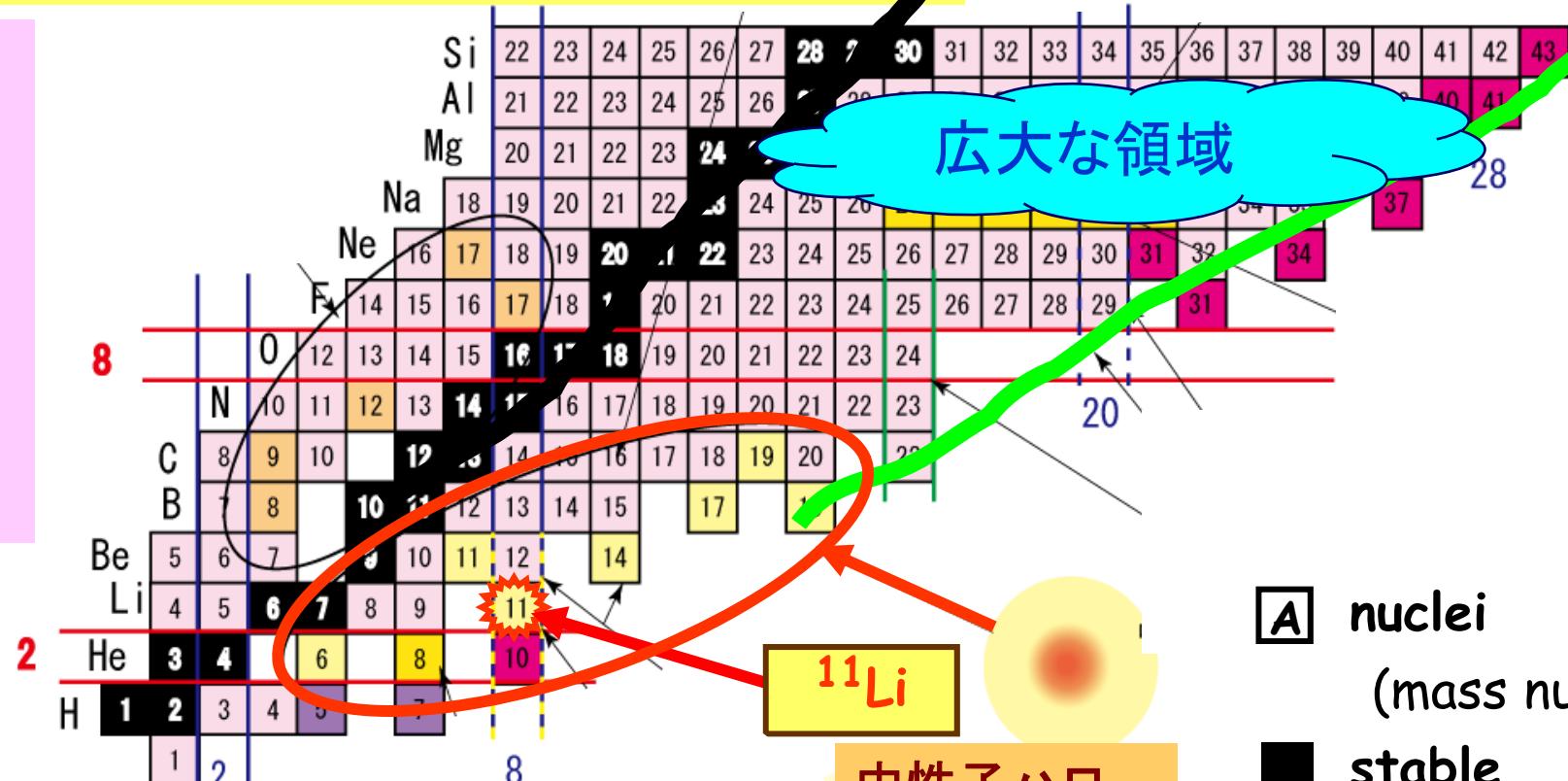


21世紀になり、第3世代 RI ビーム加速器が始まるとともに β 安定線と、ドリップラインの間の多くの不安定原子核が研究できるようになってきた

ベータ安定線

ドリップライン

↑ proton number



Neutron number →

中性子ハロー

- A nuclei
(mass number)
- stable
- exotic
- Riken's work

このように核子数を大きく変えた場合に原子核構造を変える要素の一つとして、Monopole Interaction による Effective Single-Particle Energy の変化があげられる（殻進化、Shell Evolution）

Effective single particle energy

- Monopole part of the NN interaction

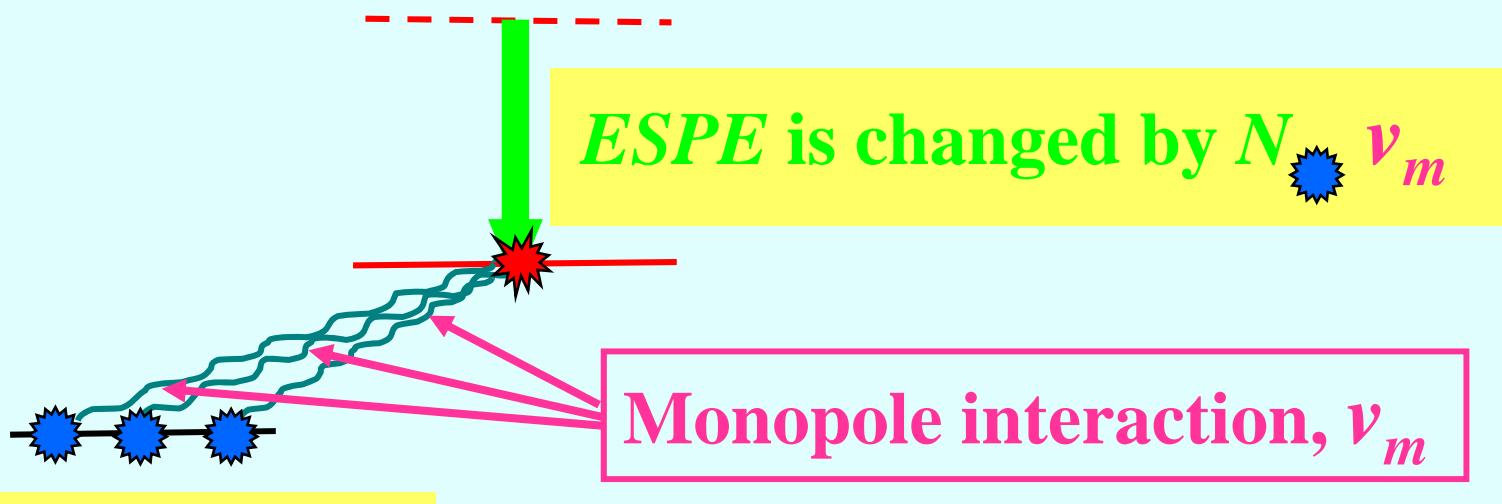
$$V_{ab}^T = \frac{\sum_J (2J+1) V_{abab}^{JT}}{\sum_J (2J+1)}$$



Angular averaged interaction

2核子の軌道運動の間の相対角度について
一般の相互作用を平均したもの

- Effective single-particle energy (*ESPE*)



$N \star$ particles

このような linearity は monopole interaction だけ。
他の multipole interaction は相手(上の図で●)の軌道が全部埋まると効果はゼロになる (trace=0)。一方、monopole にはしばしば 10 のオーダーの係数がかかる。

この linearity は、中性子数だけ(又は陽子数だけ)を変えることができる不安定核の物性論では大変重要な役割。
(安定核だけをやっていた時には顕著でなかった。)

問：殻構造(象徴的には魔法数)は変わり得るか？

From undergraduate nuclear physics,

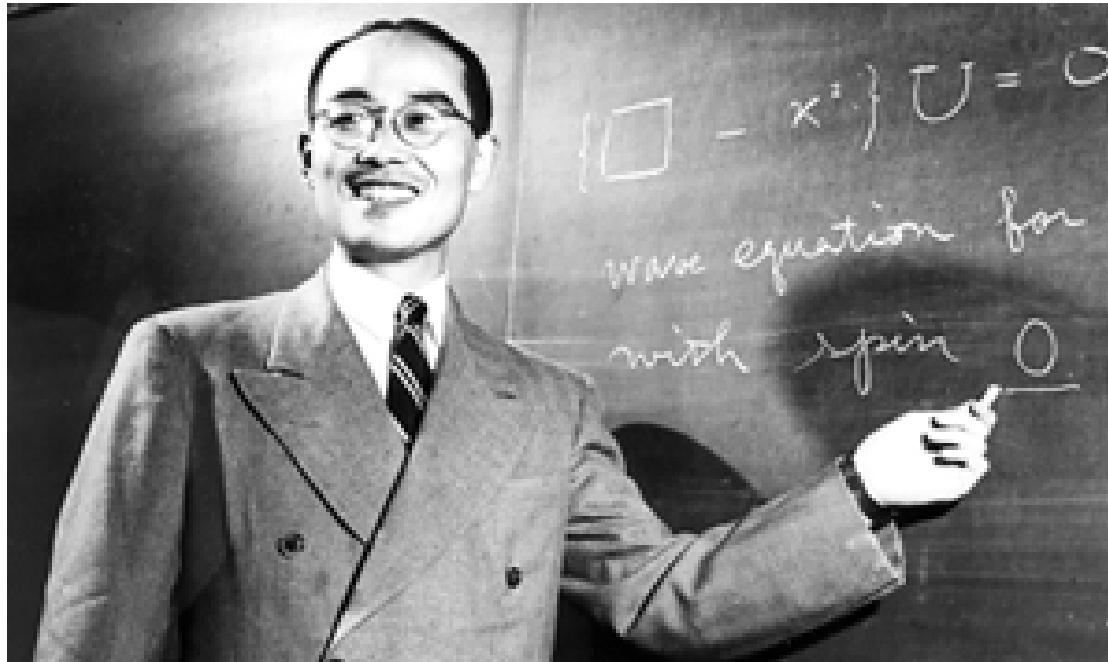
density saturation
+ short-range $\Lambda\Lambda$ interaction
+ spin-orbit splitting

→ Mayer-Jensen's magic number
with rather **constant gaps**
(except for gradual A dependence)

→ 上にないものが効かない限り変わらない

→ だから変わらない(ギャップは変わらない) 10年前は私も

鍵=テンソル力（1次の効果）

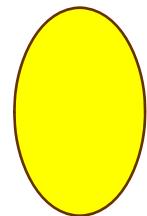


π中間子 1 個の交換 ~ テンソル力

複数のパイオンの交換 → 有効中心力 (σ 中間子)
原子核の結合エネルギーの主要部分

Monopole effects due to the tensor force

- An intuitive picture -

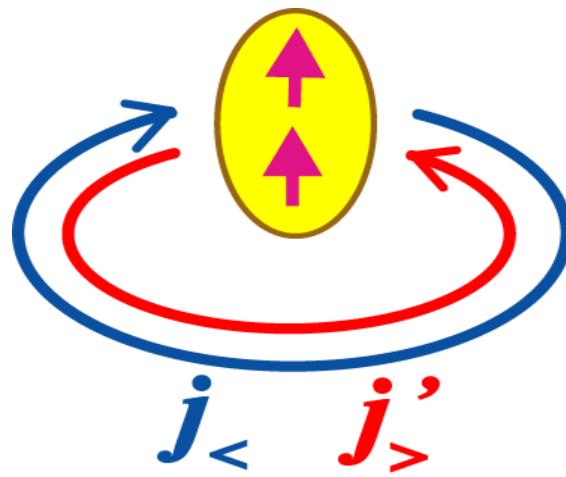


wave function of relative motion



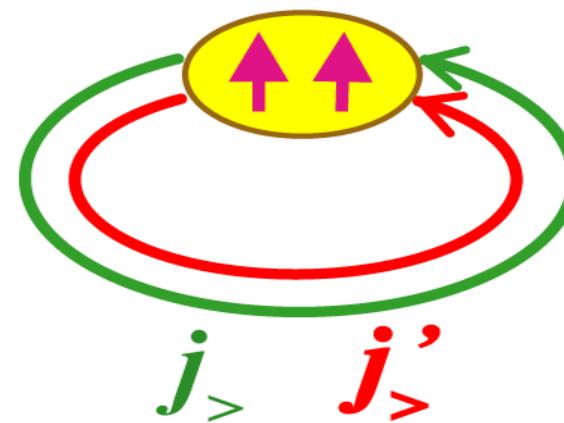
spin of nucleon

large relative momentum



attractive

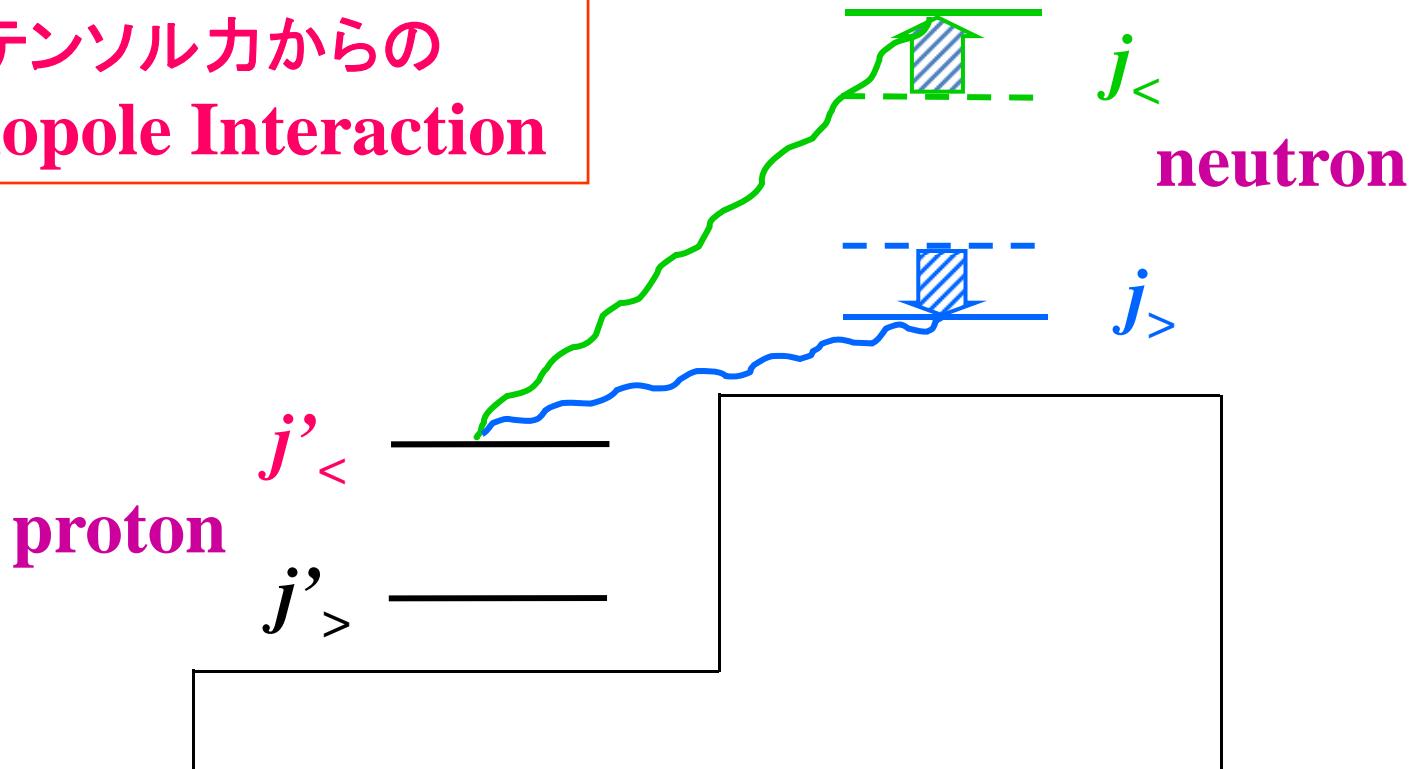
small relative momentum



repulsive

$$j_{>} = l + \frac{1}{2}, \quad j_{<} = l - \frac{1}{2}$$

テンソル力からの
Monopole Interaction



テンソル力の monopole interaction が満たす恒等式

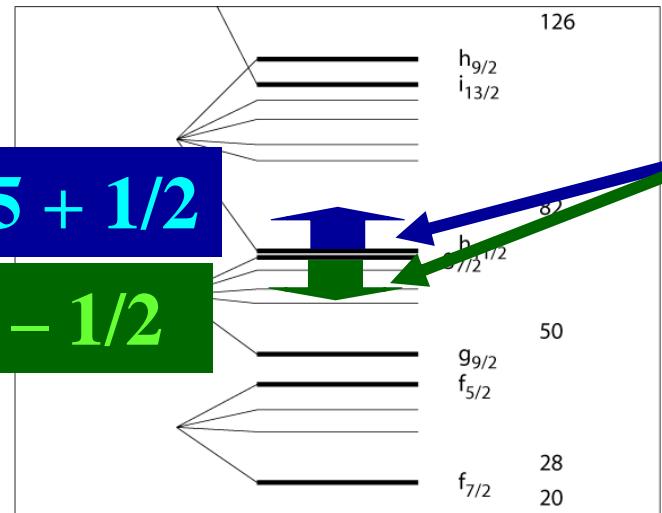
$$(2j_{>}+1) \ v_{m,T}^{(j' j_{>})} + (2j_{<}+1) \ v_{m,T}^{(j' j_{<})} = 0$$

$v_{m,T}$: monopole strength for isospin T

51Sb case

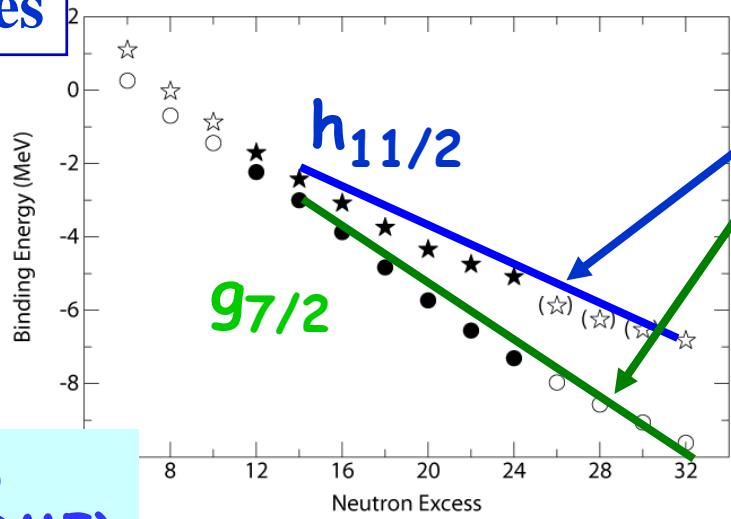
$1h_{11/2} : 11/2 = 5 + 1/2$

$1g_{7/2} : 7/2 = 4 - 1/2$



Opposite monopole effect from tensor force with neutrons in $h_{11/2}$.

