

Reaction plane determination INTT at sPHENIX

M. Fujiwara,^{*1,*2} Y. Akiba,^{*2} J. Bertaux,^{*2,*3} D. Cacace,^{*4} R. G. Cecato,^{*5} A. Enokizono,^{*2} Y. Fujino,^{*2,*6} T. Hachiya,^{*1,*2} T. Harada,^{*2,*6} S. Hasegawa,^{*7} B. Hong,^{*8} J. Hwang,^{*2,*8} M. Ikemoto,^{*1,*2} Y. Ishigaki,^{*1,*2} M. Kano,^{*1,*2} T. Kato,^{*2,*6} T. Kikuchi,^{*2,*6} T. Kondo,^{*9} T. Kumaoka,^{*2} C. M. Kuo,^{*10} R. S. Lu,^{*11} N. Morimoto,^{*1,*2} I. Nakagawa,^{*2} R. Nouicer,^{*4} G. Nukazuka,^{*2} I. Omae,^{*1,*2} R. Pisani,^{*4} Y. Sekiguchi,^{*2} C. W. Shih,^{*2,*10} M. Shimomura,^{*1} R. Shishikura,^{*2,*6} W. C. Tang,^{*2,*10} H. Tsujibata,^{*1,*2} W. Xie,^{*3} and H. Yanagawa^{*2,*6}

The sPHENIX Experiment started operation in May 2023 at the Relativistic Heavy Ion Collider (RHIC) in Brookhaven National Laboratory. One of the main purposes of the sPHENIX is to study the quark gluon plasma (QGP). QGP is a high-temperature state where quarks and gluons are no longer confined inside hadrons. Au + Au collisions at a nucleon-nucleon center-of-mass energy $\sqrt{s_{NN}} = 200$ GeV were collected for Run 2023 and 2024.

When the QGP is produced in non-central nucleus-nucleus collision, elliptic anisotropy (v_2) is observed because the number of particles that are produced by the collision varies depending on the azimuthal angle to the pressure gradient. Figure 1 (left) is the concept of elliptic flow. It shows that particle production is maximal along the event plane direction ($\phi - \psi_2 = 0^\circ$) and minimal in the direction perpendicular to it ($\phi - \psi_2 = 90^\circ$). v_2 is calculated by $v_2 = \langle \cos(2(\phi - \psi)) \rangle$ where ψ is the reaction plane (RP) angle and ϕ is the azimuthal angle of the particle.

The Intermediate Tracker (INTT) is one of the tracking detectors in the sPHENIX experiment. The INTT that covers a pseudorapidity range of $|\eta| < 1.1$ and the full azimuth consists of two barrel-shaped layers located in the radius of 7 to 10 cm from the beam axis. sPHENIX Event Plane Detector (sEPD) covers $2.0 < |\eta| < 4.9$ is a detector to determine event plane.

To measure v_2 at forward rapidity, the RP should be determined at a different rapidity by the INTT. If the v_2 and RP are measured at the same rapidity, v_2 signal is swept out by a significant auto-correlation between the v_2 and RP. Observed RP angle ψ_2^{obs} is determined by

$$\psi_2^{obs} = \frac{1}{2} \tan^{-1} \frac{\sum_i \sin(2\phi_i)}{\sum_i \cos(2\phi_i)}, \quad (1)$$

^{*1} Department of Mathematical and Physical Sciences, Nara Women's University

^{*2} RIKEN Nishina Center

^{*3} Department of Physics and Astronomy, Purdue University

^{*4} Physics Department, Brookhaven National Laboratory

^{*5} Instrumentation Division, Brookhaven National Laboratory

^{*6} Department of Physics, Rikkyo University

^{*7} Advanced Science Research Center, Japan Atomic Energy Agency

^{*8} Department of Physics, Korea University

^{*9} Information Systems Technology Division, Tokyo Metropolitan Industrial Technology Research Institute

^{*10} Department of Physics, National Central University

^{*11} Department of Physics, National Taiwan University

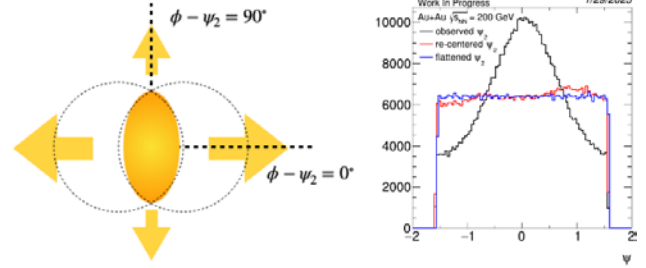


Fig. 1. Left: Elliptic flow illustration showing azimuthal anisotropy relative to the event plane (ψ_2), with maximum particle yield at $\phi - \psi_2 = 0^\circ$ and minimum at 90° . Right: Observed (black), re-centered (red), and flattened (blue) reaction plane distributions.

where the i represents each cluster of the INTT and ϕ_i is the azimuthal angle of the i -th cluster.

Black line in Fig. 1 (right) shows the observed ψ_2^{obs} distribution, determined using the 200 GeV Au + Au data without the magnetic field. The ψ distribution should be flat because the RP is randomly oriented. There is a strong oscillation that supposes to be caused by the effects of detector acceptance and the beam offset from the center of the INTT barrels. These detector effects need to be corrected for further analysis. The red line in the Fig. 1 (right) shows re-centered ψ distribution (ψ_2^{rec}) in which the beam effect has been corrected.¹⁾ ψ_2^{rec} is calculated by subtracting the mean of the ψ_2^{obs} distribution from ψ_2^{obs} for each event and dividing it by the standard deviation of the distribution. ψ_2^{rec} still has a non-flat distortion effect which is caused by the detector acceptance. To correct the distortion effect, the flattered ψ_2^{flat} is determined by

$$\begin{aligned} \psi_2^{flat} &= \psi_2^{rec} + \Delta\psi \\ \Delta\psi &= \sum_{k=1}^8 (A_k \cos 2k\psi_2^{rec} + B_k \sin 2k\psi_2^{rec}) \\ A_k &= -\frac{2}{k} \langle \sin 2k\psi_2^{rec} \rangle, \quad B_k = \frac{2}{k} \langle \cos 2k\psi_2^{rec} \rangle \end{aligned} \quad (2)$$

where k represents the harmonic order in the Fourier series expansion, which in this analysis sums up to the 8th order in this analysis. The blue line in Fig. 1 (right) shows that the RP has been successfully calibrated by correcting the distortion effect.

Reference

- 1) I. Selyuzhenkov, S. Voloshin, Phys. Rev. C **77**, 034904 (2008).