

# Disentangling transverse single spin asymmetries for forward neutrons in high-energy polarized-proton–nucleus collisions

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It is reported from the PHENIX experiment at BNL-RHIC that the transverse single spin asymmetry, denoted as  $A_N$ , for forward neutrons measured in transversely polarized-proton–nucleus ( $pA$ ) collisions at  $\sqrt{s_{NN}} = 200\text{ GeV}$  is far different from that in proton–proton ( $pp$ ) collisions at  $\sqrt{s} = 200\text{ GeV}$ <sup>1)</sup> (see panel (b) of Fig. 1.)

In this report, I present an important but rather unknown mechanism: ultra-peripheral  $pA$  collisions (UPCs, also known as Primakoff effects). UPCs contribute to the measured  $A_N$  modestly in  $pAl$  collisions and significantly in  $pAu$  collisions. UPCs occur when the impact parameter  $b$  is larger than the sum of the radii of each colliding particle, namely  $b > R_p + R_A$  ( $R_p$  and  $R_A$  are the radius of the proton and nucleus, respectively). In UPCs, virtual photons ( $\gamma^*$ ) emitted from the relativistic nucleus interact with the polarized protons and then produce the neutrons and other particles.

The differential cross section for single pion and neutron production, dominant among many other channels, in UPCs is given by

$$\frac{d\sigma_{\text{UPC}(pA \rightarrow \pi^+n)}^4}{dW db^2 d\Omega_n} = \frac{d^3 N_{\gamma^*}}{dW db^2} \frac{d\sigma_{\gamma^* p \rightarrow \pi^+n}(W)}{d\Omega_n} \overline{P}_{\text{had}}(b), \quad (1)$$

where  $d^3 N_{\gamma^*}/d\omega_{\gamma^*}^{rest} db^2$  is the double differential photon flux due to the fast-moving nucleus,  $W$  is the  $\gamma^* p$  center-of-mass energy,  $d\Omega_n = \sin\Theta d\Theta d\Phi$  with the neutron scattering polar angle  $\Theta$  and azimuthal angle  $\Phi$  in the  $\gamma^* p$  center-of-mass frame, and  $\overline{P}_{\text{had}}(b)$  is the probability of having no hadronic interactions in  $pA$  collisions at a given  $b$ . Single neutron and pion productions from the  $\gamma^* p$  interaction are simulated following the differential cross sections predicted by the MAID 2007 model.<sup>2)</sup> The cross section of the  $\gamma^* p \rightarrow \pi^+ n$  interaction is approximated as

$$\frac{d\sigma_{\gamma^* p \rightarrow \pi^+ n}}{d\Omega_\pi} \propto R_T^{00} \left( 1 + \cos\Phi \frac{R_T^{0y}}{R_T^{00}} \right), \quad (2)$$

where  $R_T^{00}$  and  $R_T^{0y}$  are the response functions for pion photoproduction.  $A_N$  for forward neutrons in UPCs (hereafter  $A_N^{\text{UPC}}$ ) inherits the target asymmetry  $T(\pi - \Theta) \equiv R_T^{0y}/R_T^{00}$  in Eq. (2), which is  $\sim 0.7$  at  $W < 1.3\text{ GeV}$  and  $\sim -0.2$  at  $W > 1.3\text{ GeV}$  within the PHENIX detector acceptance. Owing to the virtual photon flux leading to low-energy photons and the pion photoproduction cross section via a  $\Delta(1232)$  resonance, UPCs accordingly provide  $A_N^{\text{UPC}} \sim 0.35$  for forward neutrons.

Figure 1 (a) shows the differential cross section in  $pAu$  collisions,  $d\sigma/d\Phi$ , as a function of  $\Phi$ , for UPCs (dashed

red line) and a one-pion exchange model (OPE) that represents hadronic interactions occurring at  $b < R_p + R_A$  (solid black line).  $A_N$  originating in OPE well explains the PHENIX result in  $pp$  collisions but does not in  $pA$  collisions. In this study, the Glauber multiple scattering model is applied to OPE to account for nuclear effects. Here we find that UPCs have a positive and large  $A_N^{\text{UPC}}$  compared with  $A_N^{\text{OPE}} = -0.05$  of hadronic interactions.

In Fig. 1 (b), filled black circles indicate  $A_N$  inclusively measured by the PHENIX zero-degree calorimeter<sup>1)</sup>. These  $A_N$  values can be compared with open red circles that correspond to the sum of UPCs and OPE MC simulations, denoted as  $A_N^{\text{UPC+OPE}}$  and calculated as

$$A_N^{\text{UPC+OPE}} = \frac{\sigma_{\text{UPC}} A_N^{\text{UPC}} + \sigma_{\text{OPE}} A_N^{\text{OPE}}}{\sigma_{\text{UPC}} + \sigma_{\text{OPE}}}, \quad (3)$$

where  $\sigma_{\text{UPC}}$  and  $\sigma_{\text{OPE}}$  are the cross sections of UPCs and OPE, respectively. In  $pAu$  collisions, since  $\sigma_{\text{UPC}} \simeq \sigma_{\text{OPE}}$ , we obtain  $A_N^{\text{UPC+OPE}} = 0.16$ , which is consistent with the PHENIX result. Consistency between our simulation result  $A_N^{\text{UPC+OPE}} = -0.02$  and the PHENIX data is also found in  $pAl$  collisions, where  $\sigma_{\text{UPC}}$  is 8% of  $\sigma_{\text{OPE}}$ .

In the MC simulations discussed in this report, electromagnetic effects (UPCs) and hadronic effects (OPE) are taken into account independently. However, the interference between these two effects, called the Coulomb-nuclear interference, would have nonzero amplitudes in the very small momentum-transfer region. The implementation of the Coulomb-nuclear interference will be a topic of future investigation.

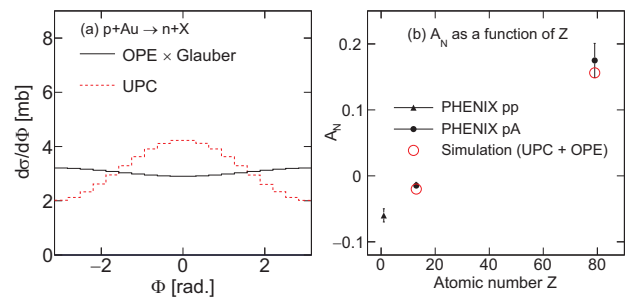


Fig. 1. Left:  $d\sigma/d\Phi$  distributions for UPCs and OPE. Right: Comparison of  $A_N$  as a function of  $Z$ .

## References

- 1) I. Nakagawa et al. (PHENIX Collaboration), J. Phys. Conf. Ser. **736**, 012017 (2016).
- 2) D. Drechsel, S. S. Kamalov, L. Tiator, Eur. Phys. J. A **34**, 64 (2007).

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