Z=51 isotopes



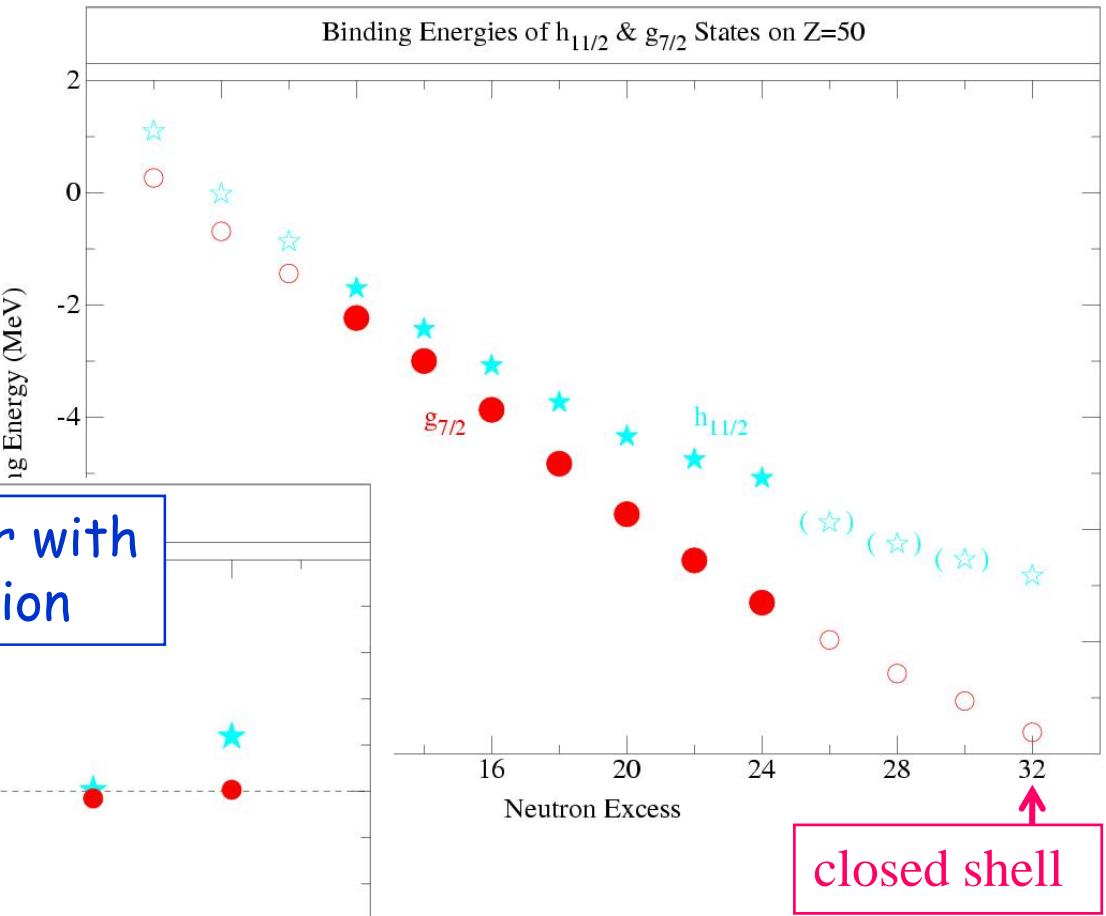
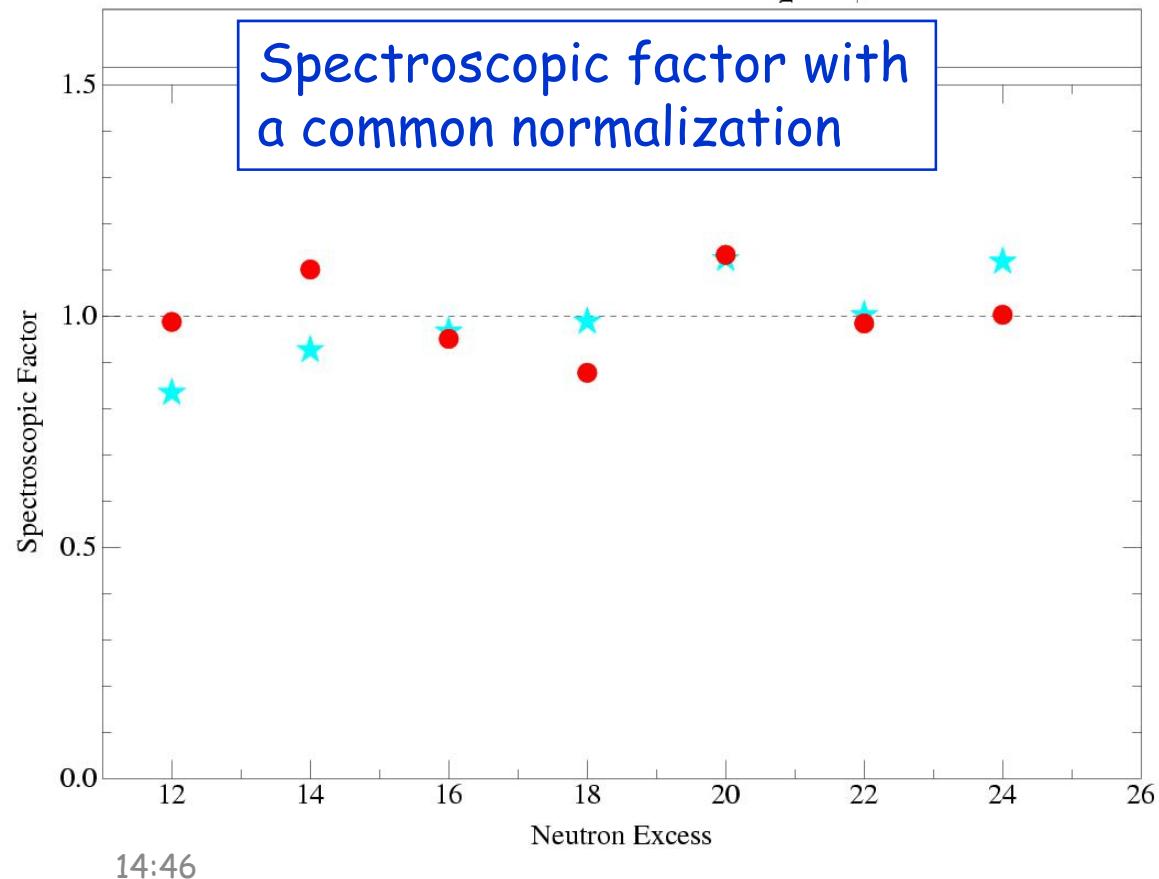
Tensor by
 $\pi + \rho$ meson exchange

+ common effect
(Woods-Saxon)

1h_{11/2} neutrons

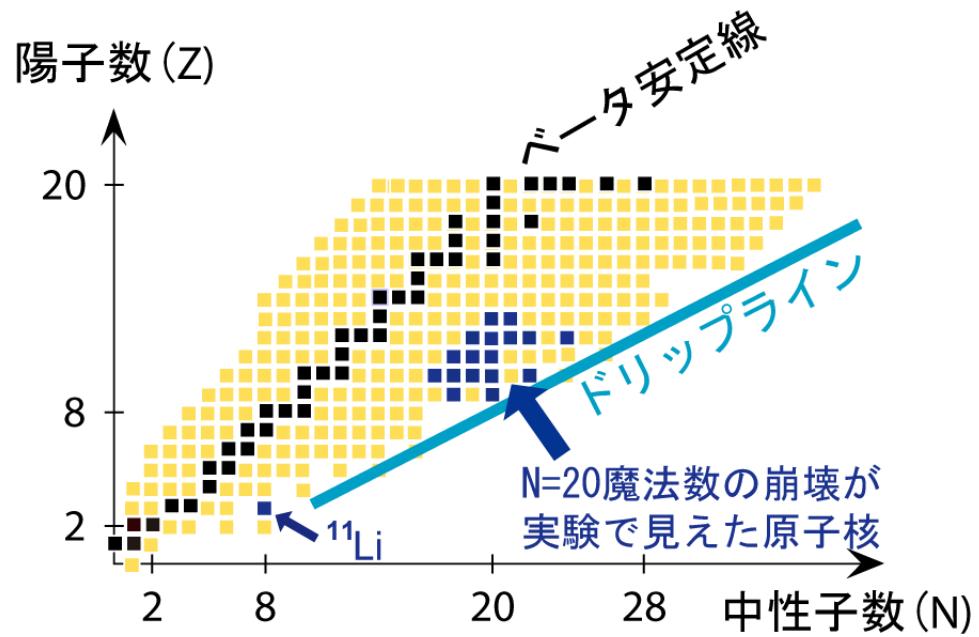
どの平均場理論でも
(Skyrme, Gogny, RMF)
説明できなかった

These states are
of single-particle nature
to a similar extent.

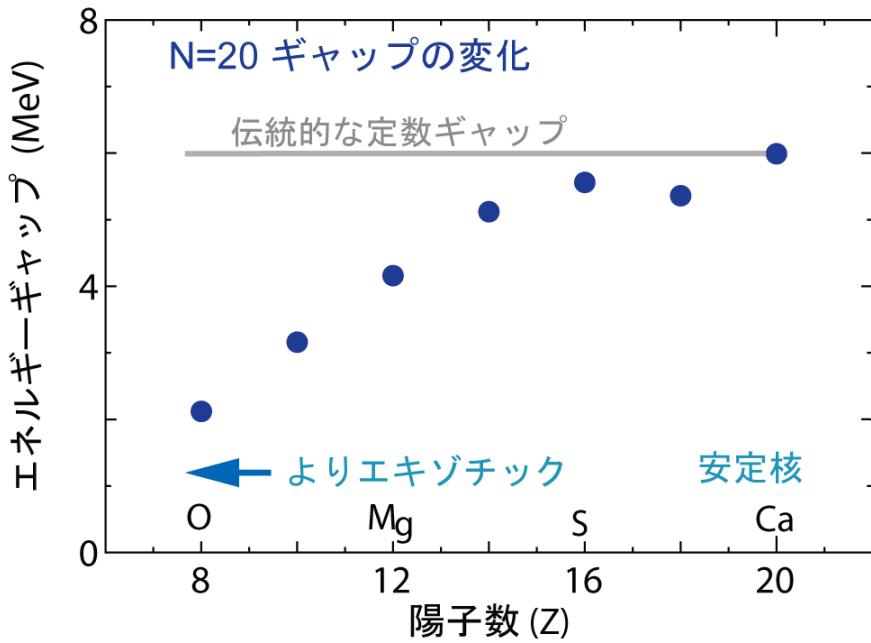


Correlation effects do not
seem to change the story.

From the lecture note
by Schiffer for the 2008
CNS-EFES summer school



Island of Inversion の
隠された姿は？



テンソル力の monopole interaction の検証で
最近判明した例：

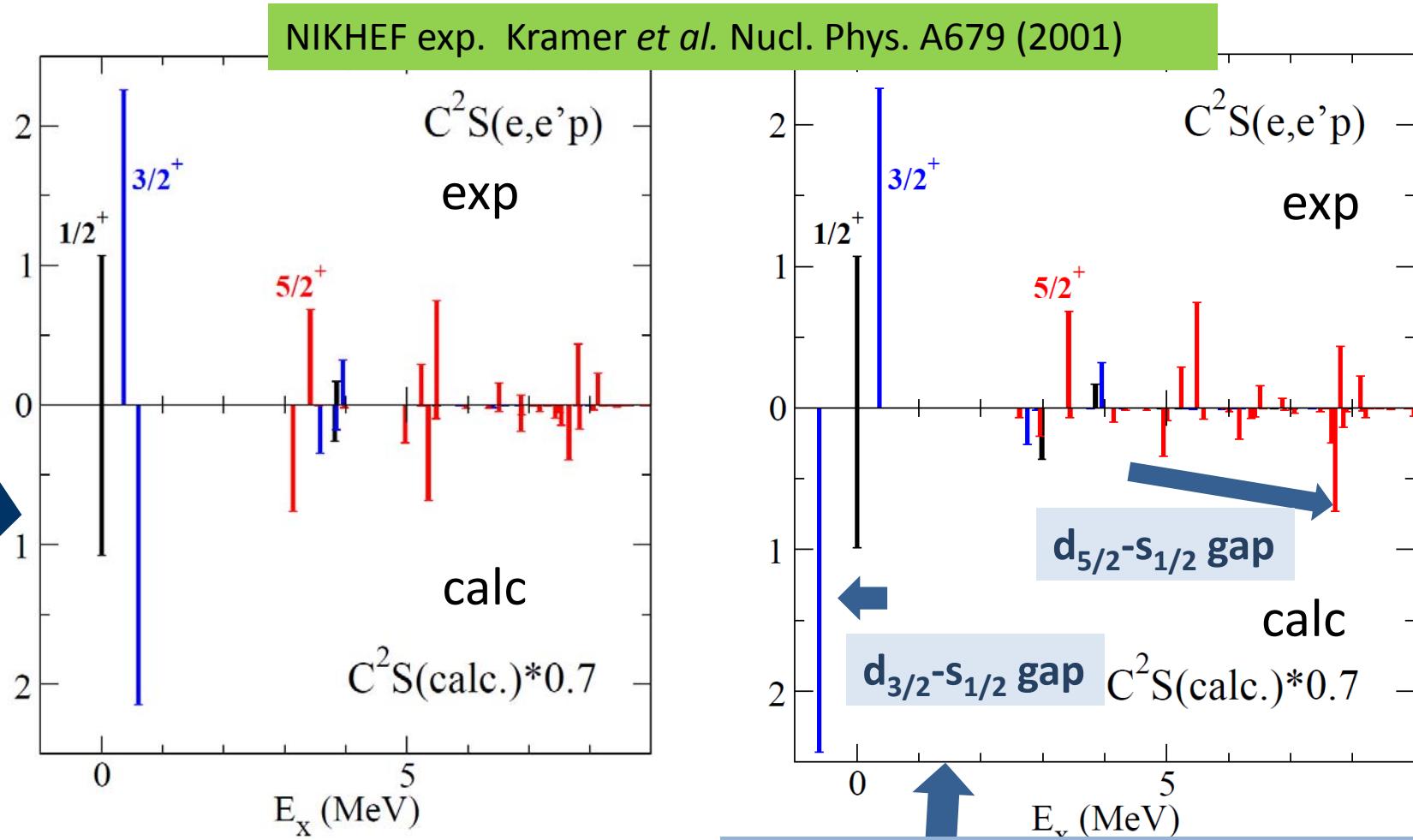
... ^{48}Ca ($e, e' p$)

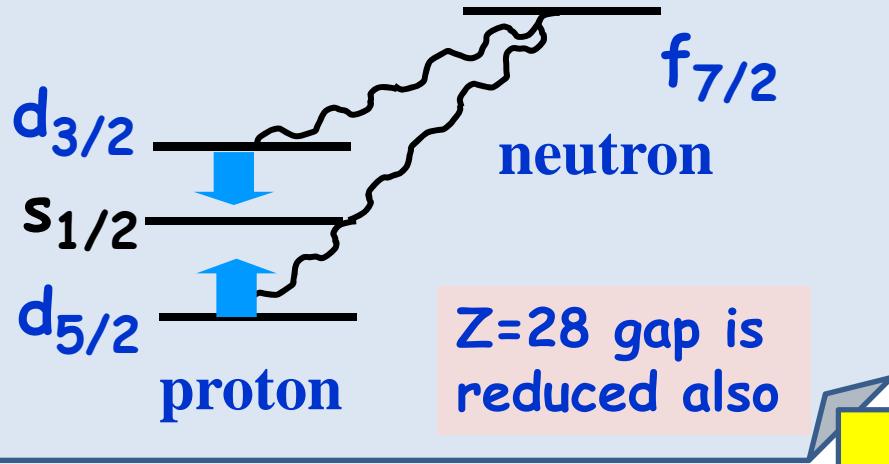
テンソル力によるスピン軌道分岐の変化そのものを
見た(含むfragmentation)最初の例？

電子散乱実験 ⇒ 不安定核での殻進化
J-Labなどができるといいのだが
RIビームができるのはだいぶ先の話？

Spectroscopic factor for -1p from ^{48}Ca : *probing the change of spin-orbit splitting*

