

# Half-life determination of nuclear excited states of $^{229}\text{Th}$ by the coincidence measurement between $\alpha$ particles and $\gamma$ rays from $^{233}\text{U}$

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Studying the nuclear structure of  $^{229}\text{Th}$  is important in that reflection-symmetric and reflection-asymmetric shapes are expected to coexist<sup>1)</sup> and the first nuclear excited isomeric state of  $^{229}\text{Th}$  ( $^{229\text{m}}\text{Th}$ ) has an extremely low energy of 8.3 eV,<sup>2)</sup> which would enable the development of nuclear laser spectroscopy and an ultraprecise nuclear clock. In this study, we measured the half-lives of nuclear excited states of  $^{229}\text{Th}$ , populated through the alpha decay of  $^{233}\text{U}$ , for further understanding of the nuclear properties of  $^{229}\text{Th}$ . This work was also aimed at determining the half-life of the 29.2-keV state (Fig. 1). This half-life was one of the key parameters for exciting  $^{229}\text{Th}$  from the ground state to the 29.2-keV state with synchrotron radiation to actively produce  $^{229\text{m}}\text{Th}$ , which was recently realized by our group.<sup>3)</sup>

For the half-life determination, we performed a coincidence measurement between  $\alpha$  particles and  $\gamma$  rays from a  $^{233}\text{U}$  source prepared by electrodeposition. The  $^{233}\text{U}$  source was placed inside a vacuum chamber enclosed by 5-cm Pb blocks for background  $\gamma$ -ray reduction. Emitted  $\alpha$  particles from the source were measured with a passivated implanted planar silicon detector mounted inside the chamber, and  $\gamma$  rays that were emitted from the source and passed through a polyimide window out of the vacuum were measured with a  $\text{LaBr}_3(\text{Ce})$  scintillator and photomultiplier. After amplification,  $\alpha$ -particle

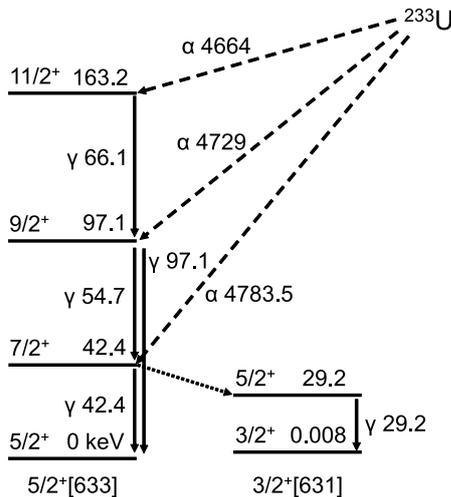


Fig. 1. Excited states of  $^{229}\text{Th}$  the half-lives of which were determined, and  $\alpha$  and  $\gamma$  transitions used for determining the half-lives (unit: keV).

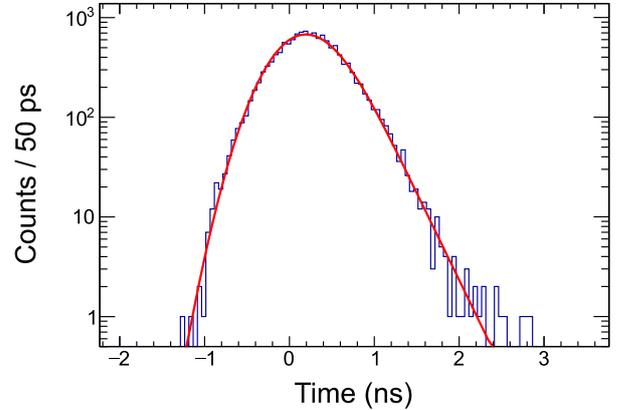


Fig. 2. Time spectrum of  $\gamma$ -ray signals following  $\alpha$ -particle signals for 4783.5-keV  $\alpha$  particles and 42.4-keV  $\gamma$  rays (histogram). A single exponential decay function convoluted with a Gaussian function is fitted to the data (curve).

and  $\gamma$ -ray signals were recorded with Computer-Aided Measurement And Control (CAMAC) modules to obtain time spectra of  $\gamma$ -ray signals following  $\alpha$ -particle signals for each combination of  $\alpha$ -particle and  $\gamma$ -ray energies.

Figure 2 shows a time spectrum for 4783.5-keV  $\alpha$  particles and 42.4-keV  $\gamma$  rays. The half-life obtained from the fitting of a single exponential decay function convoluted with a Gaussian function was almost the same as the reference value (172(6) ps).<sup>4)</sup> From a time spectrum for 4783.5-keV  $\alpha$  particles and 29.2-keV  $\gamma$  rays, we obtained the half-life of the 29.2-keV state, which is consistent with that obtained from our synchrotron excitation of  $^{229}\text{Th}$  (82(4) ps).<sup>3)</sup> These coincidences with the previous results indicate that our experiments and data analysis were correctly performed. The half-life of the 97.1-keV state was obtained from  $\gamma$  rays of 54.7 and 97.1 keV detected in coincidence with 4729-keV alpha particles. The obtained half-life was seemingly shorter than the reference value (147(12) ps).<sup>4)</sup> This may be because our measurement has better and time resolutions than the previous measurement.<sup>4)</sup>

We could clearly observe the peak of 66.1-keV  $\gamma$  rays by gating the  $\gamma$ -ray spectrum with  $\alpha$  particles of approximately 4664 keV, which allowed us to obtain the half-life of the 163.2-keV state for the first time. We will continue detailed analysis to precisely determine each half-life and discuss the nuclear properties of  $^{229}\text{Th}$  based on the half-lives determined in this study.

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