

# Na dynamics in the quasi-one-dimensional ionic conductor $\text{NaM}_2\text{O}_4$ ( $M=\text{Ti}$ and $\text{V}$ )

I. Umegaki,<sup>\*1</sup> Y. Higuchi,<sup>\*1</sup> M. Månsson,<sup>\*2</sup> H. Sakurai,<sup>\*3</sup> I. Kawasaki,<sup>\*4</sup> I. Watanabe,<sup>\*4</sup> and J. Sugiyama<sup>\*1</sup>

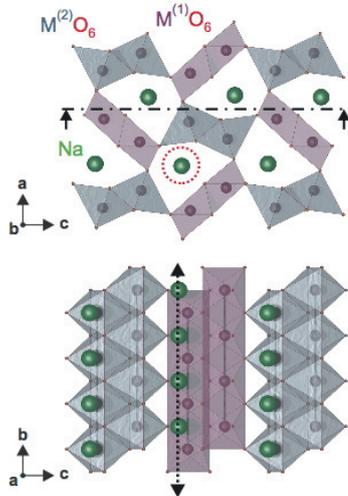


Fig. 1. Crystal structure of  $\text{NaM}_2\text{O}_4$ .

In the  $\text{NaM}_2\text{O}_4$  lattice with a  $\text{CaFe}_2\text{O}_4$ -type orthorhombic structure, the  $\text{Na}^+$  ions are located at the center of a one-dimensional (1D) tunnel along the  $b$ -axis, which is formed by 1D double chains consisting of edge-sharing  $\text{MO}_6$  octahedra ( $M$ : transition metal) (see Fig. 1). The physical properties of  $\text{NaM}_2\text{O}_4$  are reported to strongly depend on  $M$ . In particular, it is very important to clarify their  $\text{Na}^+$ -ion conductivity ( $\sigma_{\text{Na}}$ ) and/or  $\text{Na}^+$ -ion diffusion coefficient ( $D_{\text{Na}}$ ) when using  $\text{NaM}_2\text{O}_4$  as a solid electrolyte in an all-solid-state Na-ion battery.

Following the preliminary report on  $\text{NaV}_2\text{O}_4$ <sup>1)</sup>, we explain here in the results of  $\mu^+$ SR measurements on  $\text{NaM}_2\text{O}_4$  ( $M=\text{Ti}$  and  $\text{V}$ ). The former is a semiconductor with a small band gap<sup>2)</sup>, while the latter is a half metal with anisotropic electric conductivity<sup>3)</sup>. Both ZF- and LF- $\mu^+$ SR spectra were measured in the temperature ( $T$ ) range between 145 and 500 K. The obtained spectra were fitted by a combination of an exponentially relaxing dynamic Kubo-Toyabe signal from a sample and a non-relaxing background signal from a titanium sample holder.

Figure 2 shows the  $T$  dependences of field fluctuation rate ( $\nu$ ), field distribution width ( $\Delta$ ), and exponential relaxation rate ( $\lambda$ ) for (a)  $\text{NaTi}_2\text{O}_4$  and (b)  $\text{NaV}_2\text{O}_4$ . For  $\text{NaTi}_2\text{O}_4$ , as  $T$  increases from 150 K,  $\Delta$  slowly decreases, while  $\nu$  increases rapidly particularly above 350 K. This indicates that the local nuclear magnetic

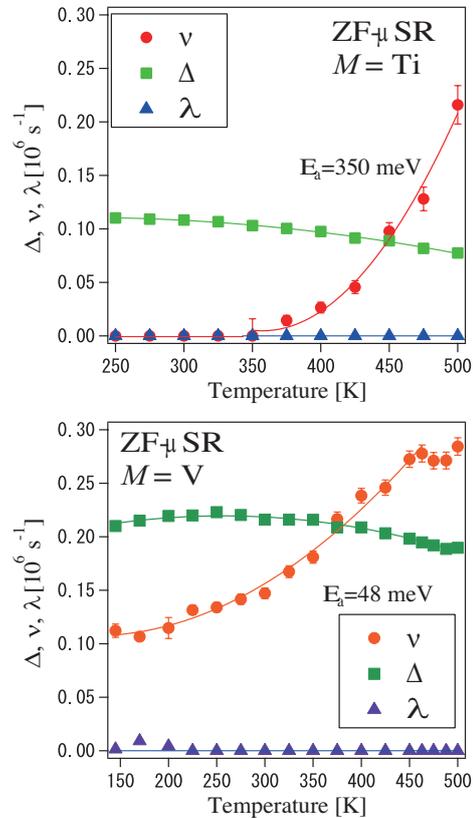


Fig. 2.  $T$ -dependences of field fluctuation rate ( $\nu$ ), field distribution width ( $\Delta$ ), and exponential relaxation rate ( $\lambda$ ) for (a)  $\text{NaTi}_2\text{O}_4$  and (b)  $\text{NaV}_2\text{O}_4$ .

field experienced by  $\mu^+$  starts to fluctuate because of  $\text{Na}^+$  diffusion. For  $\text{NaV}_2\text{O}_4$ , on the other hand, even at 150 K  $\nu$  is comparable to that for  $\text{NaTi}_2\text{O}_4$  at 450 K. This indicates that  $\text{Na}^+$  ions diffuse even at 150 K in  $\text{NaV}_2\text{O}_4$ . The anomaly around 450 K in the  $\nu(T)$  curve is probably caused by a structural phase transition.

If we assume a thermal activation process for the  $T$  dependence of  $\nu$ , the activation energy ( $E_a$ ) is estimated to be 350 meV for  $\text{NaTi}_2\text{O}_4$  and 48 meV for  $\text{NaV}_2\text{O}_4$ . Since the simple Nernst-Einstein equation states that  $\sigma_{\text{Na}} \propto D_{\text{Na}}$ , where  $D \propto \nu$ ,  $\text{NaV}_2\text{O}_4$  is expected to be a good candidate for a  $\text{Na}^+$ -ionic conductor.

## References

- 1) J. Sugiyama et al.: RIKEN Accel. Prog. Rep. 45, 197 (2012).
- 2) M. J. Geselbracht et al.: J. Solid State Chem. 179, 3489 (2006).
- 3) K. Yamaura et al.: Phys. Rev. Lett. 99, 196601 (2007).

<sup>\*1</sup> Toyota Central Research and Development Labs., Inc.

<sup>\*2</sup> École polytechnique fédérale de Lausanne and Paul Scherrer Institut

<sup>\*3</sup> National Institute for Materials Science (NIMS)

<sup>\*4</sup> RIKEN Nishina Center