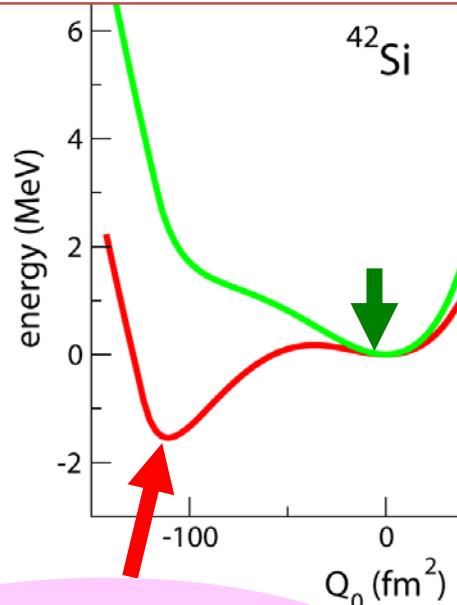
The same interaction as the one for ^{42}Si





New Shell Model Int.
based on "Anatomy"

Potential Energy Surface



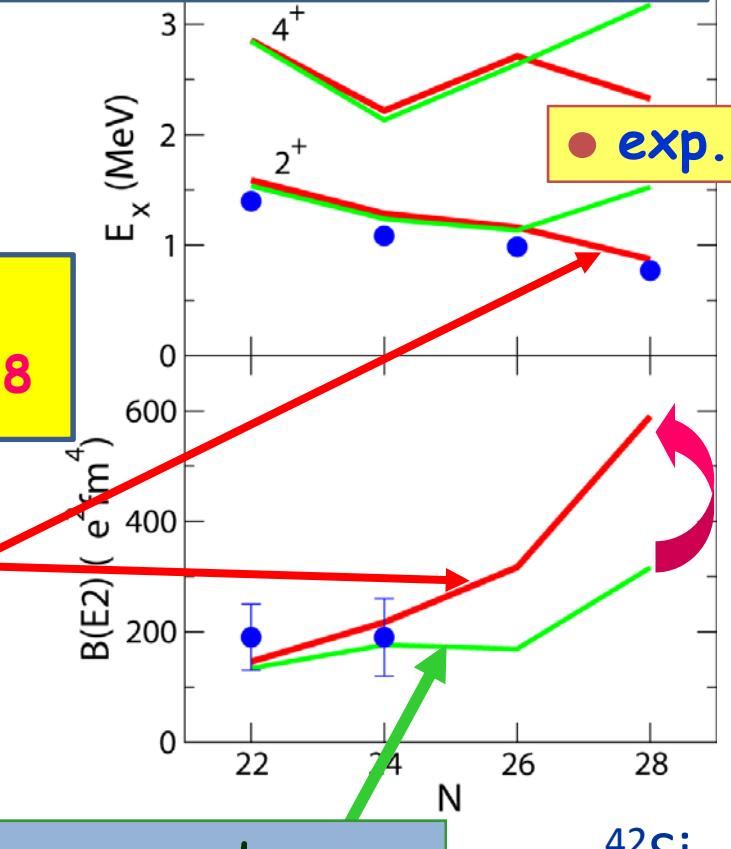
$^{42}_{\text{Si}} \text{Si}_{28}$

full

Tensor force removed
from cross-shell interaction

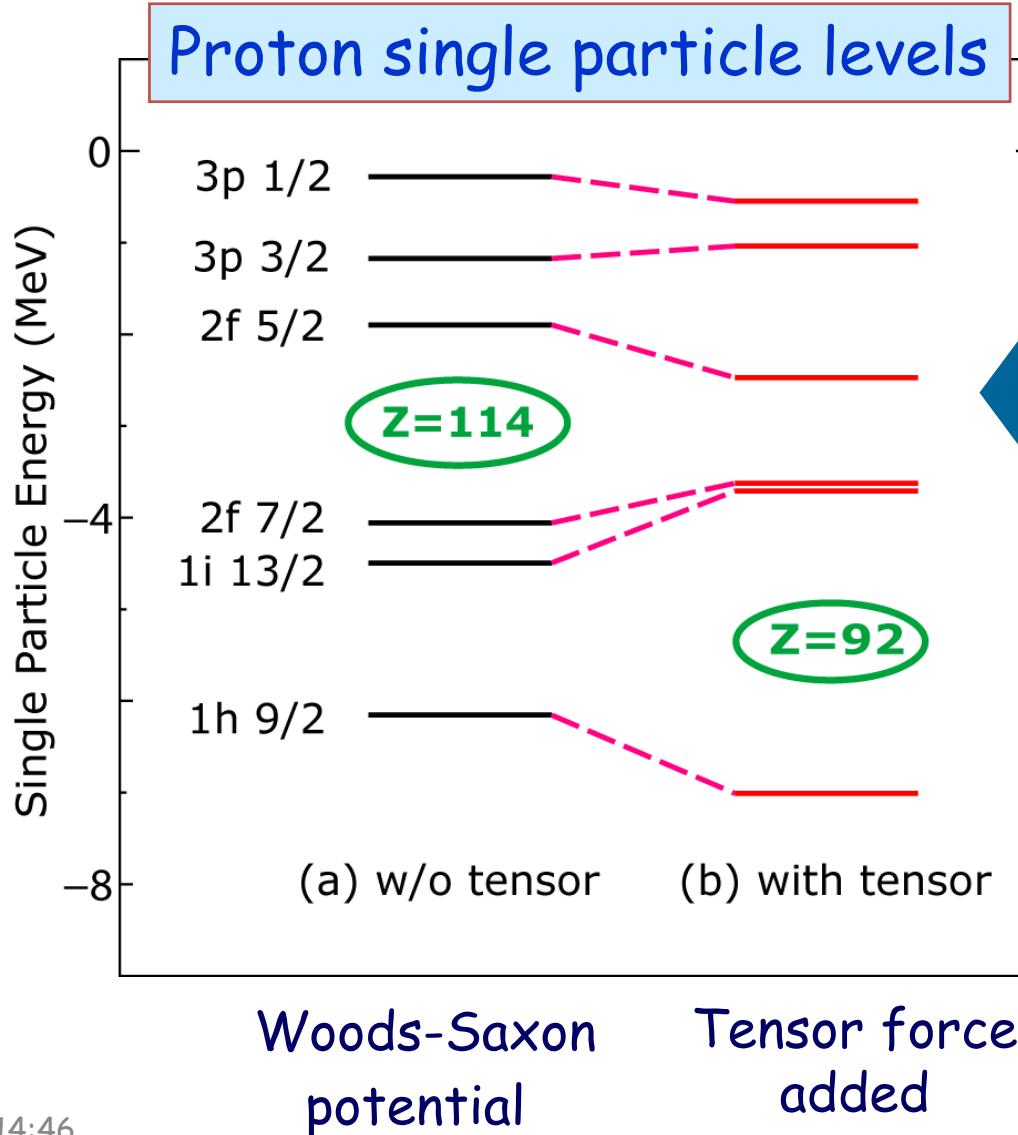
Strong oblate
Deformation ?

14:46

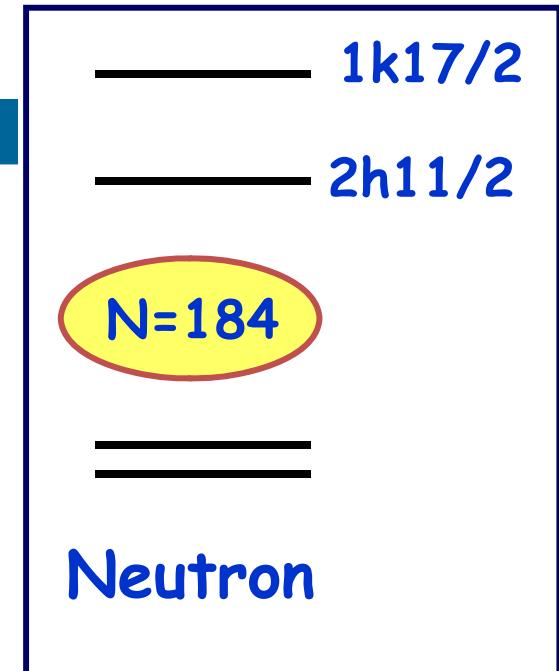


O, Suzuki and Utsuno,
Nucl. Phys. A805, 127c (2008)

Effect of tensor force on (spherical) superheavy magic numbers



Occupation of neutron
 $1k17/2$ and
 $2h11/2$



Otsuka, Suzuki and Utsuno,
Nucl. Phys. A805, 127c (2008)

戦いはまだ終わっていない . . .

PRL 100, 062501 (2008)

PHYSICAL REVIEW LETTERS

week ending
15 FEBRUARY 2008



Shell Model Description of the ^{14}C Dating β Decay with Brown-Rho-Scaled NN Interactions

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(Received 21 September 2007; published 15 February 2008)

We present shell model calculations for the beta decay of ^{14}C to the ^{14}N ground state, treating the states of the $A = 14$ multiplet as two $0p$ holes in an ^{16}O core. We employ low-momentum nucleon-nucleon (NN) interactions derived from the realistic Bonn-B potential and find that the Gamow-Teller (GT) matrix element is too large to describe the known lifetime. By using a modified version of this potential that incorporates the effects of Brown-Rho scaling medium modifications, we find that the GT matrix element vanishes for a nuclear density around 85% that of nuclear matter. We find that the splitting between the $(J^\pi, T) = (1^+, 0)$ and $(J^\pi, T) = (0^+, 1)$ states in ^{14}N is improved using the medium-modified Bonn-B potential and that the transition strengths from excited states of ^{14}C to the ^{14}N ground state are compatible with recent experiments.

In-medium ではテンソル力(による相関)は弱い
と言っている。(実は configuration space が狭いだけ?)

有効核力と Extended Weinberg Ansatz

成功している有効相互作用の解剖学

pf殻に対する GXPF1

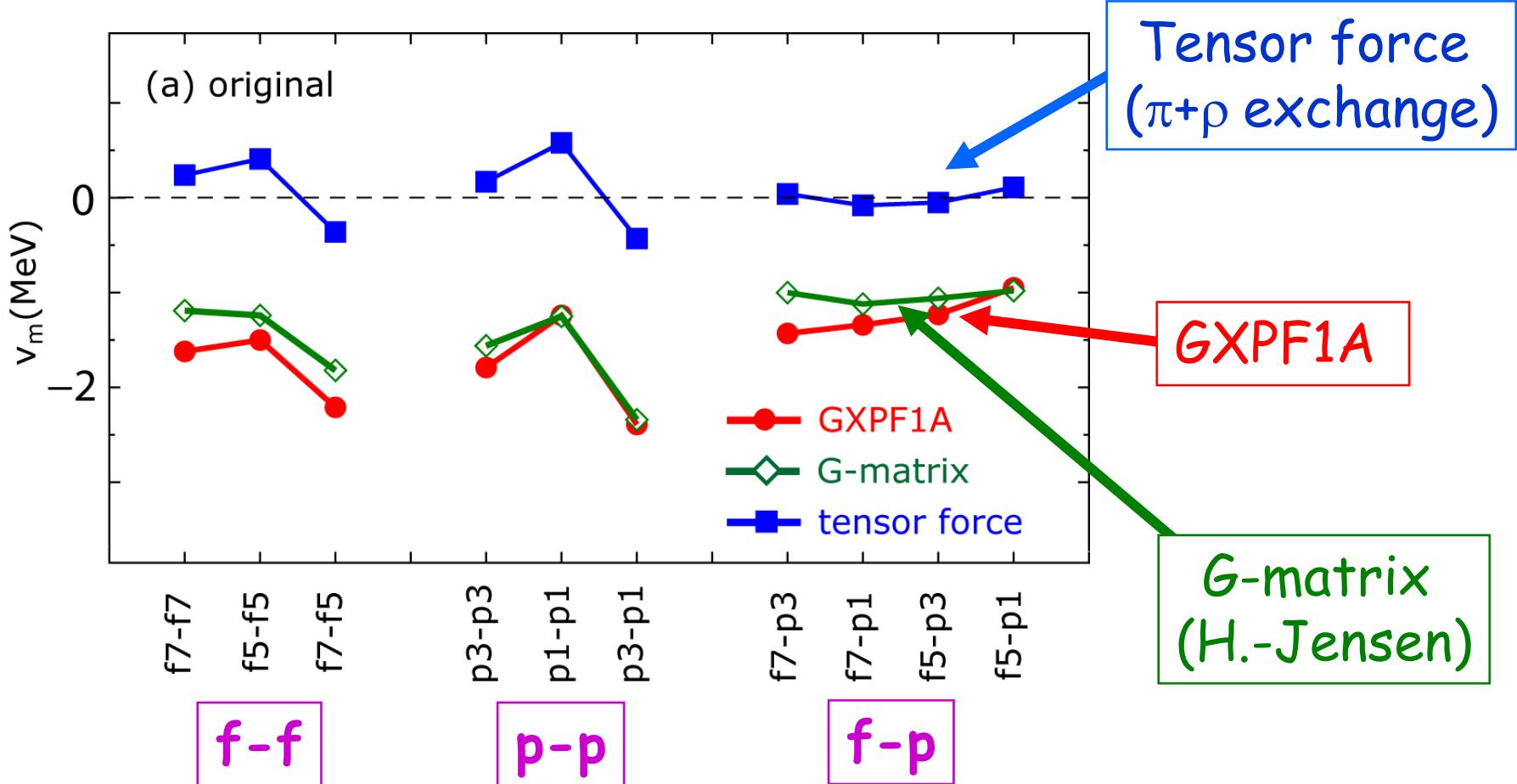
本間、他 2002年～ 最新の公開バージョン GXPF1A

G-matrix (H.-Jensen) から得られた微視的相互作用
(199個の独立な量)

その内約70個に現象論的な補正

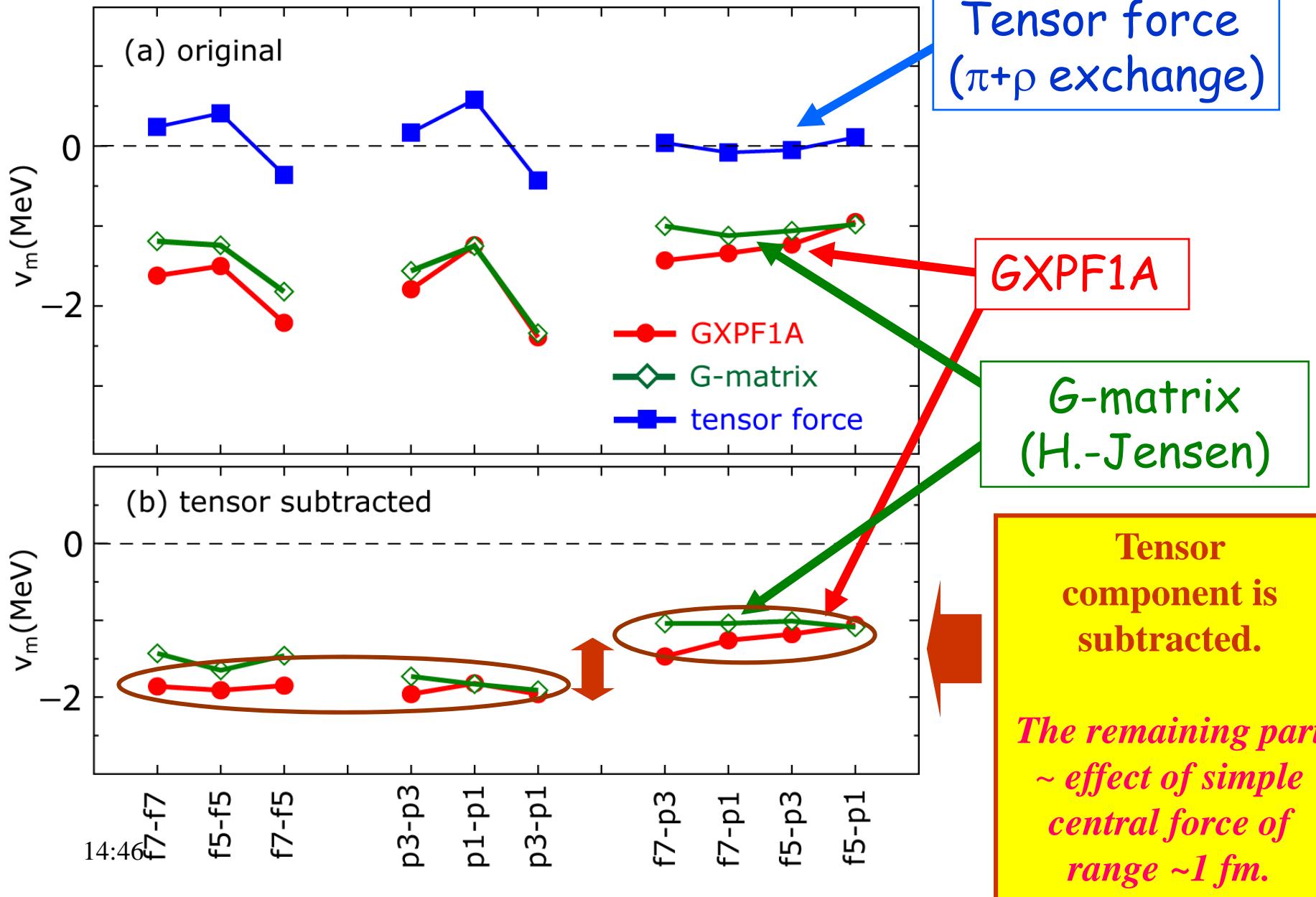
699個の実験データを使用 + その後のマイナーな改訂

T=0 monopole interactions in the pf shell



"Local pattern" ← tensor force

T=0 monopole interactions in the pf shell



Monopole interaction

1 fm 程度のレンジ(ガウス関数)を持つ中心力

結合力の源、概ね定数

繰り込みの効果

動径波動関数による → ノード数が違うと弱くなる

$\pi + \rho$ 中間子交換から来る bare のテンソル力

大きさは中心力より小さい

斥力になることがあるので、殻構造に大きな影響

2体 LS 力 (今日は省略) 限られた場合に効果

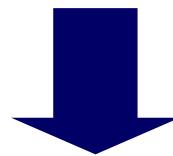
The central force is modeled by a Gaussian function

