

The ($^{16}\text{O}, ^{16}\text{F}(0^-)$) reaction to study spin-dipole 0^- states

M. Dozono,^{*1} T. Uesaka,^{*2} K. Fujita,^{*3} N. Fukuda,^{*2} M. Ichimura,^{*2} N. Inabe,^{*2} S. Kawase,^{*1} K. Kisamori,^{*1} Y. Kiyokawa,^{*1} K. Kobayashi,^{*4} M. Kobayashi,^{*1} T. Kubo,^{*2} Y. Kubota,^{*1} C. S. Lee,^{*1} M. Matsushita,^{*1} S. Michimasa,^{*1} H. Miya,^{*1} A. Ohkura,^{*3} S. Ota,^{*1} H. Sagawa,^{*2,*5} S. Sakaguchi,^{*3} H. Sakai,^{*2} M. Sasano,^{*2} S. Shimoura,^{*1} Y. Shindo,^{*3} L. Stuhl,^{*2} H. Suzuki,^{*2} H. Tabata,^{*3} M. Takaki,^{*1} H. Takeda,^{*2} H. Tokieda,^{*1} T. Wakasa,^{*3} K. Yako,^{*1} Y. Yanagisawa,^{*2} J. Yasuda,^{*3} R. Yokoyama,^{*1} K. Yoshida,^{*2} and J. Zenihiro^{*2}

We proposed the parity-transfer ($^{16}\text{O}, ^{16}\text{F}(0^-)$) reaction as a powerful tool to study spin-dipole (SD) 0^- states in nuclei¹⁾. The parity-transfer reaction has a unique selectivity to unnatural-parity states, which is an advantage over the other reactions used thus far. As the first ($^{16}\text{O}, ^{16}\text{F}(0^-)$) measurement, the experiment for a ^{12}C target was performed with the SHARAQ spectrometer. The known 0^- state at $E_x = 9.3$ MeV in $^{12}\text{B}^2)$ serves as a benchmark to verify the effectiveness of this reaction. The experimental setup and method can be found in Ref.³⁾.

The preliminary result of the excitation-energy spectrum for the $^{12}\text{C}(^{16}\text{O}, ^{16}\text{F}(0^-))^{12}\text{B}$ reaction at $\theta_{\text{lab}} = 0^\circ - 0.25^\circ$ is shown in Fig. 1. The energy resolution is 2.6 MeV in FWHM. We note that the events at $E_x \sim -10$ MeV are due to the ($^{16}\text{O}, ^{16}\text{F}(0^-)$) reaction on hydrogens in the plastic scintillator used as a reaction target. The obtained distribution is largely different from those previously obtained by other reaction probes such as (n, p) and ($d, ^2\text{He}$) (e.g., see Fig. 1 in Ref.⁴⁾). A distinct difference is that the 1^+ ground-state Gamow-Teller (GT) transition is strongly hindered. Furthermore, an enhancement at $E_x \sim 9$ MeV can be seen, which is due to the 0^- state at $E_x = 9.3$ MeV. Therefore, the obtained distribution shows high selectivity of the present reaction to 0^- states.

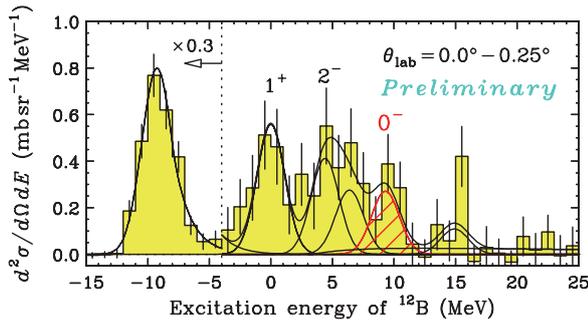


Fig. 1. Preliminary result of the excitation-energy spectrum for the $^{12}\text{C}(^{16}\text{O}, ^{16}\text{F}(0^-))^{12}\text{B}$ reaction at $\theta_{\text{lab}} = 0^\circ - 0.25^\circ$.

In order to extract the yield of the 0^- state, we performed peak fitting, and the results are shown as the solid lines in Fig. 1. Figure 2 shows the angular

distribution for the 0^- state at $E_x = 9.3$ MeV together with the results for the 1^+ g.s. and the 2^- state at $E_x = 4.4$ MeV. The solid curves denote the results calculated by the distorted-wave Born approximation (DWBA). The DWBA calculations predict the oscillatory patterns of the cross sections that are different depending on the spin-parity. The 0^- has the strong forward peaking, while the other states, 1^+ and 2^- , have the first maximum at finite angles. These patterns reproduce the experimental data very well. Thus, the oscillatory pattern of the angular distribution allows a clear spin-parity determination for the unnatural-parity state. Further analysis is underway to finalize experimental results.

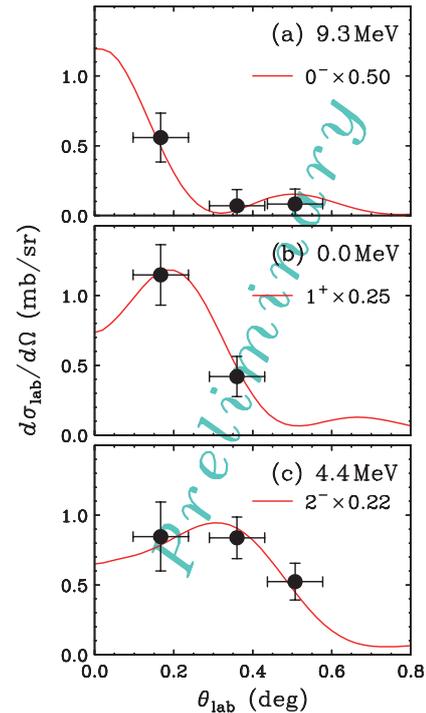


Fig. 2. Measured and calculated differential cross sections for $E_x = 9.3$ MeV, 0.0 MeV, and 4.4 MeV states.

References

- 1) M. Dozono et al.: RIKEN Accel. Prog. Rep. **45**, 10 (2012).
- 2) H. Okamura et al.: Phys. Rev. C **66**, 054602 (2002).
- 3) M. Dozono et al.: RIKEN Accel. Prog. Rep. **48**, 58 (2015).
- 4) H. Okamura et al.: Phys. Lett. B **345**, 1 (1995).

^{*1} Center for Nuclear Study, University of Tokyo

^{*2} RIKEN Nishina Center

^{*3} Department of Physics, Kyushu University

^{*4} Rikkyo University

^{*5} Center for Mathematics and Physics, University of Aizu