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 $^{45}\mathrm{Ti}~(T_{1/2}=184.8~\mathrm{min})$  is an appropriate positron emitter isotope  $(E_{\beta^+}=439~\mathrm{keV},~I_{\beta^+}=84.8\%)$  for positron emission tomography (PET). This radioisotope can be produced in the deuteron-induced reaction on a scandium-45 target at cyclotrons. However, the quality of experimental data on the cross sections of the  $^{45}\mathrm{Sc}(d,2n)^{45}\mathrm{Ti}$  reaction is not satisfactory. Therefore, we aim to measure the cross sections of the  $^{45}\mathrm{Sc}(d,2n)^{45}\mathrm{Ti}$  reaction and to investigate a route for  $^{45}\mathrm{Ti}$  production.

The stacked-foil activation technique and  $\gamma$ -ray spectrometry were adopted to determine the cross sections. The stacked target included metallic foils of  ${
m ^{45}Sc}$  (thicknesses of 25.8 and 250  $\mu$ m with a purity of 99.0%), <sup>27</sup>Al (18.5  $\mu$ m, 99.6%), and <sup>nat</sup>Ti (20.2  $\mu$ m, 99.6%). The target was irradiated for 30 min with a 24-MeV deuteron beam from the RIKEN AVF cyclotron. The incident beam energy was measured using the time-offlight method. The energy degradation in the stacked target was calculated using the SRIM code.<sup>1)</sup> The beam intensity was measured using a Faraday cup and doublechecked with the <sup>nat</sup>Ti $(d, x)^{48}$ V monitor reaction.<sup>2)</sup> By comparing the monitor reaction, the measured intensity  $(180 \pm 9 \text{ nA})$  was corrected by decreasing it by 2% to  $176 \pm 9$  nA. The  $\gamma$ -ray spectra of the irradiated foils were measured using a high-resolution and a high-purity germanium (HPGe) detector. The detector was calibrated using a mixed  $\gamma$ -ray point source. In the measurements, the dead time was kept below 7%.

Subsequently, the activation cross sections of <sup>44,45</sup>Ti and <sup>44g,44m,46</sup>Sc were determined. The measurements of the 719.6-keV  $\gamma$ -ray ( $I_{\gamma} = 0.154\%$ ) from the <sup>45</sup>Ti decay were used to derive the cross sections of the <sup>45</sup>Sc(d, 2n)<sup>45</sup>Ti reaction. Figure 1 shows our measured excitation function of the <sup>45</sup>Sc(d, 2n)<sup>45</sup>Ti reaction in comparison with previous experimental data<sup>3)</sup> and the theoretical estimation retrieved from TENDL-2019.<sup>4)</sup> The derived excitation function is consistent with the data reported by Hermanne *et al.*;<sup>3)</sup> however, it is less scattered. The peak position of the TENDL-2019 data is slightly shifted to a lower energy.

The physical yield of  $^{45}$ Ti was deduced from the measured excitation function and is shown in Fig. 2. The

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450 400 45Sc(d,2n)45Ti 350 300 Cross section (mb) 250 200 150 This work Hermanne 2012 100 50 TENDL-2019 0 10 15 20 25 Energy (MeV)

Fig. 1. Excitation function of the  ${}^{45}Sc(d, 2n){}^{45}Ti$  reaction.



Fig. 2. Physical yield of  $^{45}$ Ti.

present yield curve of <sup>45</sup>Ti is slightly higher than the experimental data obtained by Dmitriev *et al.*<sup>5)</sup> at 22 MeV. <sup>44</sup>Ti is the only one co-produced radioactive isotope of titanium in our experiment and can be formed by (d, 3n) reaction on <sup>45</sup>Sc above 15 MeV. Therefore, isotopically pure <sup>45</sup>Ti production is possible in (d, 2n) reaction on <sup>45</sup>Sc in an energy range of 8–15 MeV.

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