$$V = V_0 \exp(-r/\mu)^2 \quad (S,T \text{ dependences})$$

with $V_0 = -166 \text{ MeV}$, $\mu = 1.0 \text{ fm}$,

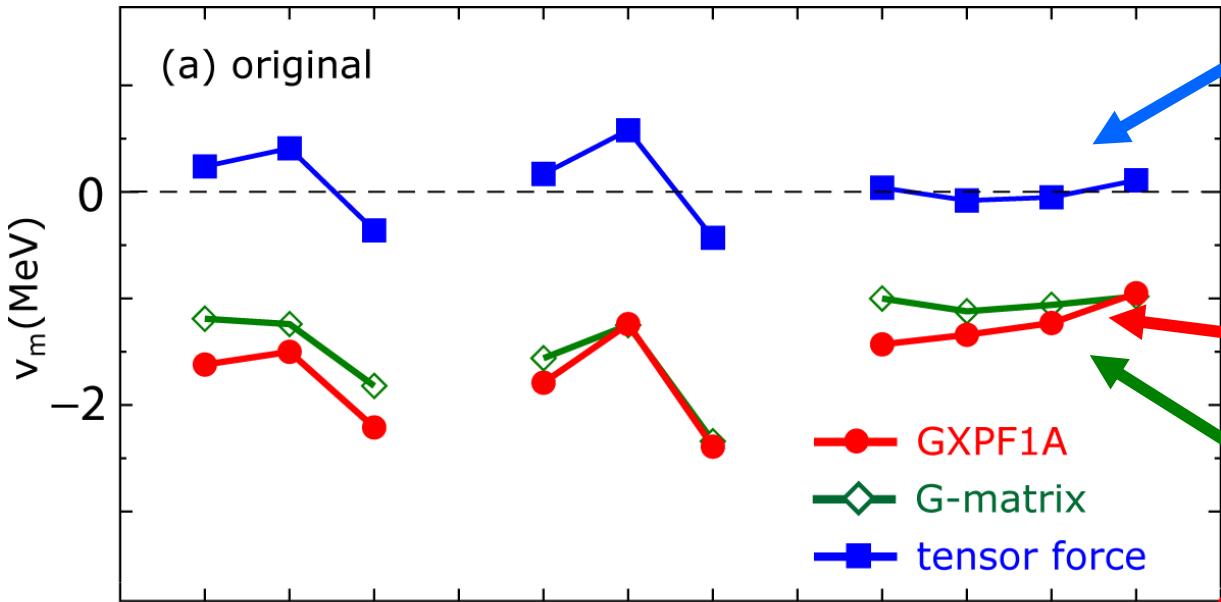
(S,T) factor (0,0) (1,0) (0,1) (1,1)

relative strength 1 1 0.6 -0.8

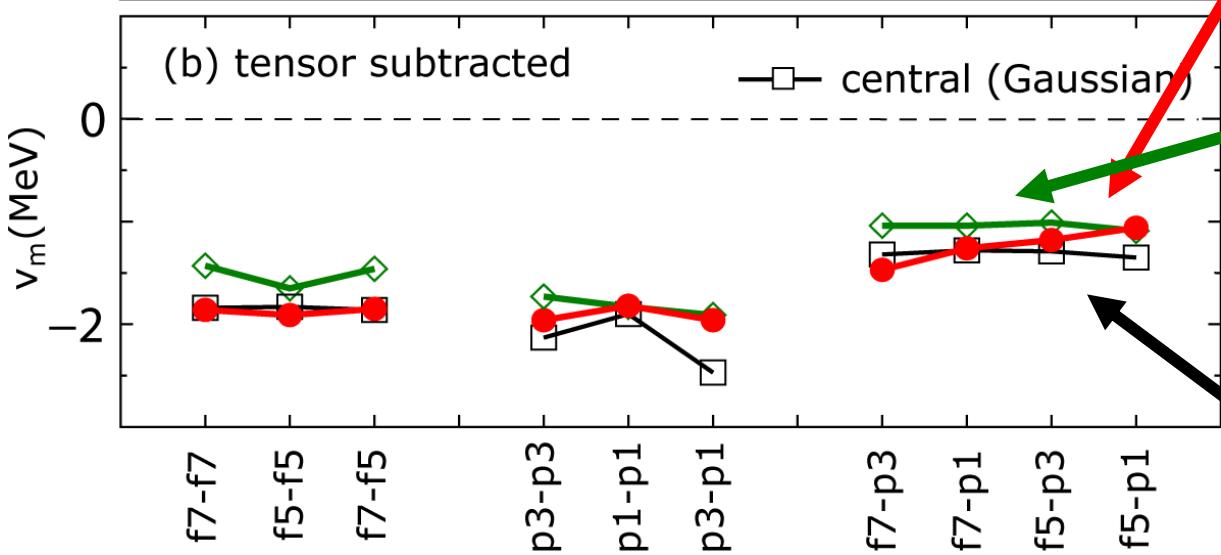


Can we explain the difference between f-f/p-p and f-p ?

T=0 monopole interactions in the pf shell



Tensor force
($\pi+p$ exchange)



central (Gaussian)

G-matrix
(H.-Jensen)

Central (Gaussian)
- Reflecting
radial overlap -

14:46

f-f

p-p

f-p

This is not a very lonely idea → Chiral Perturbation of QCD

S. Weinberg,
PLB 251, 288 (1990)

Short range central forces
have complicated origins and
should be adjusted.

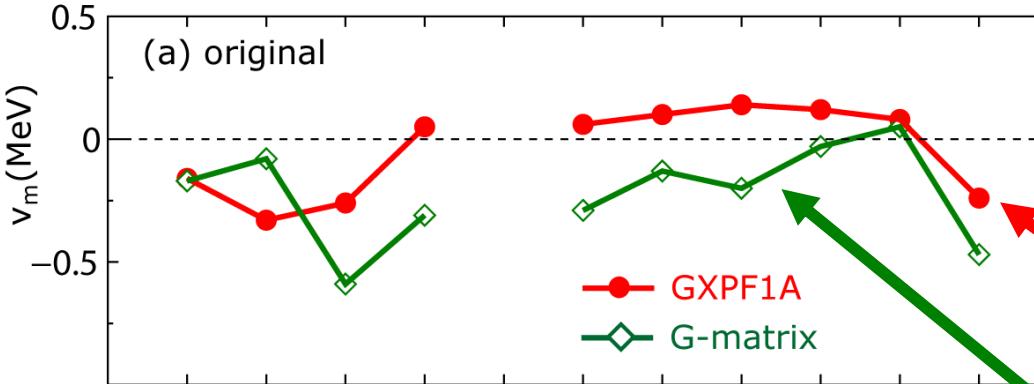
tive potential gives a local coordinate-space two-nucleon potential:

$$V_{2\text{-nucleon}} = 2(C_S + C_T \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) \delta^3(\mathbf{x}_1 - \mathbf{x}_2) \\ - \left(\frac{2g_A}{F_\pi} \right)^2 (\mathbf{t}_1 \cdot \mathbf{t}_2) (\boldsymbol{\sigma}_1 \cdot \nabla_1) (\boldsymbol{\sigma}_2 \cdot \nabla_2) Y(|\mathbf{x}_1 - \mathbf{x}_2|) \\ - (1' \leftrightarrow 2'),$$

Tensor force is explicit

where $Y(r) \equiv \exp(-m_\pi r)/4\pi r$ is the usual Yukawa potential. [Throughout it should be understood that these are local potentials, containing a delta function factor like $\delta^3(\mathbf{x}'_1 - \mathbf{x}_1)$ for each nucleon.]

T=1 monopole interaction

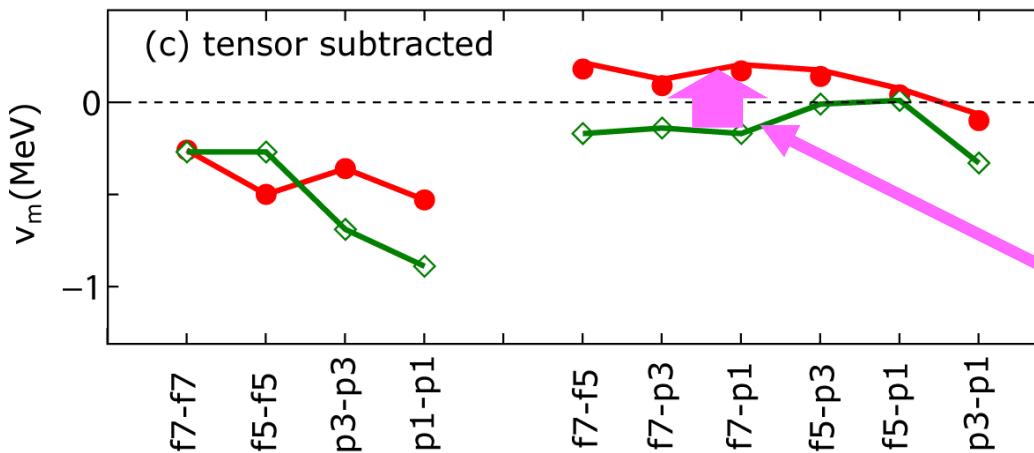
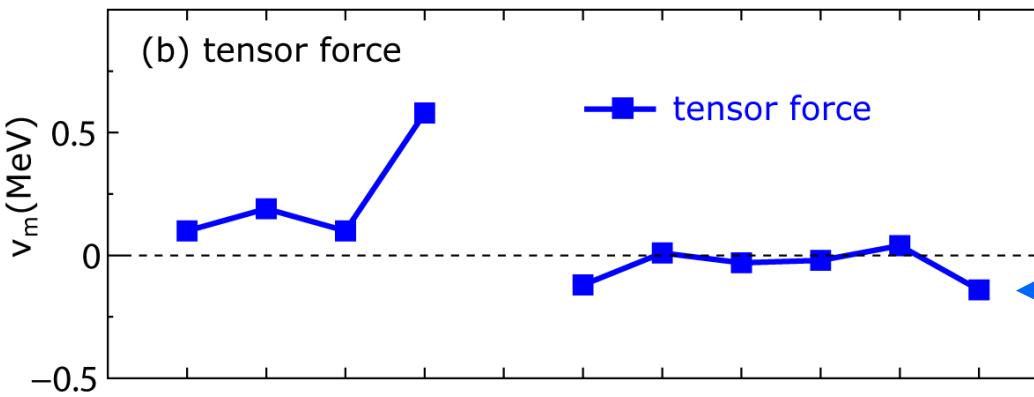


