

Three-body model calculation of the 2^+ state in $^{26}\text{O}^\dagger$

K. Hagino^{*1,*2} and H. Sagawa^{*3,*4}

We discuss the 2^+ state of ^{26}O using a three-body model of an $^{24}\text{O}+n+n$ system with full account of the continuum. The decay energy spectrum for a given angular momentum I can be evaluated as

$$\frac{dP_I}{dE} = \sum_k |\langle \Psi_k^{(I)} | \Phi_{\text{ref}}^{(I)} \rangle|^2 \delta(E - E_k), \quad (1)$$

where $\Psi_k^{(I)}$ is a solution of the three-body model Hamiltonian with angular momentum I and energy E_k , and $\Phi_{\text{ref}}^{(I)}$ is the wave function for a reference state with the same angular momentum. For a reference state we use the uncorrelated state of ^{27}F with the neutron $[[1d_{3/2} \otimes 1d_{3/2}]^{(IM)}]$ configuration, which is dominant in the ground state of ^{27}F .

With a contact interaction, the continuum effects on the decay energy spectrum can be taken into account in terms of the Green's function. Notice that Eq. (1) can be expressed as

$$\begin{aligned} \frac{dP_I}{dE} &= -\frac{1}{\pi} \Im \sum_k \langle \Phi_{\text{ref}}^{(I)} | \Psi_k^{(I)} \rangle \frac{1}{E_k - E - i\eta} \langle \Psi_k^{(I)} | \Phi_{\text{ref}}^{(I)} \rangle, \\ &\equiv -\frac{1}{\pi} \Im \langle \Phi_{\text{ref}}^{(I)} | G^{(I)}(E) | \Phi_{\text{ref}}^{(I)} \rangle, \end{aligned} \quad (2)$$

where \Im denotes the imaginary part and η is an infinitesimal number and $G^{(I)}(E)$ is the correlated Greens's function. The correlated Greens's function will be constructed using the unperturbed Green's function.

The upper panel of Fig. 1 shows the decay energy spectrum of ^{26}O for $I=0$ (dashed line) and $I=2$ (solid line). For presentation purposes, we set η in Eq. (2) to be a finite value, i.e., $\eta = 0.21 \text{ MeV}^1$. For comparison, we also show the spectrum for the uncorrelated case with a dotted line, which gives the same spectrum both for $I=0$ and $I=2$. For the uncorrelated case, the spectrum has a peak at $E = 1.54 \text{ MeV}$, which is twice the single-particle resonance energy, 0.77 MeV . With the pairing interaction between the valence neutrons, the peak energy shifts towards lower energies. The energy shift ΔE is larger in $I=0$ than in $I=2$, i.e., the peak in the spectrum appears at $E = 0.148 \text{ MeV}$ ($\Delta E = -1.392 \text{ MeV}$) for $I=0$ and at $E = 1.354 \text{ MeV}$ ($\Delta E = -0.186 \text{ MeV}$) for $I=2$.

We have shown that the 2^+ state appears at approximately $E = 1.35 \text{ MeV}$. This 2^+ energy is close

to, but slightly smaller than, the unperturbed energy, $E = 1.54 \text{ MeV}$, and thus the energy shift from the unperturbed energy is much smaller than the energy shift for the 0^+ state. We have argued that this is a typical spectrum well understood by the single- j model with the pairing residual interaction. Many shell model calculations such as the ab initio³⁾ and USDA and USDB⁴⁾ calculations have predicted the excitation energy of the 2^+ state in ^{26}O in the opposite trend, i.e., they have predicted a higher energy than the unperturbed energy. The energy of the 2^+ state needs to be urgently confirmed experimentally⁵⁾ in order to clarify the validity of nuclear models and effective interactions in nuclei on and beyond the neutron drip-line.

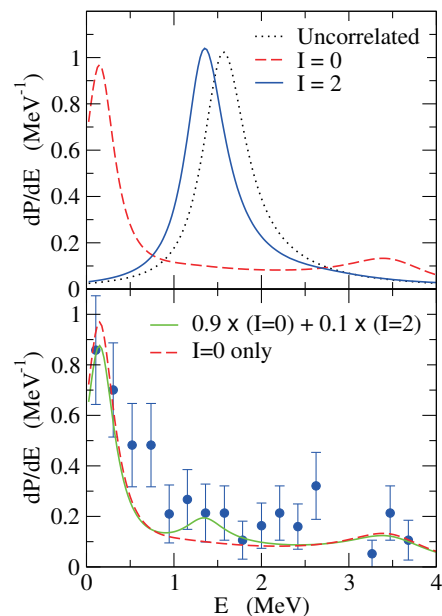


Fig. 1. (upper panel) The decay energy spectrum for the two-neutron emission decay of ^{26}O . The dashed and solid lines represent the 0^+ and 2^+ states, respectively. The dotted line shows the uncorrelated spectrum obtained by ignoring the interaction between the valence neutrons. (lower panel) The decay energy spectrum obtained by superposing the $I=0$ and $I=2$ components. The dashed line is the decay energy spectrum for the pure $I=0$ configuration. The experimental data, normalized to the unit area, are taken from Ref.²⁾.

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^{*1} Department of Physics, Tohoku University

^{*2} Research Center for Electron Photon Science, Tohoku University

^{*3} RIKEN Nishina Center

^{*4} Center for Mathematics and Physics, University of Aizu

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