## Search for $\bar{K}NN$ bound state at J-PARC

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The possible existence of strongly bound  $\bar{K}$  nuclear states has been widely discussed as a consequence of the strongly attractive  $\bar{K}N$  interaction in the I = 0channels.<sup>1)</sup> The investigation of those exotic states will provide unique information of the  $\bar{K}N$  interaction below the threshold, which is still not fully understood. Among the  $\bar{K}$  nuclear-states, the simplest  $\bar{K}$ -nuclear cluster  $\bar{K}NN$ , is of special interest because it is the lightest S = -1  $\bar{K}$  nucleus. Many theoretical calculations based on the  $\bar{K}NN - \pi\Sigma N - \pi\Lambda N$  coupled channels have predicted the existence of the bound state.<sup>2)</sup> However, the calculated properties of the bound state, such as the binding energy (B.E.) and decay width ( $\Gamma$ ), strongly depend on  $\bar{K}N$  interaction models.

Experimentally, there are several reports on the observation of peak structure with a B.E. of  $\sim$  100 MeV.<sup>3–5)</sup> In contrast, null results were also reported by several experiments,  $^{6,7)}$  and therefore the existence of the  $\bar{K}NN$  bound state has remained controversial. The J-PARC E15 experiment was performed at the K1.8BR beam line using the  $K^{-}$  + <sup>3</sup>He reaction at  $p_{K^{-}}$ = 1.0 GeV/c.<sup>8)</sup> The first physics run of the E15 experiment was conducted in May 2013 with  $5.3 \times 10^9$ kaons incident on <sup>3</sup>He. The results of a semi-inclusive  ${}^{3}\text{He}(K^{-}, n)X$  measurement can be found in Ref<sup>9</sup>. To measure an expected decay mode of  $\bar{K}NN \to \Lambda p$ , an exclusive  ${}^{3}\text{He}(K^{-},\Lambda p)n$  measurement was also performed by identifying two protons and one negative pion in a cylindrical detector system (CDS) surrounding a liquid <sup>3</sup>He target system.<sup>10</sup> The  $p\pi^-$  pair associated with a specific  $\Lambda$  decay in the  $pp\pi^-$  event was reconstructed with a log-likelihood method, and a missing neutron was identified kinematically.

The distribution of the  $\Lambda p$  invariant mass and the neutron emission angle of the obtained  $\Lambda pn$  final state are shown in Fig. 1 (b) and (c), respectively, and a scatter plot of the two is given in Fig. 1 (a). A bump structure is observed at the mass threshold of  $K^- + p + p$ in the  $\Lambda p$  invariant mass spectrum, which is not reproduced by a global fit based on multi-nucleon absorption processes of  $K^-$ . To explain the excess, a simple Swave Breit-Wigner structure over the three-body phase space of the  $K^-$  <sup>3</sup>He  $\rightarrow \Lambda pn$  reaction is assumed. A  $\chi^2$ test was performed between the experimental data and the simulated pole together with the multi-nucleon absorption processes, and the minimum  $\chi^2$  point was obtained with  $M_X = 2355^{+6}_{-8}(stat.) \pm 12(syst.) \text{ MeV}/c^2$ ,  $\Gamma_X = 110^{+19}_{-17}(stat.) \pm 27(syst.) \text{ MeV}/c^2$ .



Fig. 1. (a) Distribution of the  $\Lambda$  invariant mass and the emission angle of the missing neutron, (b)  $\Lambda$  invariant mass with simulated spectra obtained by the global fit, and (c) angular distribution of the missing neutron.<sup>10</sup>

A naive interpretation of the structure would be the  $\bar{K}NN$  bound state, since the pole position is located below the mass threshold of  $K^- + p + p$  (2370 MeV/ $c^2$ ). It could also be a shallow bound or unbound resonance of the  $\Lambda(1405)p$  system followed by  $\Lambda p$  conversion, the threshold of which is located at 2343 MeV/ $c^2$  if the mass of  $\Lambda(1405)$  is assumed to be 1405 MeV/ $c^2$ . To explore whether or not the observed structure is the  $\bar{K}NN$  bound state in more detail, the second run was carried out in 2015 with the accumulation of 30 times more data on the  $\Lambda pn$  final state and the analysis is in progress.

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