$T=1$ monopole interactions in the pf shell

GXPF1A

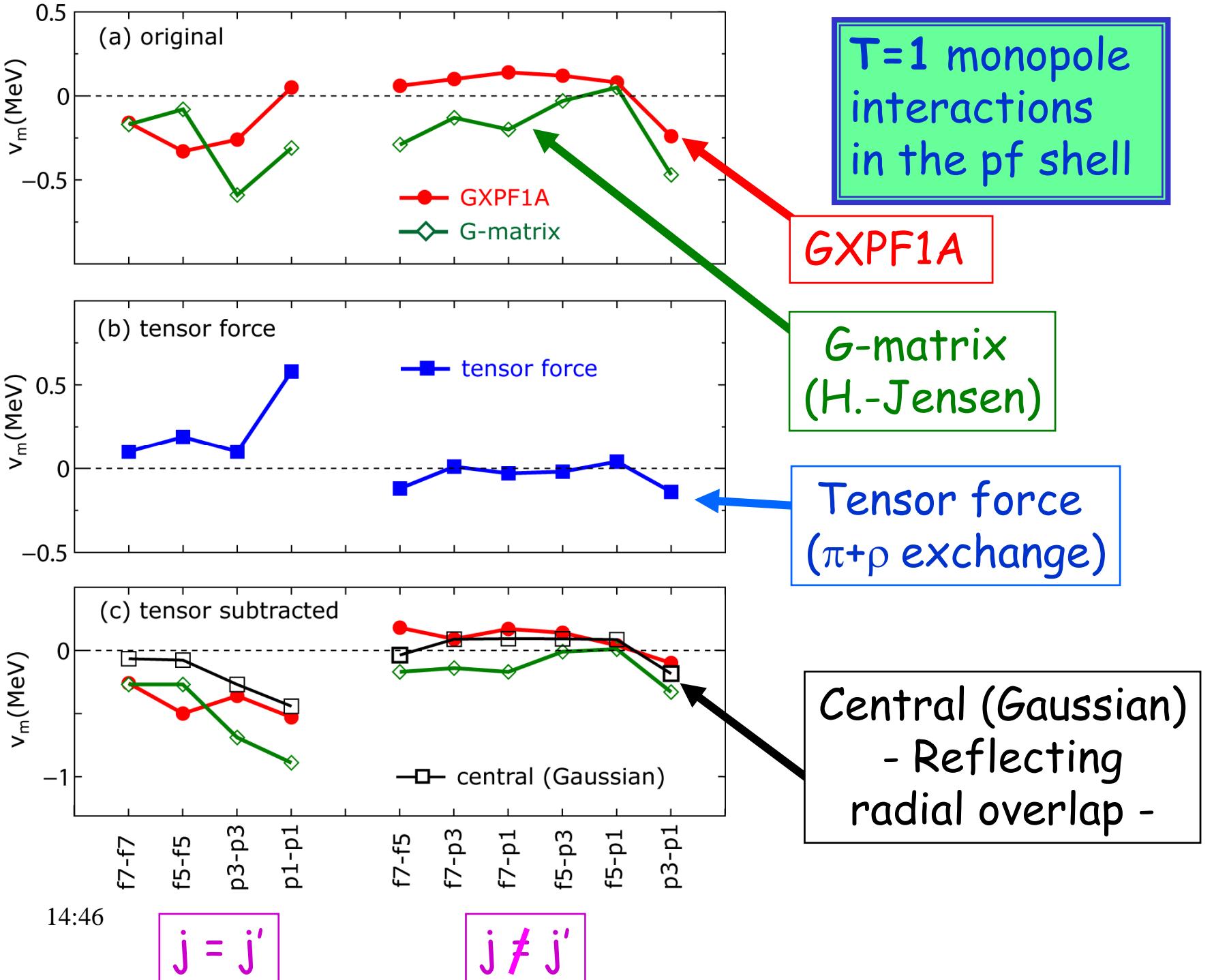
**G-matrix
(H.-Jensen)**

**Tensor force
($\pi+\rho$ exchange)**



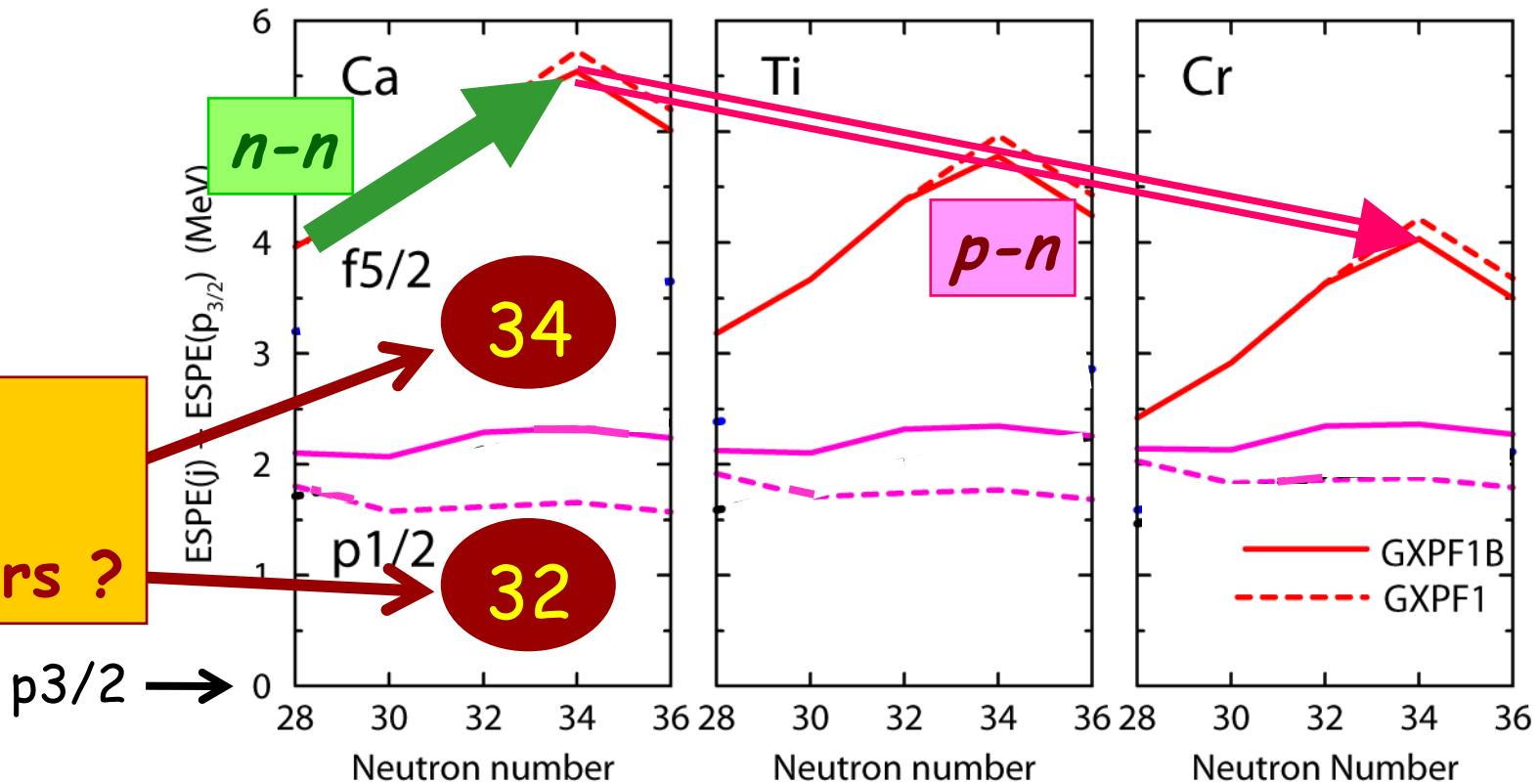
Basic scale
 $\sim 1/10$ of $T=0$

**Repulsive
corrections
to G-matrix**



(effective) single-particle energies relative to $p_{3/2}$

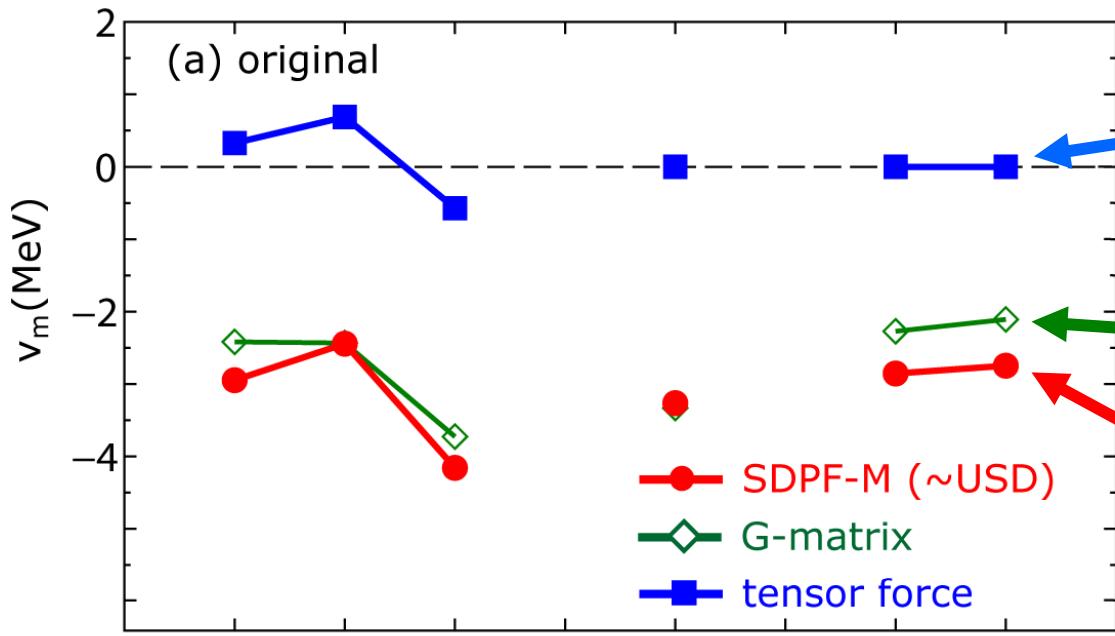
new
magic
numbers ?



Reduction of $f_{5/2}$ - $p_{3/2}$ from Ca to Cr p - n force:
 ~ 1.5 MeV = 1.0 MeV (tensor) + 0.5 MeV (central)

Rising of $f_{5/2}$ from ^{48}Ca to ^{54}Ca nn force:
 $\sim p_{3/2}$ - $p_{3/2}$ attraction + $p_{3/2}$ - $f_{5/2}$ repulsion

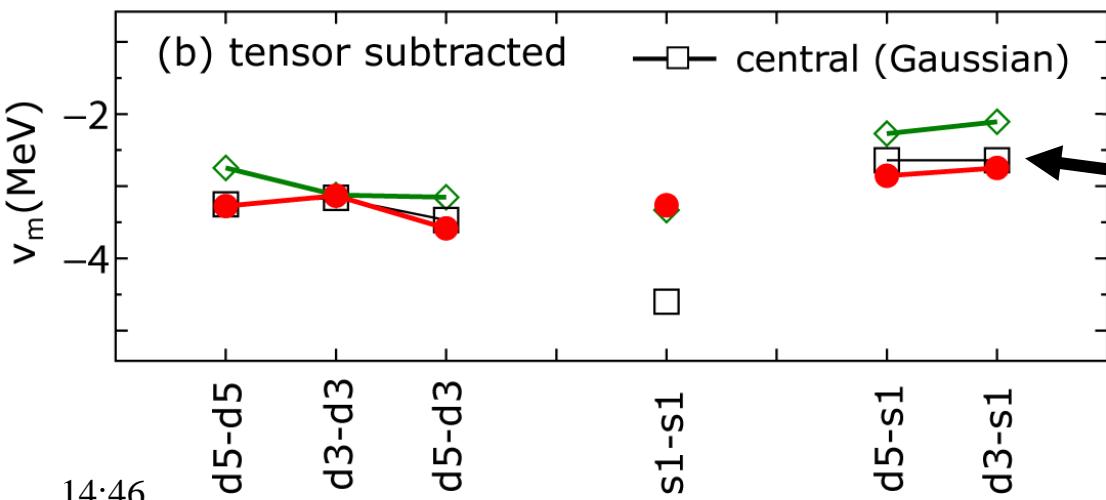
T=0 monopole interactions in the *sd* shell



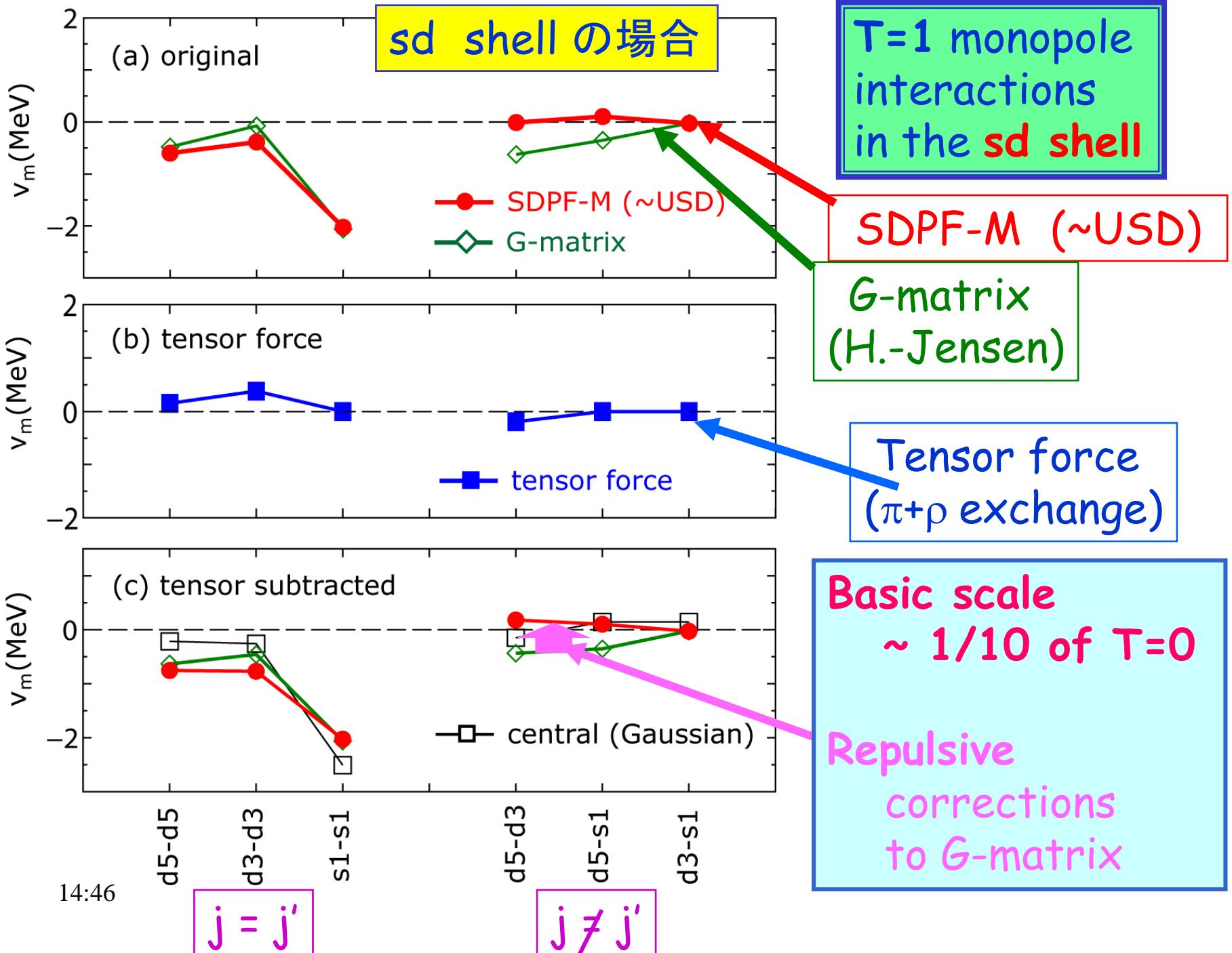
Tensor force
($\pi+\rho$ exchange)

G-matrix
(H.-Jensen)

SDPF-M
(~USD)

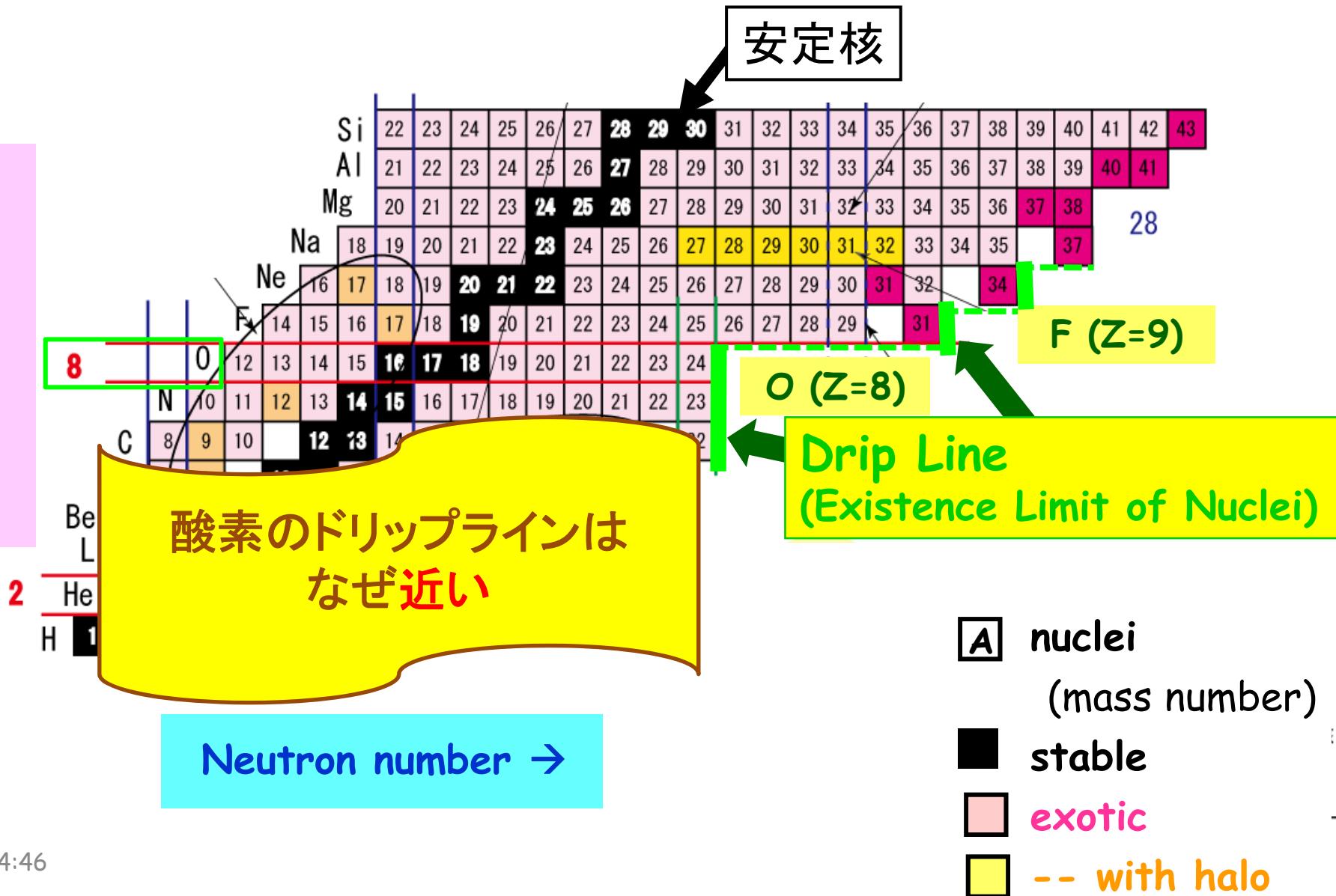


Central (Gaussian)
- Reflecting
radial overlap -



T=1 の斥力補正とドリップライン

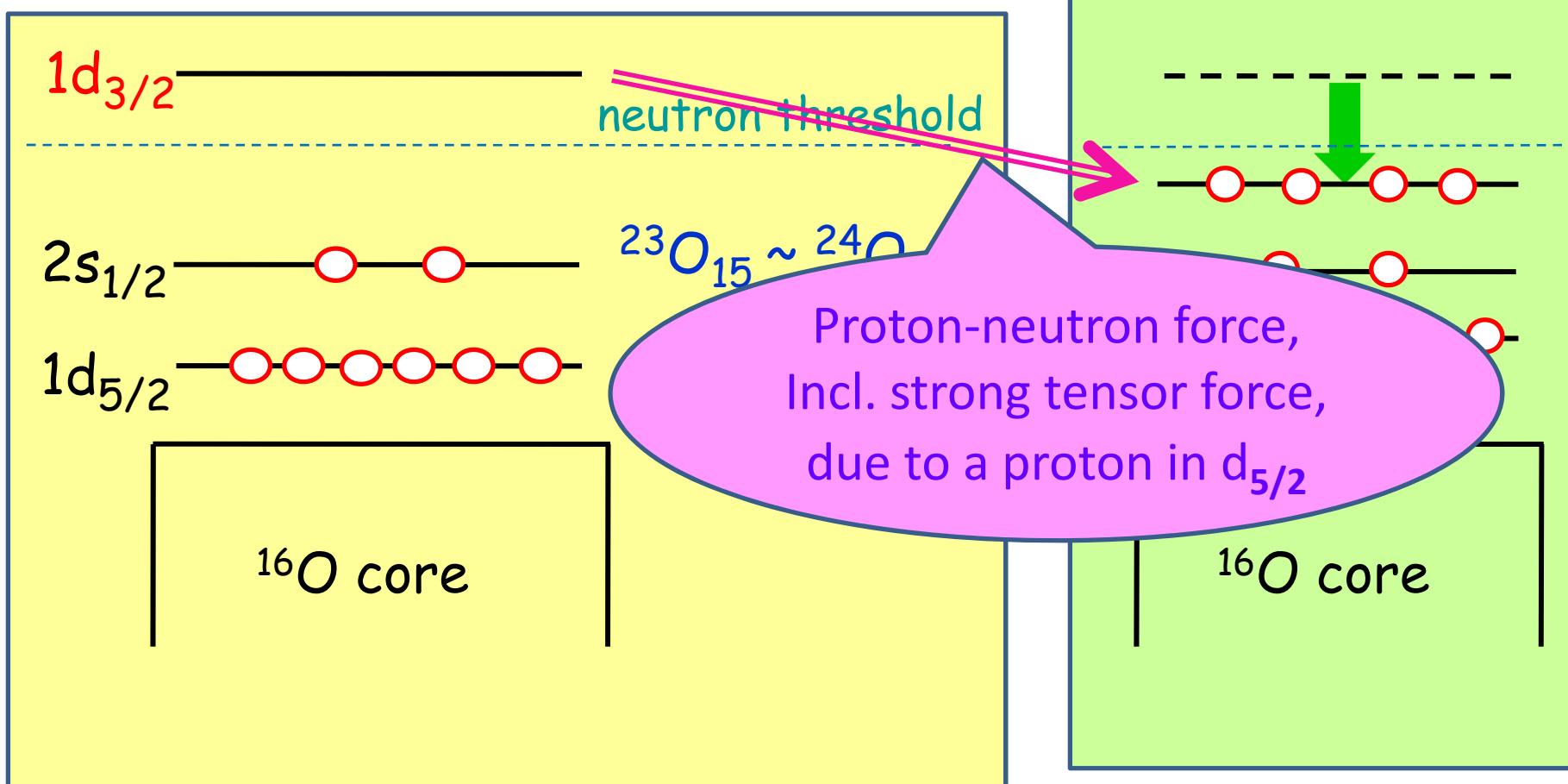
Proton number →



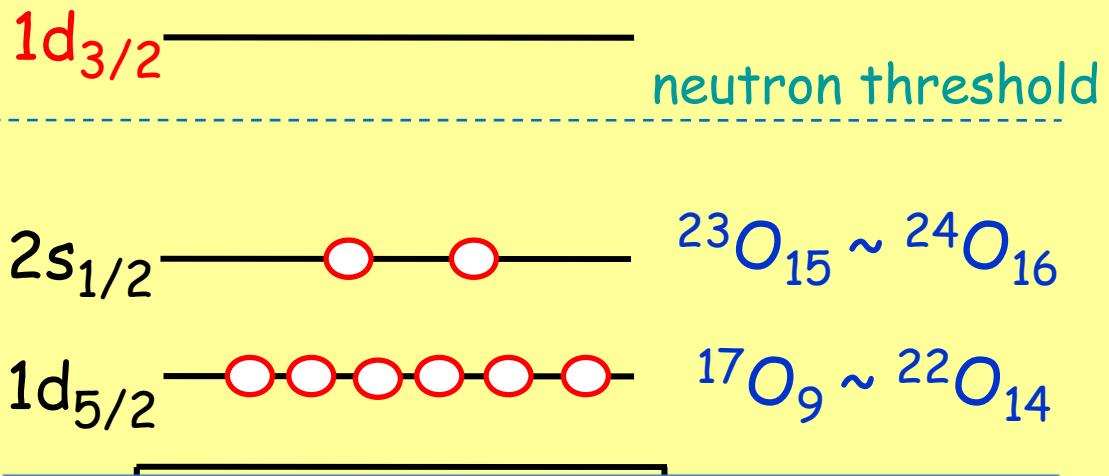
This is because the neutron $d_{3/2}$ orbit is high for Oxygen.

