

## Beta-NMR study of $^{58}\text{Cu}$ in Si

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Cu impurities in Si devices are considered serious contaminants. The short-lived  $\beta$  emitter  $^{58}\text{Cu}$  ( $I^\pi = 1^+$ ,  $T_{1/2} = 3.2$  s) is attractive for studying the behavior of Cu impurities in Si using the  $\beta$ -NMR technique, which will provide unique information on the mechanism of fast Cu diffusion<sup>1)</sup> or the property of Cu-dopant complex that is related to the gettering technique.<sup>2)</sup> The  $N = Z$  odd-odd nucleus  $^{58}\text{Cu}$ , consisting of  $^{56}\text{Ni}$  plus one proton and one neutron, is also interesting in terms of the nuclear moment, from which we can study the proton-neutron interaction in  $pf$ -shell nuclei.<sup>3)</sup>

We detected an NMR signal of  $^{58}\text{Cu}$  in Si in 2010, and the magnetic dipole moment  $|\mu[^{58}\text{Cu}]| = (0.46 \pm 0.03)\mu_N$  was obtained.<sup>4)</sup> In 2011, Vingerhoets et al. greatly improved the measurement accuracy using collinear laser spectroscopy, achieving  $\mu[^{58}\text{Cu}] = +(0.570 \pm 0.002)\mu_N$ ,<sup>5)</sup> which is about 20% larger than ours. One possibility for the discrepancy is the existence of an electric field gradient (EFG) which could be generated if some defects are formed at a  $^{58}\text{Cu}$  site in Si, though a cubic symmetry site without EFG is expected in terms of the crystal structure of Si. If the EFG exists, the NMR spectrum should split into two lines with frequencies of  $\nu_{\pm} = \nu_0 \pm \nu_Q/2$  in the case of  $I = 1$ , where  $\nu_0$  and  $\nu_Q$  are the carrier frequency and the quadrupole splitting frequency, respectively. In this case, our previous NMR line may have originated from  $\nu_-$ . In the present study, we have applied the multi-radiofrequency (RF) ( $\beta$ -NQR) technique<sup>6)</sup> to the  $\beta$ -NMR measurement of  $^{58}\text{Cu}$  in Si to search for a quadrupole splitting in order to verify the above picture and to solve the discrepancy problem.

The experimental method is similar to our previous one.<sup>4)</sup> Spin-polarized  $^{58}\text{Cu}$  nuclei were produced through the charge exchange reaction of  $^{58}\text{Ni}$  by im-

pinging a  $^{58}\text{Ni}$  primary beam at 63 MeV/u, provided by the RIKEN ring cyclotron with a typical intensity of 100 particle nA, on a 0.5-mm thick Be target. Fully stripped  $^{58}\text{Cu}^{29+}$  ions were separated by the RIKEN projectile fragment separator (RIPS) and were implanted into a single crystal sample of B-doped Si at 15 K with the crystal (001) orientation set parallel to the external magnetic field  $B_0 = 0.93$  T, the same condition as in the previous experiment. A pair of resonance frequencies  $\nu_{\pm}$  was searched for by changing both  $\nu_0$  and  $\nu_Q$ , using the  $\beta$ -NQR technique in which two frequencies were applied in series as  $\nu_- \rightarrow \nu_+ \rightarrow \nu_-$  during the RF duration to inverse spin polarization of  $^{58}\text{Cu}$ .

The resonance was found at  $\nu_0 \sim 4.1$  MHz and  $\nu_Q \sim 2.6$  MHz. The  $\nu_Q$  spectrum at  $\nu_0 = 4.00\text{--}4.15$  MHz is shown in Fig. 1, from which  $\nu_Q = 4eqQ/3h(3\cos^2\theta - 1) = (2.6 \pm 0.4)$  MHz was obtained. The EFG  $q$  will be obtained using the known  $Q$  moment of  $^{58}\text{Cu}$ <sup>5)</sup> after determining the angle  $\theta$  between the main axis of the EFG and  $B_0$  from the crystal orientation dependence of  $\nu_Q$ .  $|\mu[^{58}\text{Cu}]| = (0.58 \pm 0.01)\mu_N$  was obtained from  $\nu_0$ , which is in agreement with the data reported by Vingerhoets et al.<sup>5)</sup> The difference between the present  $\mu[^{58}\text{Cu}]$  and the previous one<sup>4)</sup> is mostly explained by the quadrupole splitting.

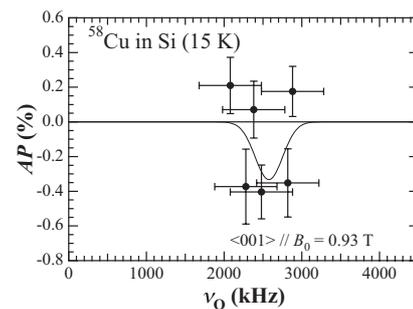


Fig. 1. Beta-NQR spectrum of  $^{58}\text{Cu}$  in Si.

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