Neutron orbits in Oxygen isotopes

Neutron orbits in Fluorine isotopes

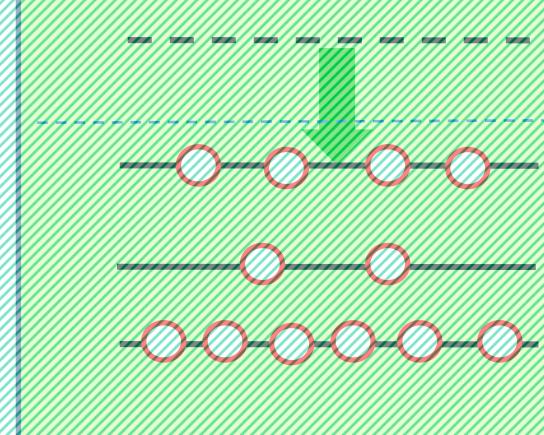


Neutron orbits in Oxygen isotopes



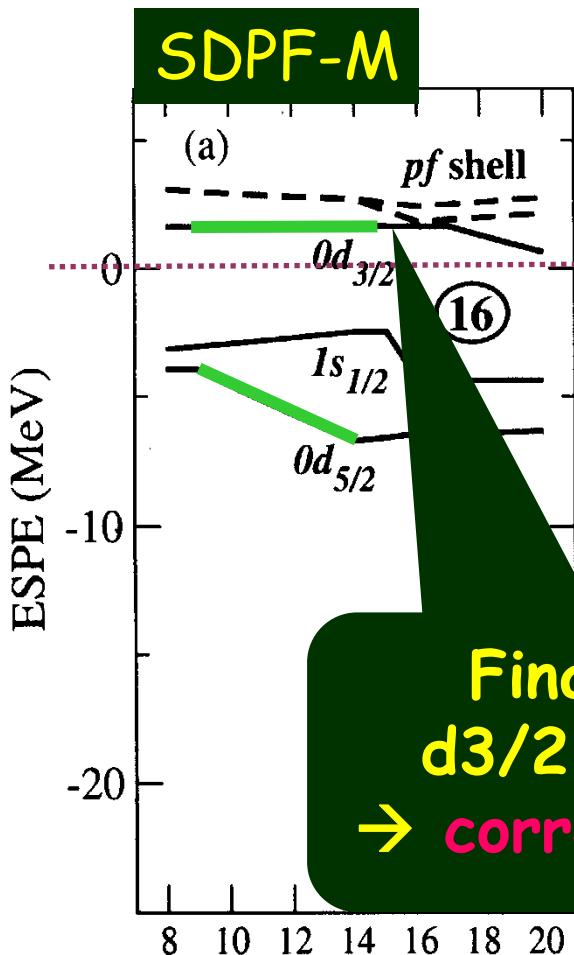
Why do those neutrons **NOT**
pull down $d_{3/2}$?

Neutron orbits in Fluorine isotopes

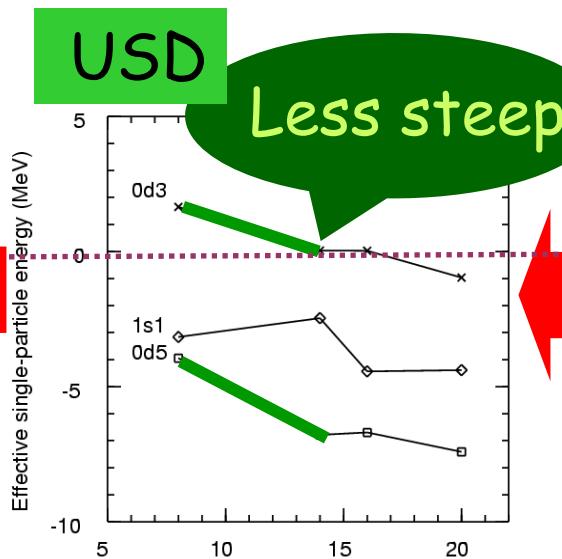


Effective Single-Particle Energy for Oxygen isotopes

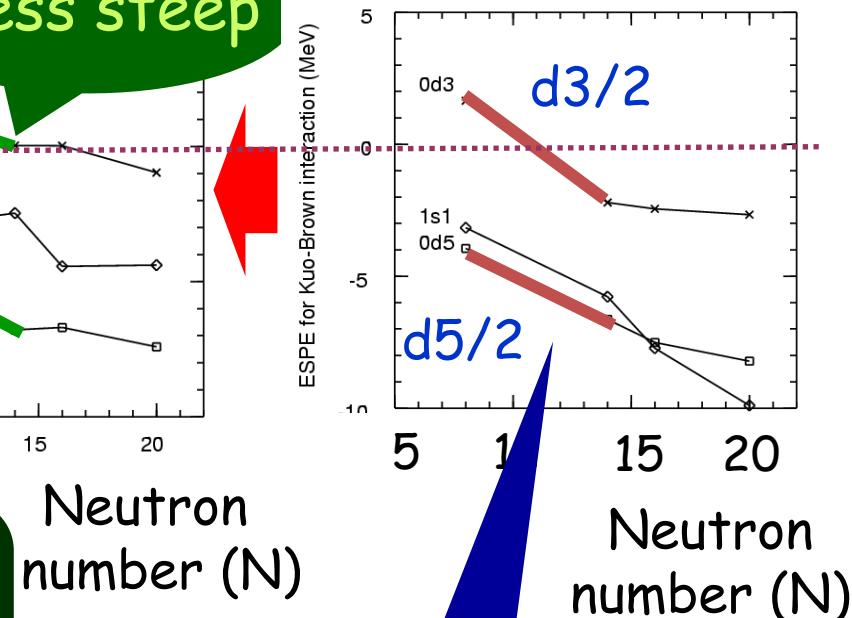
Final correction



Empirical correction



Kuo-Brown
G-matrix
+ core-pol.



Finally flat,
d_{3/2} kept high
→ correct drip line

narrowing

斥力補正の起源は、これだけ2体力を考えても
見つからないという状況証拠からも

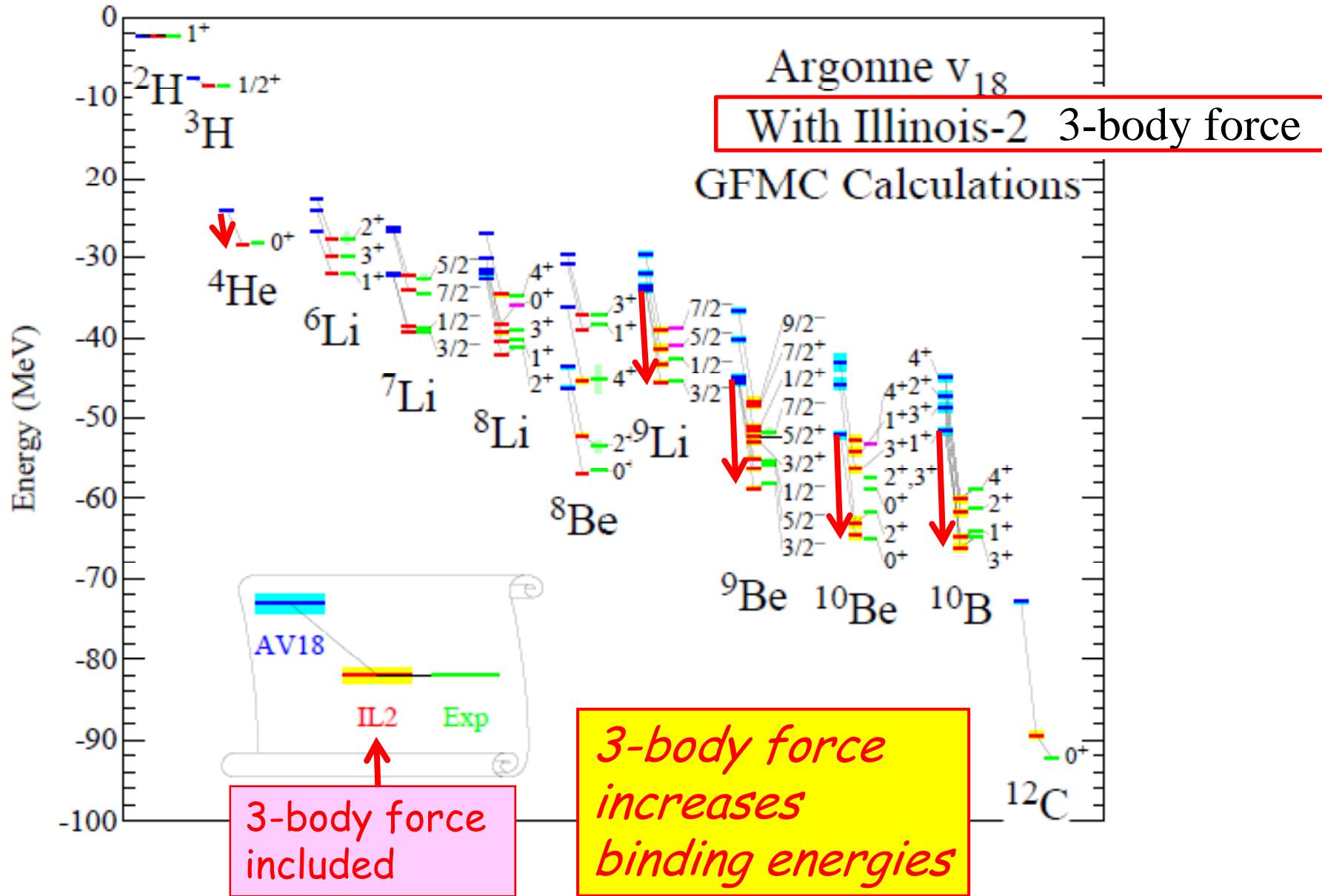
→ 3体力に起源？

3NF → attractive effects

systematics in results of GFMC, NCSM

CC (Hagen et al., Phys. Rev. C76, 034302 (2007))

GFMC (Green Function Monte Carlo) by Argonne group



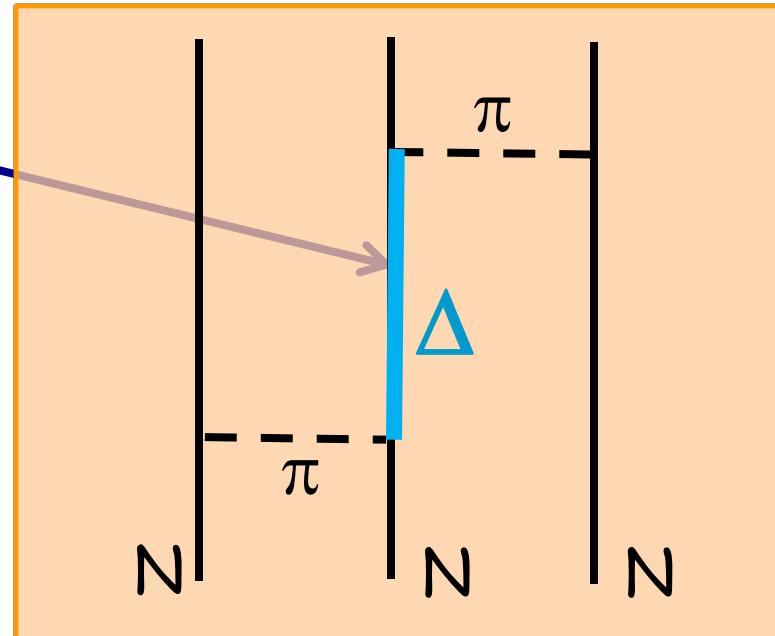
The key : Fujita-Miyazawa 3N mechanism (Δ -hole excitation)

Progress of Theoretical Physics, Vol. 17, No. 3, March 1957

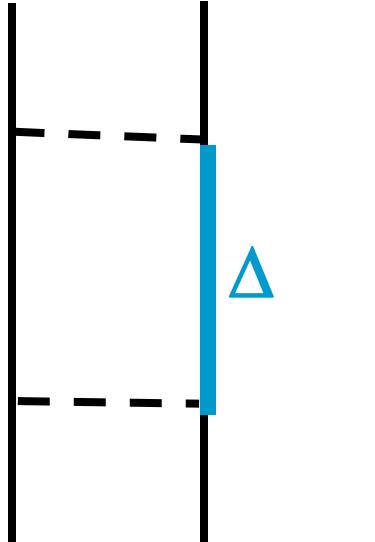
Pion Theory of Three-Body Forces

Jun-ichi FUJITA and Hironari MIYAZAWA

Δ particle
 $m=1232 \text{ MeV}$
 $S=3/2, I=3/2$



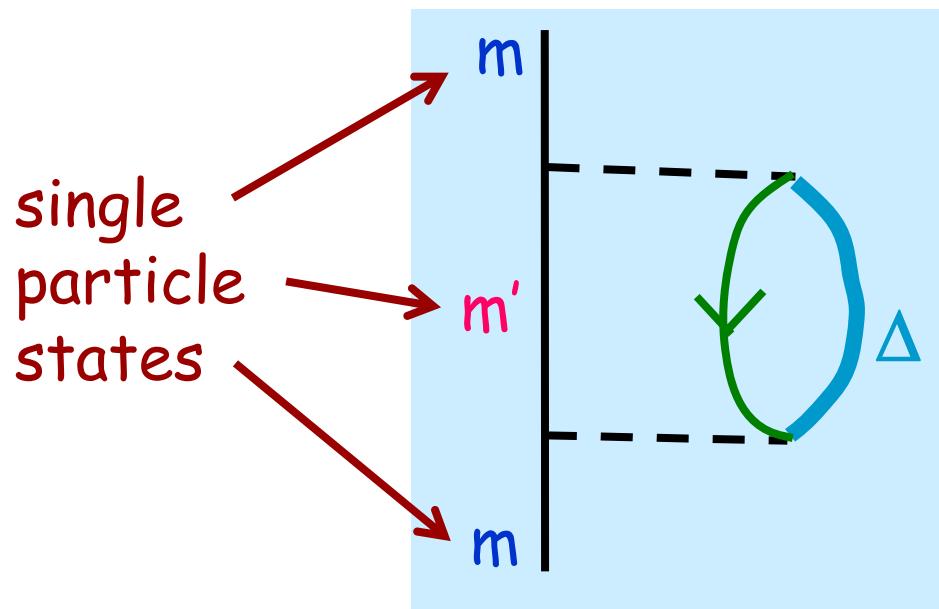
Renormalization of NN interaction due to Δ excitation in the intermediate state



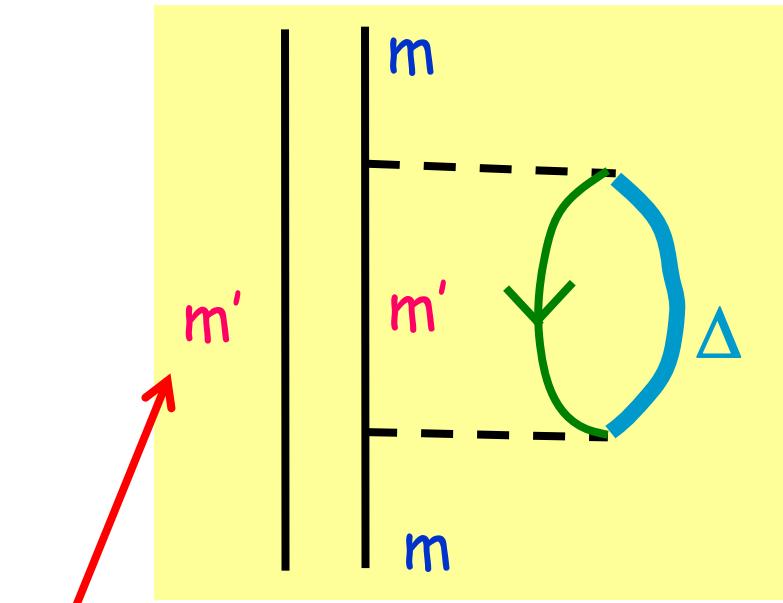
Modification to
bare NN interaction
(for NN scattering)

$T=1$
attraction
between NN
effectively

Pauli blocking effect on the renormalization of single-particle energy



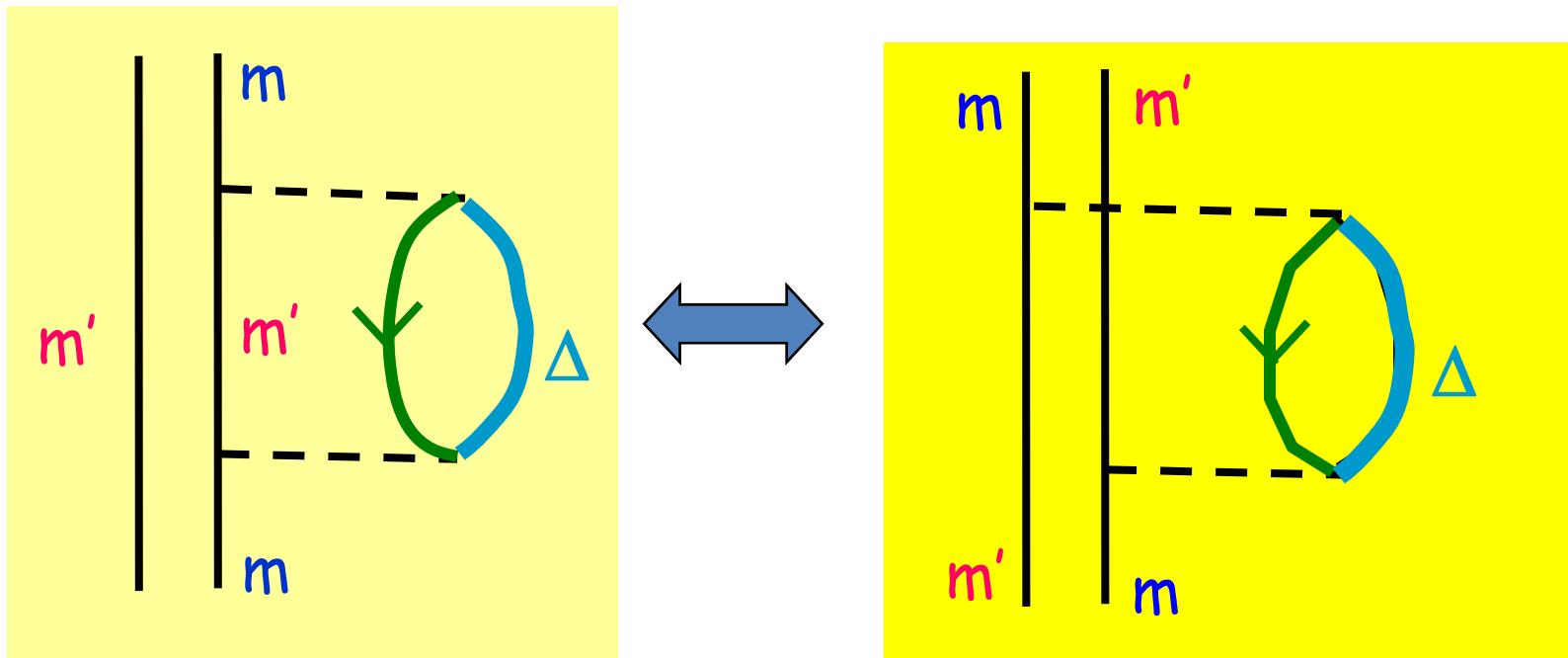
Renormalization of
single particle energy
due to
 Δ -hole excitation
→ more binding (attractive)



Another valence
particle in state m'

Pauli Forbidden
→ The effect is
suppressed

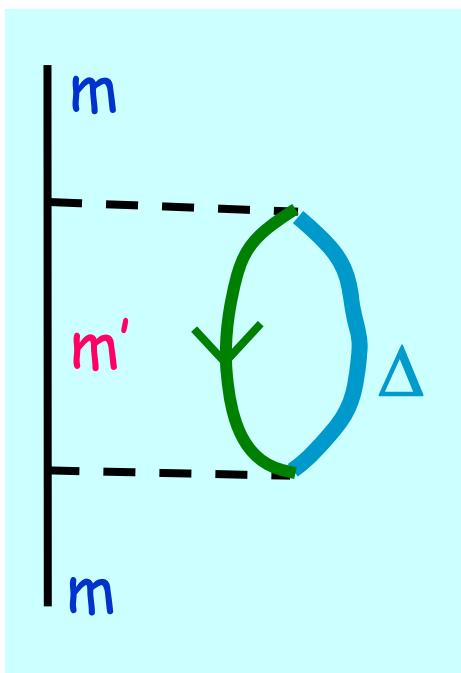
Inclusion of Pauli blocking



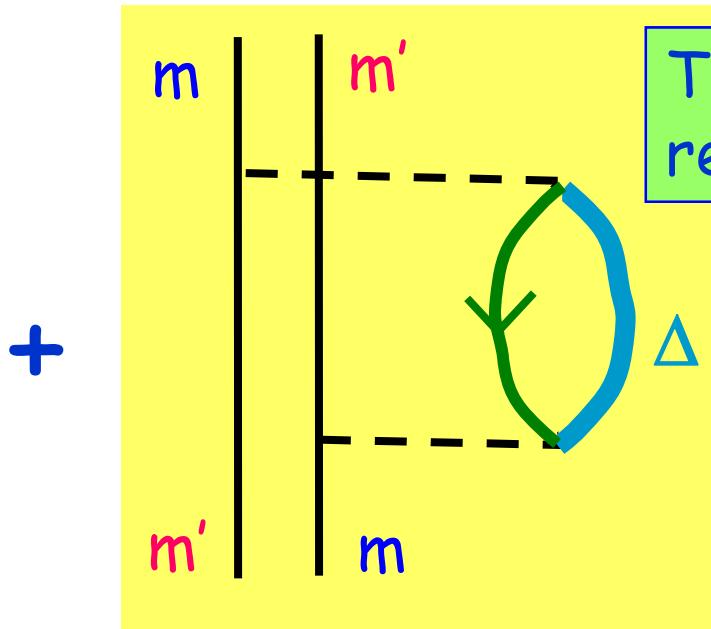
Pauli forbidden
(from previous page)

This Pauli effect is
included automatically
by the exchange term.

Realization in terms of 3-body interaction



Renormalization
of single particle
energy

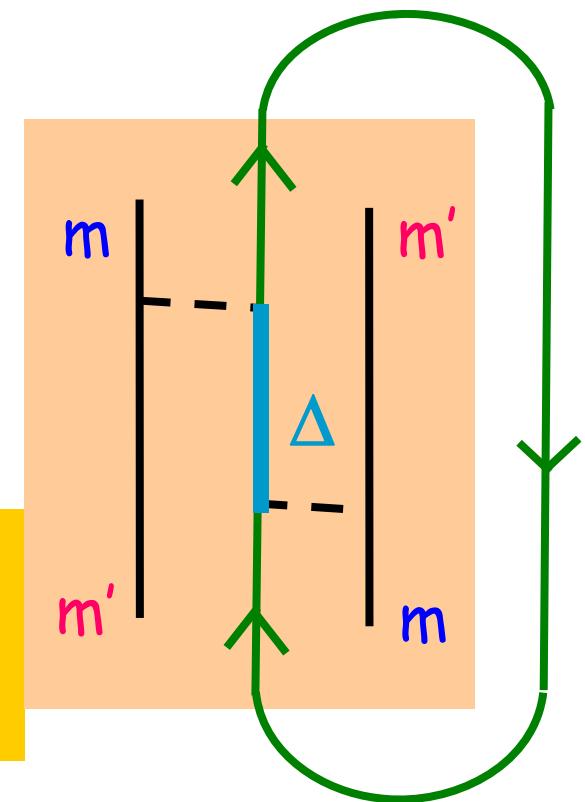


T=1 effective monopole
repulsive interaction

↑ *T=0 does not survive
3rd order (Akaishi et al.)*

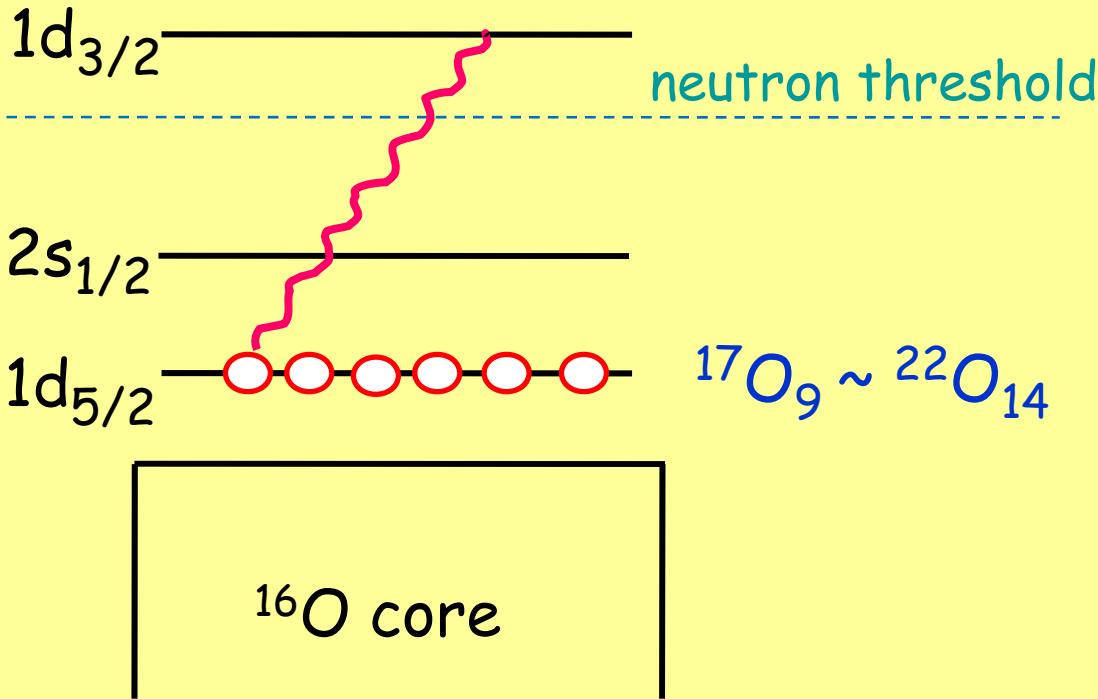
Pauli blocking
same

Monopole part of
Fujita-Miyazawa
3-body force



Back to the question of high-lying $d_{3/2}$

Neutron orbits in Oxygen isotopes



Central :
attractive
(generally)

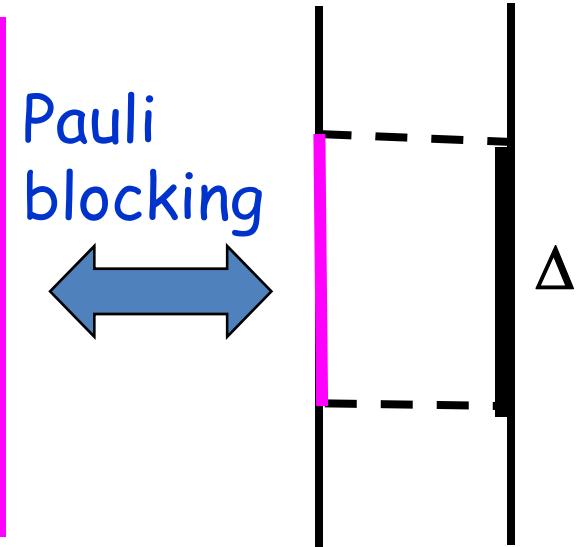
Tensor :
attractive
- 0.9 MeV
(next page)

Δ -hole induced
repulsion
Next page

Other diagram included

Particle in the inert core

Pauli blocking



T=1
interaction
between
valence
particles

Related effect
was discussed by
Frisch, **Kaiser** and
Weise for
neutron matter

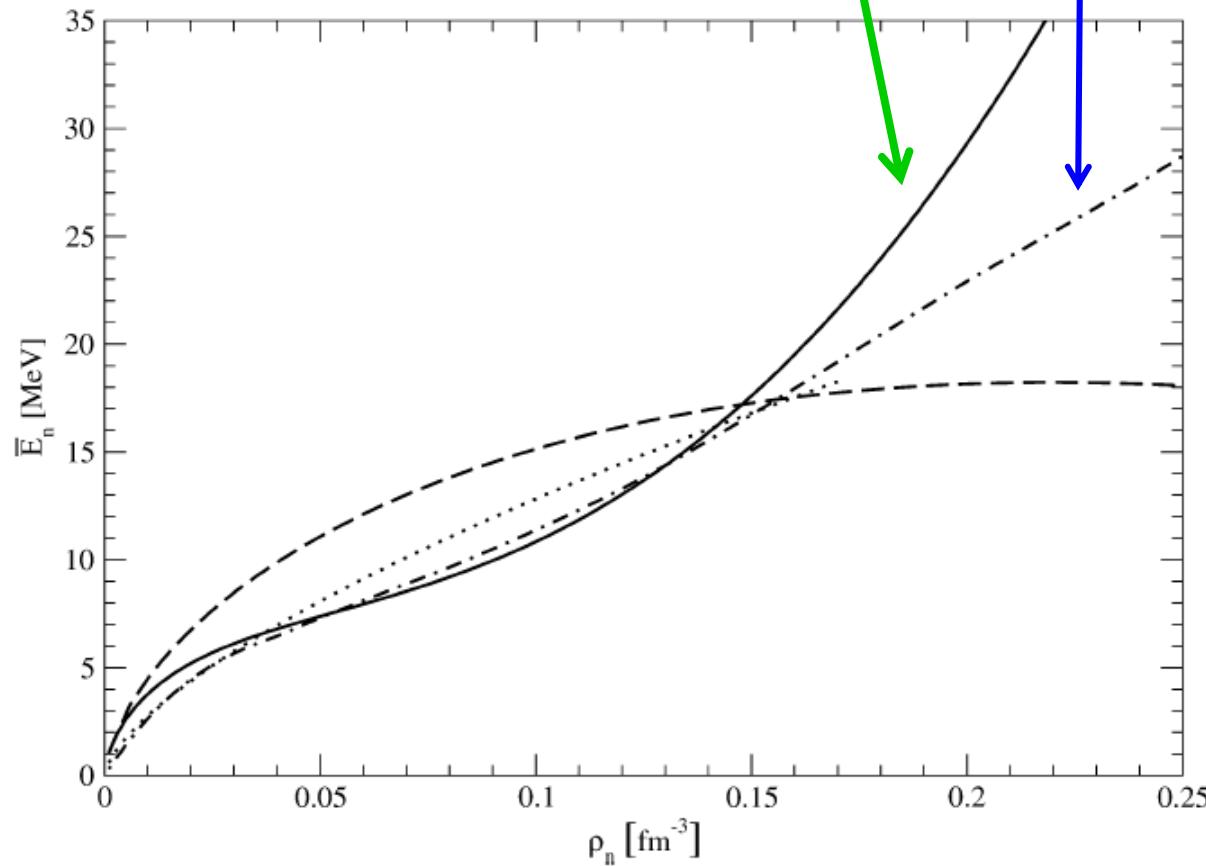
See also **Nishizaki**,
Takatsuka and **Hiura**
PTP 92, 93 (1994)

Δ -hole excitation may be crucial to neutron matter property

Chiral Perturbation incl. Δ : Frisch, Kaiser and Weise

A. Akmal, V.R. Pandharipande, D.G. Ravenhall, Phys. Rev. C 58 (1998) 1804,

S. Fritsch et al. / Nuclear Physics A 750 (2005) 259–293



Repulsive effective monopole interaction assuming ^{16}O core

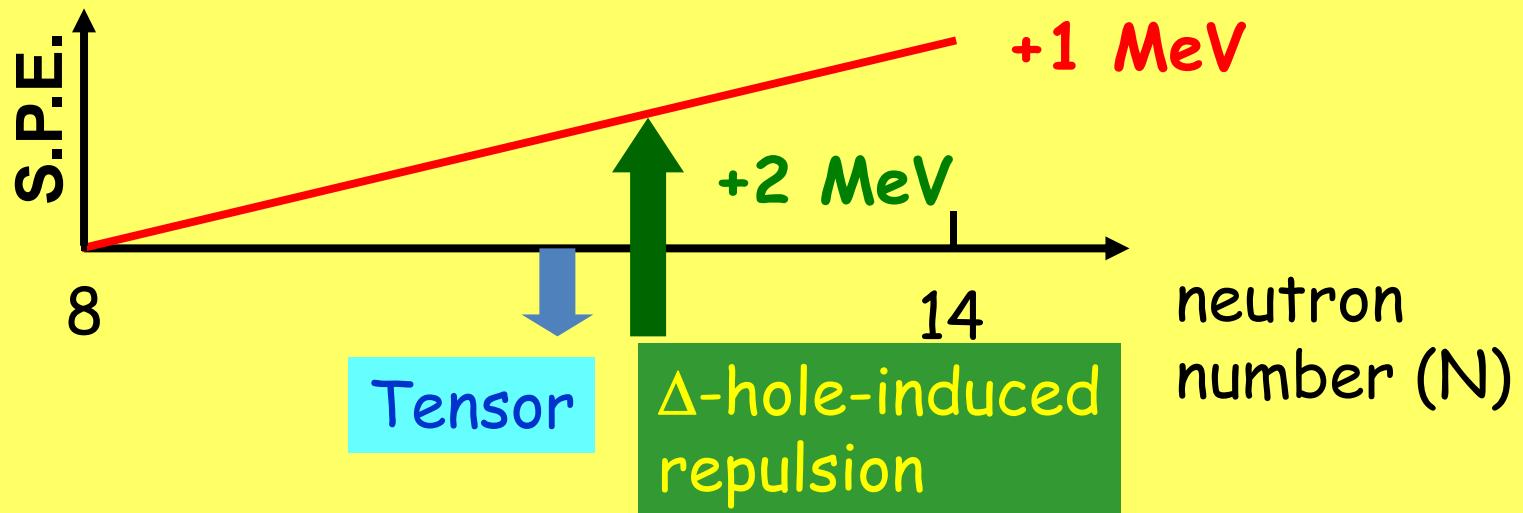
π exchange with radial cut-off at 0.7 fm , $\Delta E = 293 \text{ MeV}$

$$f_{\{\pi N \Delta\}} / f_{\{\pi NN\}} = \sqrt{9/2}$$

Monopole interaction

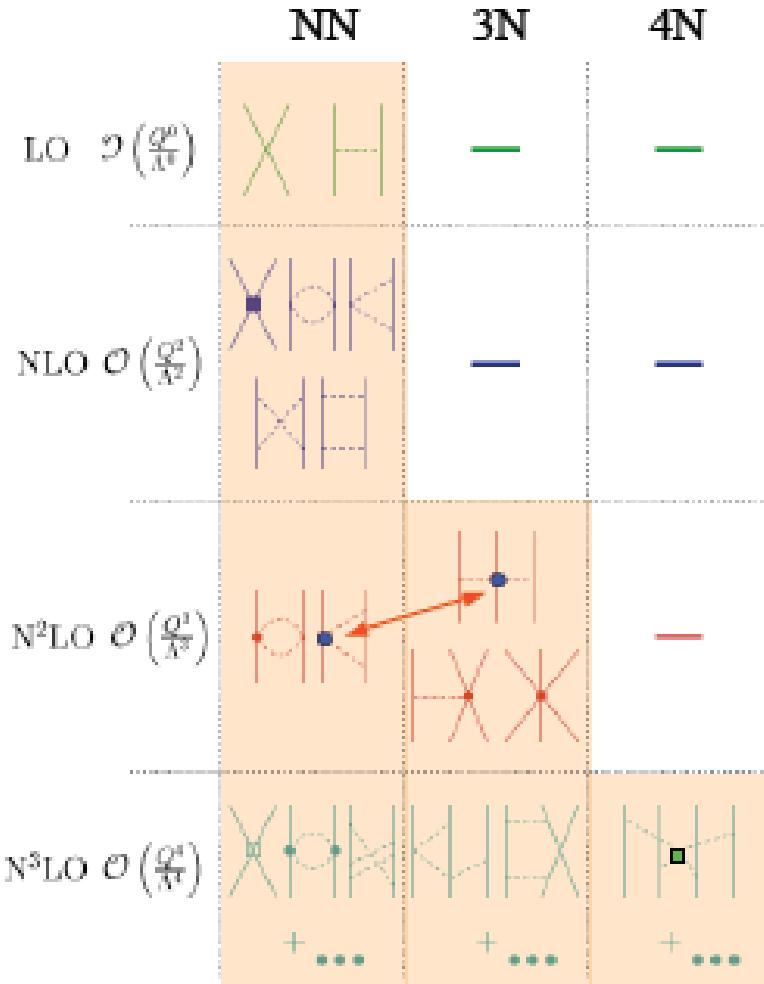
j j'
 $d_{5/2}$ $d_{3/2}$ pion tensor
 250 keV

$d_{3/2}$ single-particle energy relative to $N=8$



Chiral EFT for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale Λ_b



explains pheno hierarchy:

NN > 3N > 4N > ...

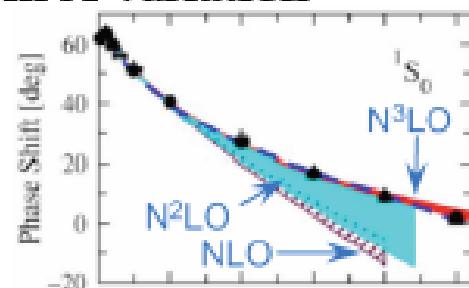
NN-3N, πN , $\pi\pi$, electro-weak,...
consistency

3N,4N: 2 new couplings to N^3LO

resolution/ Λ -dependent couplings

error estimates from truncation order,
lower bound from Λ variation

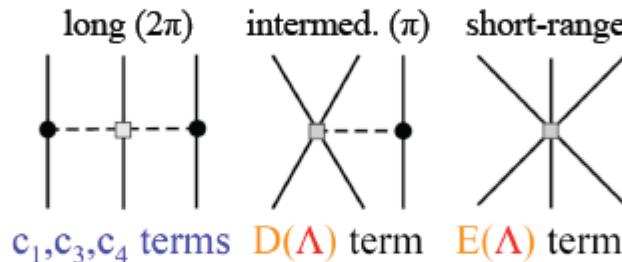
open problems:
power counting
 Δ -less vs. full



Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner, Nogga, Machleidt,...

Low-momentum 3N interactions

from leading N²LO chiral EFT $\sim (Q/\Lambda)^3$ van Kolck (1994), Epelbaum et al. (2002)



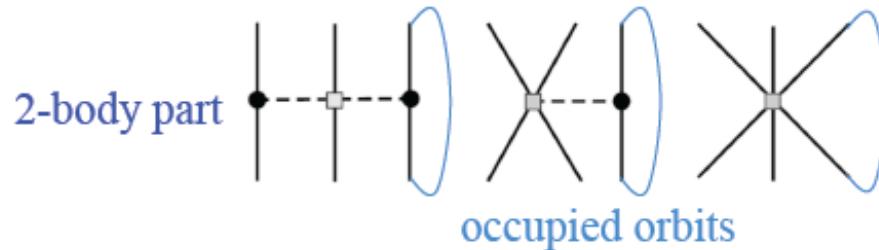
c_i from πN , consistent with NN

$$c_1 = -0.9^{+0.2}_{-0.5}, \quad c_3 = -4.7^{+1.2}_{-1.0}, \quad c_4 = 3.5^{+0.5}_{-0.2}$$

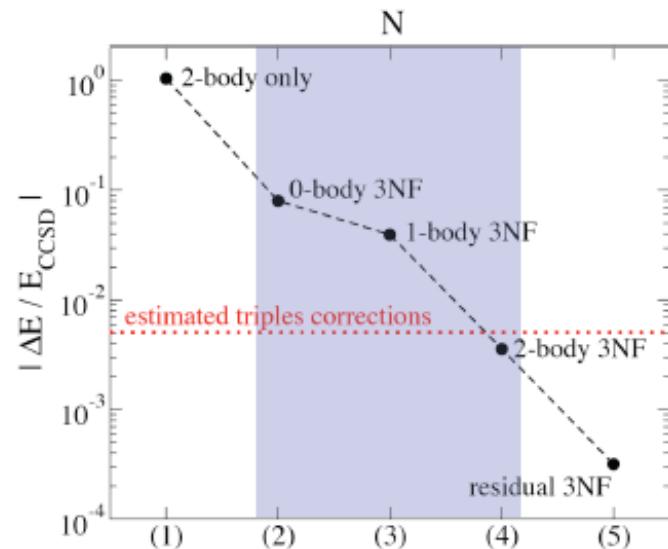
Meissner (2007)

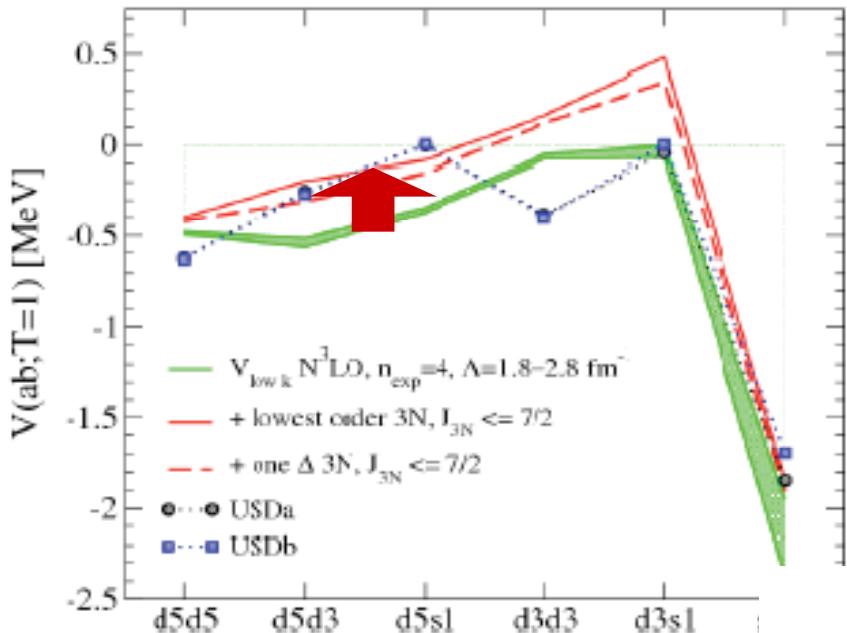
c_3, c_4 important for structure, large uncertainties at present

Results show that 0-, 1- and 2-body parts of 3N interaction dominate



residual 3N interaction can be neglected
very promising and practical

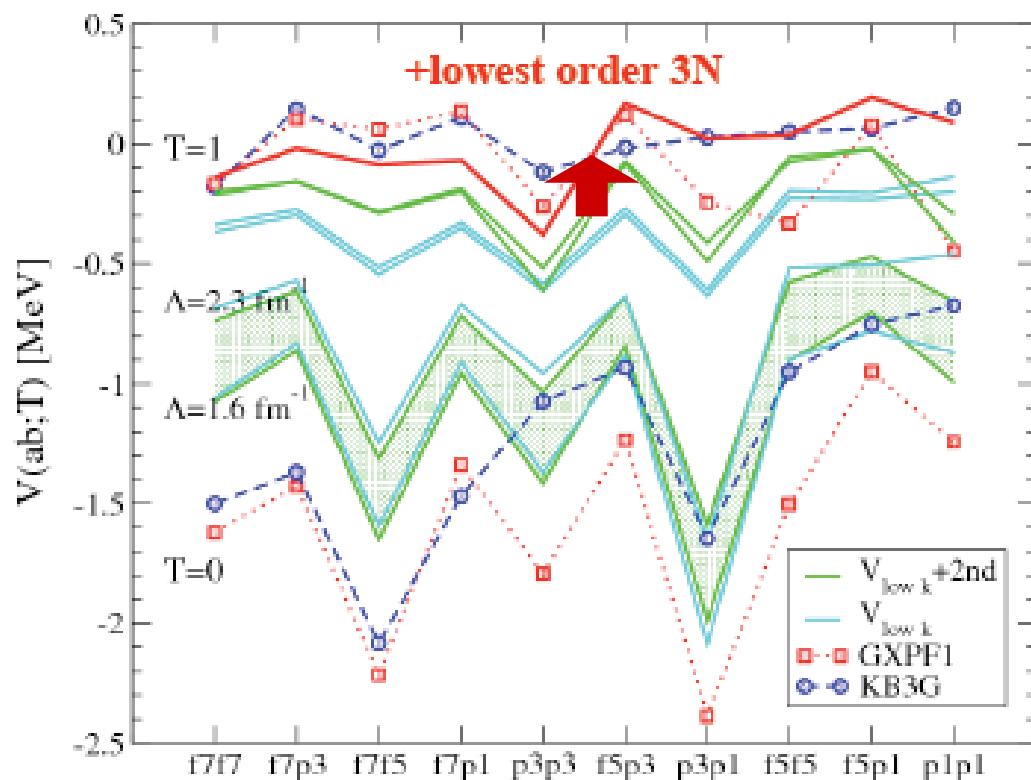




From EFT (Effective Field Theory)

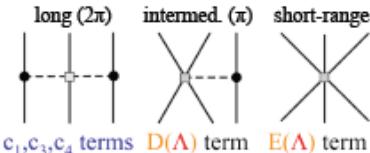
← sd shell

pf shell



Low-momentum 3N interactions

from leading $N^2\text{LO}$ chiral EFT $\sim (Q/\Lambda)^3$ van Kolck (1994), Epelbaum et al. (2002)

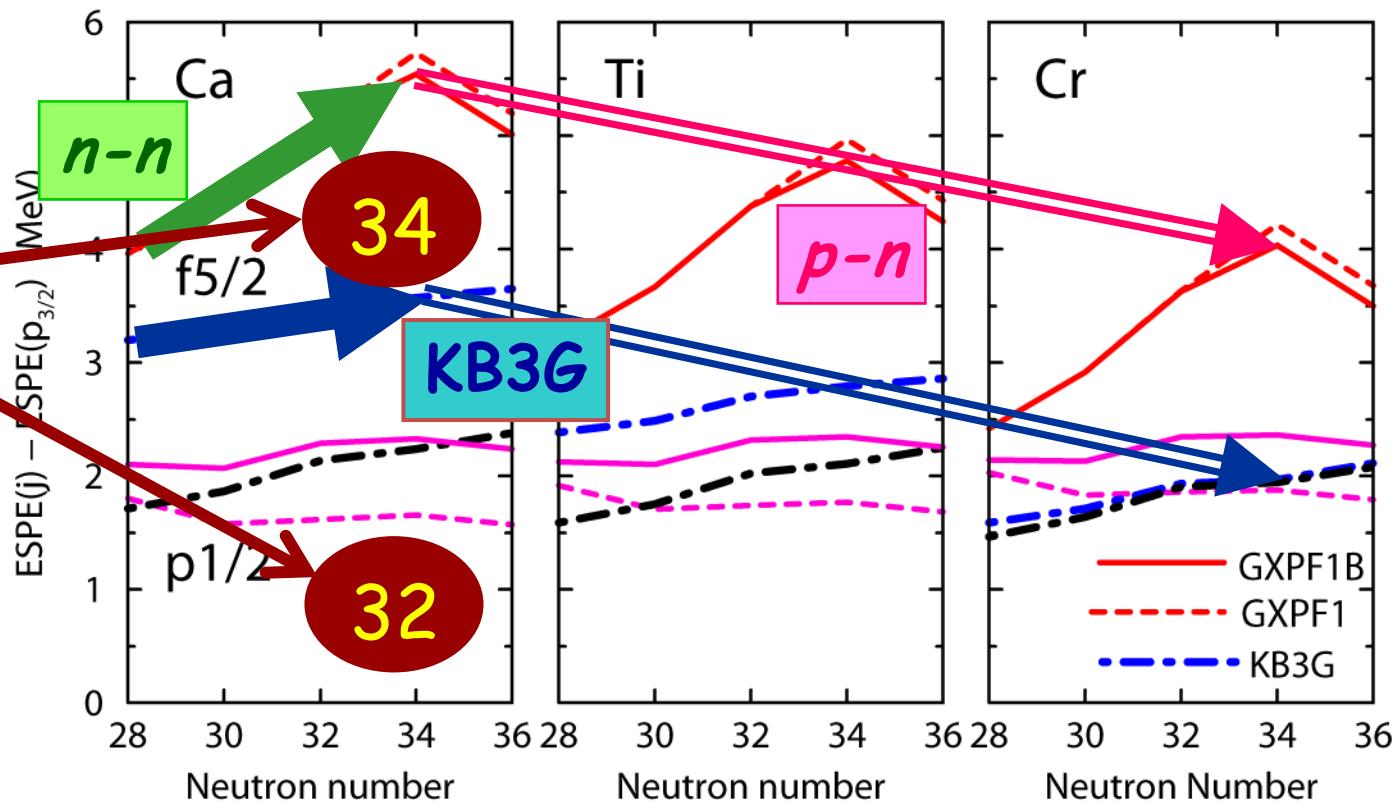


c_i from πN , consistent with NN
Meissner (2007)

$c_1 = -0.9^{+0.2}_{-0.5}$	$c_3 = -4.7^{+1.2}_{-1.6}$	$c_4 = 3.5^{+0.5}_{-0.2}$
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(Effective) single-particle energies

new
magic
numbers ?

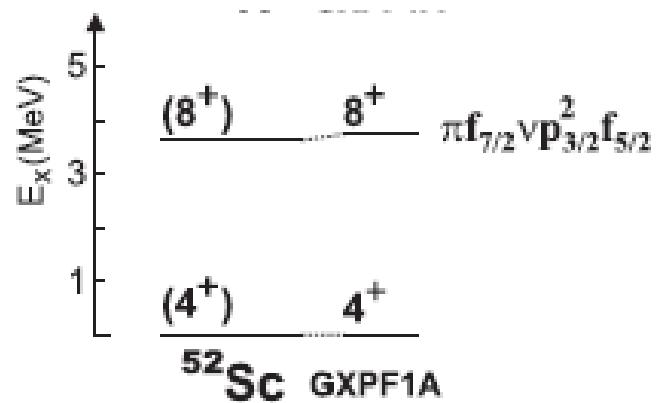


Lowering of $f5/2$ from Ca to Cr :
 $\sim 1.6 \text{ MeV} = 1.1 \text{ MeV (tensor)} + 0.5 \text{ MeV (central)}$

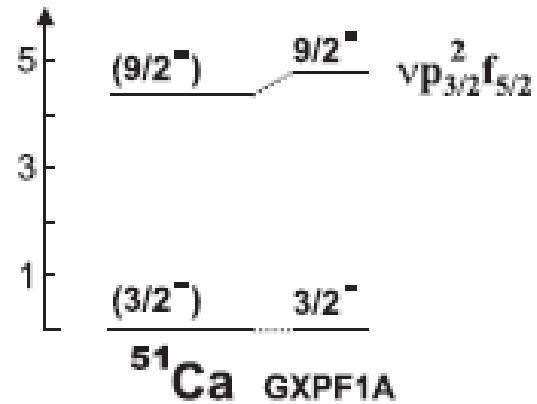
Rising of $f5/2$ from ^{48}Ca to ^{54}Ca :
 $p3/2-p3/2$ attraction $\leftrightarrow p3/2-f5/2$ repulsion

- N=34 magic number

The same mechanism widens N=34 gap.



Recently more (indirect)
experimental indications, e.g.,
Fornal et al. PR C77 (2008)



まとめ

殻構造は原子核構造の根幹

例：低励起での変形はヤーンテラー効果 → 殻効果

殻構造は不安定核で変わり、核力（ハドロン間相互作用）の特定の成分が特徴的なパターンで関与

1. テンソル力 ($\rightarrow ls$ splitting の変化: hw と同程度のオーダー、グルーピング)
2. 中心力 (\rightarrow 動径波動関数のノード)
3. 3体力 (藤田一宮沢力の Monopole 成分 = robust に斥力
 \Rightarrow バレンス核子間の $T=1$ 斥力補正)

帰結



$N=20$ 魔法数の消滅 \times ギャップは大きい今まで、ギャップ越えの励起が大
(WBB の Island of Inversion のアイデア)
○ ギャップそのものが変動し、小さくなる



14:46 3NF ~ EFT ~ neutron matter 不安定核がハドロン力理論の検証の場